TOWARDS THE UNBOUND STELLAR POPULATION IN THE SLOAN DIGITAL SKY SURVEY

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To my family.

TABLE OF CONTENTS

LIST	OF TAB	LES	v
LIST	OF FIGU	URES	vii
Chapt	er		
I.	INTRO	DUCTION	1
	1.1.	Historical Perspectives on Galactic Structure 1.1.1. 1.1.1. Stellar Halo 1.1.2. Disk 1.1.3 Bulge	$2 \\ 2 \\ 6 \\ 9$
	1.2.	Current View of Galactic Structure 1.2.1. Bulge and Supermassive Black Hole 1.2.2. Thin & Thick Disks 1.2.3. Stellar Halo 1.2.4. Dark Halo	13 13 14 17
	1.3.	1.2.4. Dark Halo The Sloan Digital Sky Survey 1.3.1. Sloan Extension for Galactic Understanding and Explo-	$\frac{18}{32}$
	1.4.	ration	33 34
II.	IDENT DISTAN	IFYING HIGH METALLICITY M GIANTS AT INTRAGROUP NCES WITH SDSS	38
	 2.1. 2.2. 2.3. 2.4. 2.5. 2.6. 	Abstract	38 38 42 47 49 50 52
III.	HYPER SAMPL	VELOCITY STAR CANDIDATES IN THE SEGUE G & K DWARF E	59
	3.1. 3.2. 3.3. 3.4.	AbstractIntroductionIntroductionIntroductionIdentifying HVS CandidatesIntroductionEstimating the Fidelity of Our CandidatesIntroduction3.4.1.Proper-Motion Quality Cuts3.4.2.Monte Carlo Sampling	59 60 61 64 64 67

	3.5. Orbits of HVS Candidates 6 3.5.1. Galaxy Model 6 3.5.2. Origins 7 3.6. Discussion 7 3.6.1. Chemical Tracing 7 3.6.2. Alternative Origins 7 3.6.3. Follow-up Analysis 7 3.6.4. Constraints on the Initial Mass Function 7	9 9 12 2 2 6 7 7
	3.7. Summary and Conclusions	9
IV.	WORK IN PROGRESS	2
	4.1. A Global Search for G/K-type Hypervelocity Stars 8 4.1.1. Motivation 8 4.1.2. Candidate Selection 8 4.1.3. Preliminary Results 8 4.1.4. Discussion 8 4.2. F-type Hypervelocity Star Candidates 8 4.2.1. Motivation 8 4.2.2. Candidate Selection 8 4.2.3. Preliminary Results 8 4.2.4. Discussion 8	2234566778
V.	CONCLUSION	5
Append	DIX	
А.	ESTIMATION OF EXPECTED HVS VIA THE BINARY DISRUPTION MECHANISM	7
В.	EXPANSION ON THE EXPECTED IGS CALCULATION 9	9
С.	IGS CANDIDATES REMAINING AFTER ALL CRITERIA CUTS 10	0
D.	EXPANSION ON THE EXPECTED HVS CALCULATIONS	7
BIBLIO	GRAPHY	8

LIST OF TABLES

TABLE		Pac	ΞE
II.1.	IGS candidates remaining after all criteria cuts [Complete version pro- vided in Appendix C]		53
III.1.	Stellar and kinematic parameters for the 20 HVS candidates	•	75
C.1.	Complete and unabridged	. 1	01

LIST OF FIGURES

FIGURE		Page	
1.1.	CMD of 47 Tuc		10
1.2.	UV excess vs. Eccentricity, Velocity, and Angular Momentum		11
1.3.	Density of F/G-type stars vs. Distance		12
1.4.	Schematic of Milky Way Structure		12
1.5.	Disruption of a Satellite		22
1.6.	Galaxy Disk During a Merger		23
1.7.	Galactic Disks of varying Heights		24
1.8.	Evolution of a Clumpy Galaxy		25
1.9.	Number of Stars vs. Rotational Velocity		26
1.10.	Metallicity of Disk Stars		27
1.11.	Stellar Density vs. Scale Height		28
1.12.	Mass Density of Disk Stars		28
1.13.	Rotational Velocity of Stars with Metallicity $<$ -2		29
1.14.	Disruption of Satellite Galaxies into Streams		30
1.15.	SDSS Map of Halo Structure		31
1.16.	COBE Image of Milky Way		31
1.17.	Intracluster Light in Virgo Cluster		36
1.18.	Snapshot of Hills Mechanism		36
1.19.	Ejection Velocity vs. Kick Velocity		37
2.1.	CCDs of Covey et al. (2007) sample data		40
2.2.	CCDs including L & T Dwarfs and SDSS Stripe 82 \hdots		43
2.3.	Distinctions between M giants and M dwarfs		55

2.4.	CCD with M giant Standards from the BPGS Atlas	56
2.5.	Comparison Spectra	57
2.6.	Position of IGS Candidates vs. HVS	58
3.1.	Transverse vs. Radial Velocity for HVS Candidates	62
3.2.	Velocity Distribution for Most and Least Bound HVS Candidates $\ . \ . \ .$	73
3.3.	Comparison of Orbit in Spherical and Triaxial Halos	73
3.4.	Orbits of 7 Most Unbound HVS Candidates	74
3.5.	Metallicity Distribution of HVS Candidates	81
4.1.	Positions of Preliminary Global G/K-dwarf HVS Candidates	89
4.2.	Transverse vs. Radial Velocity of Global G/K-dwarf HVS Candidates	90
4.3.	Velocity Distribution of Global G/K-dwarf HVS Candidates $\ . \ . \ . \ .$	91
4.4.	Metallicity Distribution of Global G/K-dwarf HVS Candidates	92
4.5.	Velocity Distribution of Preliminary F-dwarf HVS Candidates	93
4.6.	Metallicity Distribution of Preliminary F-dwarf HVS Candidates	94

Chapter I

INTRODUCTION

In this dissertation, I present my work to identify and understand the unbound stellar population within our Milky Way using data from the Sloan Digital Sky Survey. This investigation is a single facet in the broader field of Galactic structure research. As I will elaborate on in this document, this unique population can help us understand the dynamics and stellar content of the Galactic center, the mass of the dark matter halo, and the evolution of the Milky Way.

The study of Galactic structure isn't simply compiling a map of the different components that comprise the Milky Way. While this is important in itself, it is only the first stepping stone to creating a complete understanding of galaxy formation, star formation, and the distribution of dark matter through the Universe (Rockosi et al., 2009). Studying the density, kinematics, and chemistry of stars and their preferred locations within the Galaxy introduces new possibilities for understanding the formation and evolutionary histories therein (e.g., Freeman & Bland-Hawthorn, 2002; Rockosi et al., 2009).

To place this work in the context of Galactic Astronomy, I need to first provide some background information. It is the goal of this chapter to paint a picture of our current understanding of the structure of the Milky Way and how that view has changed over time. To that end, I will begin, in Section 1.1, with a description of Galactic structure as it was roughly 20 years ago and then move to a presentation of our current perception in Section 1.2. In Section 1.3, I will discuss the Sloan Digital Sky Survey and how it has had a tremendous impact on our understanding of Milky Way structure. Finally, in Section 1.4, I will discuss ways in which understanding the unbound stellar population can provide further insight into the larger context of Galactic structure, formation, and evolution.

1.1 Historical Perspectives on Galactic Structure

Astronomical studies focusing on the evolutionary history of the Milky Way only became feasible during the 1950's when theoretical and observational advances were finally able to join structure, kinematics, chemistry, and age into a single context (Majewski, 1993). The first development linking these parameters into a single model of the Milky Way, and perhaps the most influential, came in the form of the Eggen, Lynden-Bell, and Sandage model, discussed in detail in Section 1.1.1 below (Eggen et al., 1962; Majewski, 1993). A very simplified view of the Milky Way consists of a bulge and disk surrounded by a diffuse halo. The following sections will elaborate on this picture.

1.1.1 Stellar Halo

Before the 1990's the Milky Way halo was believed to be spheroidal, slowly rotating, and comprised of metal weak¹ stars, [Fe/H] < -1 (Freeman, 1987). Its rotational velocity,

¹The metallicity of stars, typically denoted as [Fe/H], is a measurement of the amount of heavy metals (that is, heavier than Helium) in a star's composition. Metallicity is reported with respect to that of our Sun. Therefore, stars are described as either having more (metal rich) or less (metal poor or weak) metals than the Sun. Metallicity is also a signature of age as higher metallicity implies a star was formed more recently.

 v_{rot} , was measured to be roughly 40 km/s, with a velocity dispersion² of $\sigma \sim 100$ km/s (Freeman, 1987). This halo follows an $r^{-3.5}$ density profile (Freeman, 1987).

The globular clusters that are found in the halo are also metal weak with [Fe/H] < -1, and with similar velocities, $v_{rot} \sim 50$ km/s and $\sigma \sim 110$ km/s, to that of the halo (Freeman, 1987). In addition, about 30 globular clusters were known to lie on retrograde orbits with $v_{rot} \sim -70$ km/s (Freeman, 1987).

Globular clusters are believed to be some of the oldest objects in the Universe, and therefore provide an ideal test-bed for the earliest stages of galaxy formation (e.g., Binney & Merrifield, 1998; Freeman & Bland-Hawthorn, 2002). Figure 1.1 shows the colormagnitude diagram of the nearby globular cluster 47 Tuc. The sharp knee at a magnitude³ of ~ 17 indicates the age of the cluster since the more luminous stars have evolved off the main sequence⁴. Because these clusters are the oldest observable stellar objects, they provide a lower limit on the age of the Universe and the timescale over which galaxies form.

The Eggen Lynden-Bell and Sandage (ELS) model of 1962 (Eggen et al., 1962) describes the formation of our Galaxy from a protogalactic cloud of intergalactic material. To summarize, a large cloud of rotating gas, about 100 kpc⁵ in radius, began its collapse roughly 10¹⁰ years ago. From the collapsing gas the first stars formed into globular clusters, and as the material continued to collapse more stars formed in a rotating, thin disk.

²Velocity dispersion refers to the range of velocities about the mean. Therefore, a small velocity dispersion means all stars in the system have roughly the same velocity.

³Magnitude is a measure of brightness. Greater magnitude values correspond to fainter objects, and likewise smaller values correspond to brighter objects.

⁴The main sequence is a feature on color-magnitude diagrams running diagonally from bright and blue to faint and red. Stars that are fusing Hydrogen into Helium in their cores live along the main sequence.

 $^{{}^{5}1 \}text{ kpc} = 3 \times 10^{19} \text{ m}.$

After time enough for the first O and B type $stars^6$ to fuse their core gas into heavier elements and undergo supernovae eruptions that would enrich the remaining gas at the end of their lifetimes, a second generation of stars began forming in the disk of the Galaxy (Eggen et al., 1962).

This model was theorized in order to explain correlations between ultraviolet (UV) excess⁷ and 1) orbital eccentricity, 2) W velocity (perpendicular to the plane of the Galaxy), or z_{max} (height above or below the plane), and 3) angular momentum in a sample 221 dwarf stars in the solar neighborhood (Eggen et al., 1962; Majewski, 1993). Ultraviolet excess is a signature of very old, metal poor stars. A correlation between this and any other stellar property may provide information about how these stars, and therefore the Galaxy, formed.

The first correlation, between UV excess and orbital eccentricity, is shown here in the top-left panel of Figure 1.2. This figure shows stars with large UV excess have highly eccentric orbits although stars exhibiting less UV excess live in nearly circular orbits (Eggen et al., 1962).

The second correlation is between UV excess and maximum distance above or below the plane of the disk, z_{max} . Stars with little UV excess remain in or near the plane of the disk. However, stars with high UV excess have a wider variety of z_{max} values, ranging from 0 to 10 kpc, shown in the top right panel of Figure 1.2. These two correlations,

⁶The spectral classification of stars follows the sequence: OBAFGKM. In general, the most massive, hottest, and bluest stars are denoted as O-type and transition down to the least massive, coolest, reddest M-type stars.

⁷The UV excess, $\delta(U-B)$, of a star is measured in the ultraviolet-blue wavelengths of a star's spectrum. It manifests itself in very metal poor stars that have few elemental absorption features in this region of the spectrum resulting in excess emission.

together, suggest that there is also a correlation between UV excess and the age of the stars– more UV excess in older stars, less UV excess in younger stars (Eggen et al., 1962).

Finally, the third correlation, between UV excess and angular momentum, shown in the bottom panel of Figure 1.2, suggests that older stars with more UV excess have smaller angular momenta than the stars with less UV excess (Eggen et al., 1962).

The combined results of these data leads to a concise description of the first stars in the Galaxy. These stars tend to live on highly eccentric orbits, can be observed at a wide range of heights above and below the plane, and generally have low angular momenta (Eggen et al., 1962). This new information made possible a complete theory of Galactic formation that describes how these stars came to have their observed properties.

In this model, the Galaxy forms from a smooth, hot cloud of rotating gas. As the gas begins to cool and collapse, stars begin forming in globular clusters. Then, with continued cooling, the remaining gas begins rotating faster to conserve angular momentum and collapses into a flat disk and central spheroidal component. Therefore, the stars formed during collapse have highly eccentric orbits, although the stars formed from the gas in the disk have more circular orbits. As those stars enrich the interstellar medium at the end of their lifetimes, the newly formed, younger stars will exhibit higher metallicities and lower UV excesses.

The stars used in this study were primarily chosen to be nearby, solar neighborhood, stars due to their high proper motions. Therefore, this sample, and this model, neglects metal poor stars on low eccentricity orbits (Majewski, 1993; Bond, 1970). Also, the ELS model neither accounts for nor explains the number of retrograde orbits seen in the halo of the Galaxy (e.g., Majewski, 1993; Larson, 1969; Yoshii & Saio, 1979). The ELS model was first challenged by Searle & Zinn in 1978 who noted a range in metal abundances in globular clusters independent of distances from the Galactic Center (Searle & Zinn, 1978) and suggested that the halo was built over an extended period of time from $\sim 10^8 M_{\odot}$ independent fragments (Searle & Zinn, 1978; Freeman & Bland-Hawthorn, 2002).

The subsequent Yoshii & Saio model (Yoshii & Saio, 1979) incorporated low metallicity, low eccentricity stars and retrograde orbits by allowing a slower collapse of a clumpy protogalaxy (e.g., Majewski, 1993; von Weizsäcker, 1951; Oort, 1958) over several gigayears. This model lends itself easily to the (perhaps, preferred) notion of a possible tumultuous Galactic evolution due to a series of hierarchical mergers (e.g., Majewski, 1993; Tinsley & Larson, 1979).

1.1.2 Disk

The Galactic disk is flat, rotating at 220 km/s with $\sigma \sim 20$ km/s, and comprised of the majority ($6 \times 10^{10} M_{\odot}$) of the stellar mass (Freeman, 1987). The disk also contains relatively high metallicity stars, with [Fe/H] > -0.5 (Freeman, 1987). It has an overall scale height 300-350 kpc and scale length 3.5-5.5 kpc (Freeman, 1987). A spiral arm pattern was also discovered in the Milky Way in the early 1960's. There are at least 4 arms beyond 3 kpc from the Galactic center, and each is roughly 2 kpc from the next (Rougoor & Oort, 1960). In addition, there is a thin layer of Hydrogen gas also rotating at roughly 220 km/s (Rougoor & Oort, 1960). This disk of gas is also known to "warp" at the edges, rising roughly 700 kpc above the Galactic plane on one side and similarly below the plane on the opposite side (Rougoor & Oort, 1960).

It was well known that the scale height of the disk varied with stellar population

(Mihalas & Binney, 1981). For example, young O and B stars lie very near the plane within ~ 50 pc, interstellar gas and dust have a scale height ~ 100 pc, and old M stars and white dwarfs have scale heights $\sim 500 - 1000$ pc (Mihalas & Binney, 1981).

In 1983 Gilmore & Reid used the photometric parallax of a sample of 12,500 stars complete to an I-band (near-infrared) magnitude of 18, and all within \sim 18 degrees of the South Galactic Pole. With these data they applied exponential fits of varying scale heights.

They show, in Figure 1.3, the spatial density of two subsamples of their stars (in the magnitude range $4 \leq M_V \leq 5$ and in the range $5 < M_V \leq 6$, correlating to F-G type stars and G-K type stars, respectively) as a function of distance from the galactic plane. The solid line depicts an exponential decrease with a scale height of 300 pc. This single exponential no longer fits at distances greater than 1000 pc. Therefore, they fit a second exponential, the dashed line, for stars with z > 1 kpc, corresponding to a scale height of ~ 1500 kpc for both subsamples of stars (Gilmore & Reid, 1983; Freeman, 1987). Thus, Gilmore & Reid introduced an additional "thick disk" component, convincingly unique from the thin disk population, to the existing two-component, bulge-disk, model of the Milky Way.

They propose an explanation for this apparently distinct population of thick disk stars, comprising only $\sim 2\%$ of the stars in the solar neighborhood, $\sim 10^9 M_{\odot}$, as being Galactic bulge stars that have felt the gravitational pull of the massive thin disk. This is largely due to the chemical similarities this population of stars has with the spheroidal component of our Galaxy (Gilmore & Reid, 1983).

This led to many models describing the formation of this "thick disk" population, also

referred to as "extended disk," "high velocity disk," or "flattened spheroid" depending on the preferred formation criteria (Majewski, 1993). These thick disk models tend to fall into one of two camps– "top down" or pre-thin disk, or "bottom up" or post-thin disk (Majewski, 1993).

The "top down" models form the thick disk population prior to the formation of the thin disk (Majewski, 1993). In these pre-thin disk models the thick disk is often treated as a transitionary stage between the collapse of the halo and the formation of the thin disk (e.g., Majewski, 1993; Larson, 1976; Gilmore, 1984)

The "bottom up" models typically involve some interaction or evolution of the thin disk that results in the formation of the thick disk (Majewski, 1993). These models draw from the idea that processes that can increase the velocity dispersion of, or kinematically heat, the thin disk stars may contribute to the formation of a thick disk population that is smoothly linked to the thin disk (e.g., Majewski, 1993). Thus, we would expect to see gradients in age, metallicity, and kinematics within the disks, as older, metal poor stars would have undergone more heating events. Three main heating mechanisms considered in these post-thin disk models include 1) interactions with $\sim 10^6 M_{\odot}$ molecular clouds, 2) disk instabilities such as spiral arms, and 3) interactions with fast moving, massive objects like dwarf galaxies or supermassive black holes (Majewski, 1993).

The thick disk component contains stars with [Fe/H] < -1, and has a velocity dispersion roughly double that of the thin disk, ~ 40 km/s (Freeman, 1987). 1.1.3 Bulge

The central region of the Galaxy is comprised of a spheroidal stellar bulge. These stars are metal rich with [Fe/H] > -1 and rotate at ~ 200 km/s (Freeman, 1987; Rougoor & Oort, 1960). The bulge is slightly elongated and has a radius about 2.5 kpc (Freeman, 1987).

The idea of a central bar was first introduced in the early 1960's based on observations of non-circular stellar motions near the center of the Galaxy (Rougoor & Oort, 1960; de Vaucouleurs, 1964). It was suggested to be roughly 3 kpc in radius and oriented at about 30° from our line-of-sight (Rougoor & Oort, 1960; de Vaucouleurs, 1964). However, the bar wasn't studied in detail until the 1990's, and therefore refer to Section 1.2.1 for a more detailed discussion of its properties.

At this point, the generally adopted idea of what components make up our Galaxy can be described by the schematic shown in Figure 1.4. This cartoon representation shows a large dark halo extending out to about 100 kpc, in black, and smaller stellar halo, in yellow, each with the old, metal poor globular clusters distributed throughout, marked by the filled reddish dots. It also shows a thin disk, light blue, surrounded by a puffy thick disk, dark blue, with radii about 15 kpc. Finally, in the center of the Galaxy is a stellar bulge, shown in yellow (Freeman & Bland-Hawthorn, 2002).



Figure 1.1: The color magnitude diagram of globular cluster 47 Tuc. The location of the knee directly corresponds to the age of the cluster, estimated to be approximately 12 Gyr. (Image courtesy of Michael Richmond at http://spiff.rit.edu/classes/phys301/lectures/mw_size.html.)



Figure 1.2: Top left: Ultraviolet excess as a function of orbital eccentricity. This shows stars with a greater UV excess tend to live on highly eccentric orbits, whereas stars with less UV emission have more circular orbits. Top right: W velocity as a function of UV excess. Also plotted on the right-hand vertical axis is Z_{max} above the plane. Stars with greater UV excess exhibit a much larger range of possible W and Z_{max} values. Bottom: Angular momentum as a function of UV excess, showing stars with more UV excess have little or no angular momentum. (Eggen et al., 1962)



Figure 1.3: Left: The spatial density of stars as a function of distance away from the plane for a subsample of stars with $4 \leq M_V \leq 5$ (late F- to early G-type stars). The straight lines represent exponential fits to the data. A single exponential no longer fits at distances greater than about 1.5 kpc introducing the need for a second disk component. Right: Same as left, but for a subsample of stars with $5 < M_V \leq 6$ (late G- to early K-type stars). (Gilmore & Reid, 1983)



Figure 1.4: A cartoon depicting the major components of the Milky Way. The bulge is shown by the central circle in yellow. The thin disk is in light blue and the thick disk around it is shown in darker blue. Surrounding the bulge and disks is the stellar halo in yellow. Finally, the dark halo is shown in black and the globular clusters by filled circles. (Freeman & Bland-Hawthorn, 2002)

1.2 Current View of Galactic Structure

1.2.1 Bulge and Supermassive Black Hole

Although a bar in the central region of the Milky Way was first introduced in the 1960's (Rougoor & Oort, 1960; de Vaucouleurs, 1964) with the discovery of non-circular motions in the center of the Galaxy, it wasn't until the early 1990's that it became generally established that the Milky Way is a barred spiral. The existence of a central bar has been confirmed by kinematic, photometric, and star count techniques (López-Corredoira et al., 2007; Morris & Serabyn, 1996). Its mass is ~ $10^{10}M_{\odot}$, extends roughly ~ 5 kpc from the center, rotates at about 200 km/s, and is misaligned with our line-of-sight to the Galactic center at an angle of ~ 45° (López-Corredoira et al., 2007; Morris & Serabyn, 1960).

Recently, new constraints have been made on the properties of the supermassive black hole at the center of the Galaxy. This is largely due to new precision radial velocity measurements of the high velocity stars very near the black hole (Ghez et al., 2008). Ghez et al. (2008) determined the distance to the black hole is 8.0 ± 0.6 kpc, the mass is $4.1 \pm 0.6 \times 10^6 M_{\odot}$, and a radial velocity consistent with zero.

The central parsec around the supermassive black hole is an important region of interest. In particular, the central density profile may serve as a simple test for the existence of a central black hole (Do et al., 2009). Theory states that as the orbits of stars bring them close to a central black hole, the star is destroyed and its energy is balanced by increasing the density of the central region surrounding the black hole. This results in a "cuspy" density profile, $\sim r^{-7/4}$, although in the absence of a black hole a flat or "cored" density profile is expected, ~ $r^{-1/2}$ (Bahcall & Wolf, 1976; Schödel et al., 2007; Do et al., 2009). However, studies are finding a significant lack of old, low mass, stars near the center of the Galaxy, resulting in a flat, maybe declining, density profile of ~ $r^{-0.3}$ (Do et al., 2009; Bartko et al., 2010; Schödel et al., 2007).

1.2.2 Thin & Thick Disks

The existence and current observations of the thick disk can be explained by a number of formation theories. Understanding how the thick disk formed will help to constrain theories of Galactic formation and interaction history.

One possible formation mechanism for the thick disk is through accretion of satellite galaxies. The simulations of Abadi et al. (2003), shown in Figure 1.5, imply that the majority of stars in the thick disk, $\sim 75\%$, are the debris of tidally stripped satellites. This theory suggests that the thick disk formed completely independent from a pre-existing thin disk.

The thick disk may have alternatively formed by the kinematic heating of the thin disk through minor mergers. Villalobos & Helmi (2008) consider this thick disk formation mechanism in their simulations of galaxy mergers. They merge satellites with $\sim 10\%$ of the mass of a host galaxy which already contains a disk, seen in Figure 1.6. From these mergers they find that thick disks form with similar kinematics to what is observed in the Milky Way (Villalobos & Helmi, 2008).

Another possible formation mechanism of the thick disk comes in the form of scattering, or radial migration, of stars due to asymmetric structures in the Galaxy such as the bulge or spiral arms. One study that investigates this theory was performed by Schönrich & Binney (2009). They used the chemical abundances of solar neighborhood stars and their space motions to evolve their orbits and recover the signatures of both thin and thick disks in the absence of any interactions with external objects (Schönrich & Binney, 2009). A similar study by Bird et al. (2011) simulates the effects of a bar and spiral structure compared to the effects of satellites on radial migration in disk galaxies. Figure 1.7 shows the stellar distribution in the initial disks (top with scale height, $z_d = 200$ pc, bottom with $z_d = 400$ pc) on the left, and the stellar distributions after 2.5 Gyr in isolation and in the presence of satellites, center and right panels respectively. Although a bar and noticeable spiral structure forms in the thinner disk, satellite interactions have a much more pronounced effect on the migration of the disk stars (Bird et al., 2011).

A competing theory is that the disks formed *in situ* from a 'clumpy' distribution of material as the result of a major gas rich merger (Bournaud et al., 2007). The simulations of Bournaud et al. (2007) show, in Figure 1.8, in addition to a clumpy primordial galaxy being required to form an exponential disk, the time it takes to form the disk from clumpy material is much shorter than the time it takes spiral arms to disrupt a thin disk into an exponential disk.

Carollo et al. (2010), in addition to providing support for the two-component stellar halo, see Section 1.2.3, proposed an additional metal-weak thick-disk is needed to accurately describe the thick disk stellar population. In this study, they employ the distribution of rotational velocity, V_{ϕ} , and fit a minimum number of gaussian curves to sufficiently describe the data. As a result, they claim a metal-weak thick-disk is required with V_{ϕ} of about 100-150 km/s and a peak metallicity [Fe/H] = -1.3, compared with the 'regular' thick disk which has $V_{\phi} \sim 182$ km/s and a peak metallicity [Fe/H] = -0.8 (Carollo et al., 2010). This result is depicted in Figure 1.9, where the histograms represent the number of stars with a particular rotational velocity, the red curves represent each of the components used in the model, and the blue lines are the sum of all components. The left panel shows the distribution for a subsample of stars 1-2 kpc above the plane in varying metallicity bins. Similarly, the right panel shows stars in a distance slice between 2 kpc and 4 kpc. As shown by this figure, a single component is not sufficient to recover the distribution of observed rotational velocities (Carollo et al., 2010).

In 2011, a new study has challenged the determination of the thick disk scale length, believed to be roughly the same as that of the thin disk, ~ 3 kpc. Bensby et al. (2011) add the metallicity distribution of a sample of 20 outer disk giants, further from the center than the Sun, to a sample of ~ 40 inner disk giants, between the bulge and the Sun. They then compare this to the sample of solar neighborhood stars of Alves-Brito et al. (2010) where they were able to distinguish thin disk from thick disk populations based on their metallicity determinations. What they found is a significant lack of stars with thick disk chemistry at large distances. This can be seen clearly in Figure 1.10, where the left panel shows the 40 inner disk giants, the center shows the solar neighborhood, and the right panel shows the 20 outer disk giants (Bensby et al., 2011). In each of these panels the red and blue lines show the metallicity distributions of the thick and thin disks, respectively, as they were determined from the solar neighborhood sample (Bensby et al., 2011).

A mono-abundance study of the SEGUE G dwarf sample revealed that each subpopulation can be described by a single exponential function (Bovy et al., 2012a). This is shown in Figure 1.11, where stellar density is plotted against vertical scale height for sub-populations in bins of metallicity and alpha-abundance; the black line shows the total stellar density. Hence, Bovy et al. (2012a) suggest that "thick disk" stars do not comprise a distinctly separate component from the thin disk. Similarly, Bird et al. (2013) considered mono-age sub-populations in hydrodynamic simulations of Milky Way-like disk galaxies. Figure 1.12 shows the vertical mass density profiles for the mono-age populations in the inner disk, Solar neighborhood, and outer disk, respectively. Each sub-population can be described by a single exponential function. However, the composite of all sub-populations to form the total vertical mass density profile requires a fit by two exponential profiles (Bird et al., 2013)– apparently explaining the origin of the thin disk/thick disk dichotomy.

In summary, there is still much that is undetermined about the Galactic disk despite that it is the dominant stellar component in the solar neighborhood. The idea that there is a "thick disk" in the Milky Way is generally accepted. However, the nature, formation, and even the exact definition of what constitutes the thick disk is the subject of much debate.

1.2.3 Stellar Halo

In the previous section, the halo component of the Galaxy refers to a smooth stellar halo hosting low metallicity globular clusters (Eggen et al., 1962). This component only contains ~1% of the stellar mass, ~ $10^9 M_{\odot}$, in the Milky Way and has close to zero angular momentum compared to the disk and bulge components (Morrison, 1993; Freeman, 1987; Freeman & Bland-Hawthorn, 2002). Now, we know this metal poor stellar halo component is riddled with substructure, and it is believed that it formed, in part, from the accretion of small, somewhat evolved, metal poor satellite galaxies (Searle & Zinn, 1978; Freeman & Bland-Hawthorn, 2002). Carollo et al. (2007, 2010) added to the substructure of the halo by providing strong evidence for a two-component stellar halo– an inner halo and an outer halo shown in Figure 1.13. These two halos differ kinematically, chemically, and in spatial density suggesting that they formed from entirely different mechanisms. (Carollo et al., 2007, 2010).

The inner halo exhibits a small net prograde rotation of 0 to 50 km/s, large orbital eccentricities, a metallicity peak at [Fe/H] = -1.6, and dominates up to distances of about 15 kpc. The outer halo, on the other hand, shows a net retrograde motion of -40 to -70 km/s, low eccentricities, a peak metallicity around [Fe/H] = -2.2, and dominates beyond 15 kpc (Carollo et al., 2007, 2010).

It is proposed that the inner halo formed from dissipational mergers (containing gas) of low-mass clumps into higher-mass clumps. Then the nearly radial merger of the resulting massive clumps resulted in highly eccentric orbits and star formation in the recently merged gas created the higher metallicity stars. Whereas, the outer halo formed through random dissipationless (gasless) mergers of low mass, low metallicity satellites (Carollo et al., 2007). It is the random orientation of the mergers that help to explain the slight retrograde motion observed in this region.

1.2.4 Dark Halo

In addition to the stellar halo component of the Galaxy as previously discussed, we know now that a massive dark matter halo surrounds the Milky Way. Navarro et al. (1997) have theorized that all dark matter halos follow the same density profile regardless of halo mass. The profile follows r^{-1} at small radii, and r^{-3} at large radii (Navarro et al., 1997).

Dark matter cosmological simulations predict that there should be many more dark matter satellites surrounding the Milky Way than the number of dwarf galaxies that had been discovered in the early 1990's (e.g., Simon & Geha, 2007; Kauffmann et al., 1993). This discrepancy has become known as the "missing satellite" problem. With the advent of the Sloan Digital Sky Survey (SDSS, see Section 1.3) the number of satellite galaxies around the Milky Way, and in the Local Group, has increased to about 50 (York et al., 2000a; Simon & Geha, 2007). Although this does not solve the "missing satellite" problem, it does dramatically alter our understanding of the Milky Way's local environment. Membership determinations of satellite galaxies require reliable distances and radial velocities (Freeman & Bland-Hawthorn, 2002). Given the intrinsic low luminosity of the objects and the current limitations of our surveys (SDSS has only imaged $\sim 1/5$ of the sky), it isn't unreasonable that many more dwarf galaxies within the Local Group simply haven't been discovered yet (Simon & Geha, 2007).

From this follows the idea that some of the satellites may be "missing" because they have been disrupted by interactions with the Galaxy. This question was addressed by Morrison et al. (2000) and the "Spaghetti" Survey. This survey was designed to map the substructure of the halo from the velocities of distant giants⁸, and ultimately to determine the fraction of mass in the halo that was accreted from satellites (Harding et al., 2001).

With this survey, in combination with modeling the orbits of the halo giant stars, they were able to recover the stripped remains of accreted satellites in the form of long streams surrounding the Galaxy (Harding et al., 2001). An example of this is shown in Figure

⁸Giant here refers to the stellar evolutionary stage that occurs after a star has fused all of its core Hydrogen into Helium. In this stage the accelerated fusion rate in the shells surrounding the core cause the star to expand and cool dramatically.

1.14, which shows the X-Y and Z-Y projections of three satellites after 1 Gyr, larger points, and 10 Gyr, smaller points, through its orbit around the Milky Way (Harding et al., 2001). Similarly, shown in Figure 1.15, images from the SDSS have uncovered a "field of streams" throughout the halo (Belokurov et al., 2006a). Typically, the mass of one of these streams will be in the range $10^7 - 10^9 M_{\odot}$ (e.g., Law & Majewski, 2010; Harding et al., 2001)

Since the gravitational potential of the Galaxy is dominated by dark matter at large radii, observations of the large tidal features in the halo can be used to constrain the shape, mass, and orientation of the dark matter in the Milky Way (Law & Majewski, 2010). Law & Majewski (2010) model the Sagittarius stream in different Milky Way dark matter halos, varying the overall shape of the potential. Axisymmetric halos are not able to reproduce the observed positions and distances of the stream. However, they suggest a triaxial halo with minor/major axis ratio, c/a = 0.72, and intermediate/major axis ratio, b/a = 0.99, to explain the observed structure.

The current picture of the Milky Way consists of a central supermassive black hole, surrounded by a somewhat box-shaped bulge. The majority of the stellar material is concentrated in a thin disk, which also contains spiral arms and a thin layer of gas that warps at the very edges. Encasing the thin disk is a diffuse thick disk population. These components all exist within an inner, prograde, stellar halo, an outer, retrograde, stellar halo, and a massive dark matter halo housing the extremely metal poor globular clusters, streams, and satellite galaxies. An image of the Milky Way from the Cosmic Background Explorer (COBE) satellite can be seen in Figure 1.16. The figure clearly shows the boxy shape of the bulge and the thin disk. More difficult to discern are the diffuse thick disk, thin Hydrogen layer, and halo material.

This picture, although significantly more elaborate that what was described in Section 1.1.3, is still evolving. The desire to understand the center, disk(s), halo, formation, and evolutionary history of the Milky Way is the driving factor for much of on-going research today. This is especially exciting with more and more data becoming available. The Sloan Digital Sky Survey (SDSS; York et al., 2000a) is only one example of a large survey dedicated to such topics. Its contributions are described in Section 1.3.



Figure 1.5: The disruption of a satellite over time seen edge-on with respect to the main galaxy (not shown). A significant fraction of the stars in the satellite contribute to a thick disk population. (Abadi et al., 2003)



Figure 1.6: The disk of a host galaxy during a merger with a satellite (not shown). Over time the merger causes the existing disk to thicken. (Villalobos & Helmi, 2008)



Figure 1.7: Left: Stellar distribution of galactic disks with different scale heights ($z_d = 200$ pc in the top panels and $z_d = 400$ pc in the bottom panels). The disk in the bottom panels more closely resembles the Milky Way disk. Middle: The disks after evolving in isolation for 2.5 Gyr. A bar and slight spiral structure form in the $z_d = 200$ pc disk. Right: The disks after evolving under satellite perturbations for 2.5 Gyr. Bars, spiral structure, and disk thickening occur in both cases. (Bird et al., 2011)



Figure 1.8: Face-on views of the evolution of a clumpy galaxy. The result is a spiral disk galaxy with central bulge and 2-component disk (not shown in this projection). (Bournaud et al., 2007)



Figure 1.9: Left: The number of stars as a function of rotational velocity in varying metallicity bins for stars between 1-2 kpc from the plane, with the combined model for the distribution of velocities depicted by the blue line. The red lines are the Gaussian distributions for each component in the model. Right: Similarly, for stars between 2-4 kpc from the plane. (Carollo et al., 2010)



Figure 1.10: The center panel shows the metallicity of solar neighborhood disk stars. Red points are thick disk stars, fit by the red line, and blue points are thin disk stars, fit by the blue line. The left and right panels show the inner disk (open circles) and outer disk (filled circles) samples, respectively. Each is compared to the trends determined by the solar neighborhood sample. There is a distinct lack of stars following the thick disk metallicity trend at distances greater than 9 kpc. (Bensby et al., 2011)



Figure 1.11: Stellar density versus scale height for mono-abundance sub-populations of SEGUE G dwarfs. Each population can be described by a single exponential function. The total stellar density is represented by the solid black line.



Figure 1.12: Each colored line represents the vertical mass density of a mono-age population in the inner disk (left), Solar neighborhood (center), and outer disk (right). Each mono-age population can be described by a single exponential function. However, the composite of all populations (black curve) requires two exponential profiles (grey dashed lines) (Bird et al., 2013).



Figure 1.13: The number of stars with [Fe/H] < -2 and $Z_{max} > 5$ as a function of rotational velocity, with the combined model for the distribution of velocities depicted by the blue line. The red lines are the Gaussian distributions for each component in the model. The parameter k is the number of components used in the model. When k=1 the fit is not descriptive of the data. However, k=2 does much better at fitting the data with only slight improvement at k=3, signifying that the 2-component model is accurate. (Carollo et al., 2010)


Figure 1.14: The X-Y and Z-Y projections of three separate satellites, identified numerically in the upper-right of each panel, at 1 Gyr and 10 Gyr (larger and smaller points, respectively–although, difficult to distinguish) throughout their orbit. This figure shows the disruption of satellites over time into large streams surrounding the Galaxy. (Harding et al., 2001)



Figure 1.15: A map of the sky as seen by SDSS images showing the structure of the halo. Circles denote the location of globular clusters and satellite galaxies discovered by the SDSS. Other structures in this image include the Sagittarius, Orphan, and Monoceros streams. (Image courtesy of http://www.sdss.org (????).)



Figure 1.16: This image of the Milky Way was taken with the COBE-DIRBE (Cosmic Background Explorer- Diffuse Infrared Background Experiment) satellite. The image shows the boxy shape of the bulge and the dust throughout the disk has a slight reddish color. Also shown are the diffuse thick disk and the spattering of globular clusters and satellites throughout the halo. (Image courtesy of Ned Wright and gsfc.nasa.gov.)

1.3 The Sloan Digital Sky Survey

Large spectroscopic surveys help to map the structure of the Milky Way both spatially and kinematically (Rockosi et al., 2009). Furthermore, measuring chemical abundances in today's stellar populations helps explore the star formation history of the Galaxy (Rockosi et al., 2009). The Sloan Digital Sky Survey (SDSS; York et al., 2000a) is one such survey that has proved itself pivotal in understanding galaxy structure and formation.

The third generation, SDSS-III, is a six year program that includes four surveys-BOSS, MARVELS, SEGUE-2, and APOGEE. The SDSS-III surveys focus on the following three themes: 1) dark energy and cosmology, 2) the structure, dynamics, and chemical evolution of the Galaxy, and 3) the construction of planetary systems (Rockosi et al., 2009). See Sections 1.3.1 below for a more detailed discussion of the SEGUE surveys.

Some of the Sloan Digital Sky Survey's many contributions to date include (SDSS-III, 2008; Rockosi et al., 2009):

- Discovery of distant quasars beyond redshift of 6, revealing supermassive black holes in the early Universe;
- Using weak gravitational lensing to map extended mass distributions around galaxies, demonstrating that dark matter halos extend to 200 kpc or more;
- Demonstration of substructure around the Milky Way, revealing new tidal streams;
- Large samples of white dwarfs, used to obtain an accurate luminosity function to study the cooling of white dwarfs and estimate the age of the Galactic disk;
- Discovery of many new satellites around the Milky Way and M31, nearly doubling

the known number of Milky Way satellites;

• Discovery of stars escaping the Galaxy, revealing encounters with the black hole at the Milky Way's center.

1.3.1 Sloan Extension for Galactic Understanding and Exploration

The Sloan Extension for Galactic Understanding and Exploration (SEGUE and SEGUE-2) surveys provide stellar parameters, kinematics, and metallicities of stellar populations from the disk and inner halo to the large distances of the outer halo (Rockosi et al., 2009). SEGUE-2 added an additional 140,000 stars to the sample size of SEGUE, totaling 380,000 stars, and doubled the number of halo stars observed in SEGUE (Rockosi et al., 2009).

The stars targeted by SEGUE were selected to be largely thick disk stars within 10 kpc of the plane, while SEGUE-2 stars were selected to consist of halo stars with distances greater than 10 kpc (Rockosi et al., 2009). The target selection for SEGUE-2 was predominantly inspired by the desire to study the transition between the inner and outer halos occurring at roughly 15 kpc (Rockosi et al., 2009).

A sample of the science goals of the SEGUE/SEGUE-2 surveys includes: detection and analysis of stellar streams in the inner and outer halos, improved estimates of the mass of the Milky Way, determination of velocities of thick disk, metal weak thick disk, inner halo, and outer halo components, and to provide constraints on models of galaxy formation (Rockosi et al., 2009).

1.4 The Unbound Stellar Population in the Milky Way

Stars can become unbound through a variety of mechanisms. We see this, for example, in galaxy cluster environments where violent interactions are common. Galaxy harassment-minor perturbations over extended timescales (Moore et al., 1996), and tidal stripping from in-falling galaxies (Mihos, 2004; Byrd & Valtonen, 1990) can wrestle stars from the underlying potential and result in a population of unbound stellar outcasts. We see these diffuse collections of stars, called intracluster light (ICL), abandoned to the space between galaxies in large clusters, such as Virgo shown in Figure 1.17 (Mihos et al., 2005). Anywhere from 10% - 70% of the total cluster luminosity may be comprised of ICL (Mihos, 2003; Murante et al., 2004). Determining the fraction of cluster luminosity contributed by the ICL provides insights into the assembly history and evolution of galaxy clusters (e.g.; Napolitano et al., 2003; Feldmeier et al., 2004).

In the Local Group, home to our Milky Way, which is devoid of major galaxy interactions, the unbound stellar population is more likely generated by three-body interactions (Holley-Bockelmann et al., 2005). The preferred method of ejecting stars, called the Hills mechanism (Hills, 1988), involves a binary star system and a supermassive black hole (SMBH). One of the stars is captured by the SMBH, angular momentum is transferred to the companion and it is flung from the galaxy with velocities approaching 1,000 km/s. Figure 1.18 follows the orbital trajectories of a pair of stars as they become disrupted via the Hills mechanism. The ejected companions are called hypervelocity stars (HVS). Semi-analytic models predict O(100) HVS as a result this mechanism (Yu & Tremaine, 2003). Alternatively, this same effect can be achieved with a single star and binary black holes.

A recent simulation study showed that it is possible to boost stars to hyper-velocities via binary disruption scenarios in the disk (Tauris, 2015), Figure 1.19. These disruption scenarios involve a binary star system, located within the Galactic disk, in which the more massive star undergoes a supernova explosion and its companion is flung from the disk. Contrarily, previous works showed that the ejection velocity resulting from a binary disruption mechanism is a relatively modest 300 - 500 km/s generating, instead, high velocity runaway stars (e.g.; Blaauw, 1961; Leonard & Dewey, 1993; Napiwotzki & Silva, 2012). The frequency of bona fide HVS ejected from disk binaries with velocities ~1,000 km/s is yet undetermined⁹.

The following chapters present initial steps towards probing the existence and origin of the unbound stellar population in the Milky Way and how it fits in the larger framework of Galactic structure, formation, and evolution.

⁹We perform a rough estimation of the expected number of low-mass HVS generated from the supernova binary disruption mechanism in Appendix A.



Figure 1.17: Left: Image of the Virgo Cluster captured by the Digital Sky Survey. Right: Long exposure image of the Virgo Cluster from Mihos et al. (2005) revealing the faint intracluster light between the galaxies.



Figure 1.18: Still of an animation depicting the Hills mechanism (Image/animation credit: Andreas Irrgang). Red and blue lines trace the orbits of a binary star system around a SMBH (grey). The red star is captured into a tight orbit around the SMBH and the blue star is ejected from the system.



Figure 1.19: Companion star ejection velocity as a function of supernova-induced kick velocity. Colored lines indicate results of the disk disruption simulations of Tauris (2015). The blue line represents the maximum ejection velocity of $\sim 1,200$ km/s, showing it is possible to eject G/K-dwarf HVSs from the disk with velocities comparable to the G/K-dwarf HVS candidates of Palladino et al. (2014).

Chapter II

IDENTIFYING HIGH METALLICITY M GIANTS AT INTRAGROUP DISTANCES WITH SDSS

Here we reprint, in its entirety, work published in the Astronomical Journal, 2012, Vol. 143, Article ID 128.

2.1 Abstract

Tidal stripping and three-body interactions with the central supermassive black hole may eject stars from the Milky Way. These stars would comprise a set of 'intragroup' stars that trace the past history of interactions in our galactic neighborhood. Using the Sloan Digital Sky Survey DR7, we identify candidate solar metallicity red giant intragroup stars using color cuts that are designed to exclude nearby M and L dwarfs. We present 677 intragroup candidates that are selected between 300 kpc and 2 Mpc, and are either the reddest intragroup candidates (M7-M10) or are L dwarfs at larger distances than previously detected.

2.2 Introduction

A significant fraction of the stellar component of a galaxy cluster is not confined to any galaxy. These stars between galaxies form luminous halos, called intracluster light (ICL), with very low surface brightness that can extend out to several hundred kiloparsecs around individual galaxies (eg., Abadi et al., 2006; Krick & Bernstein, 2007). The brightest ICL is less than 1% of the brightness of the night sky (Mihos, 2003; Feldmeier et al., 2004b), thus making a complete census of ICL very difficult to obtain. High resolution N-body simulations estimate that ICL could comprise 10%-70% of the total cluster luminosity (Mihos, 2003; Murante et al., 2004).

It is commonly thought that intracluster stars are caused by one of three main channels: 1) stripping from galaxies as the cluster assembles either via high speed galaxy encounters, tidal shocking, or a rapidly changing galaxy cluster potential (Byrd & Valtonen, 1990; Merritt, 1984), 2) long-lived, low level cluster perturbations in the form of "galaxy harassment" (Moore et al., 1996), or 3) tidal stripping within in-falling galaxy groups (Mihos, 2004; Rudick et al., 2009). These processes will generate a stellar 'debris field' that is highly inhomogeneous, with distinctly non-Gaussian velocities that reflect an unrelaxed intracluster population (Napolitano et al., 2003). Thus far, ICL has been identified via planetary nebulae (PNe) (Feldmeier et al., 2004b; Aguerri et al., 2005), Red Giant Branch stars (Durrell et al., 2002; Williams et al., 2007), intracluster globular clusters (Lee et al., 2010; Peng et al., 2011), and ultra-deep surface photometry (Feldmeier et al., 2002, 2004a; Mihos et al., 2005).

The properties of ICL may provide insights into the accretion history and evolution of galaxy clusters (Mihos, 2003, 2004; Napolitano et al., 2003; Feldmeier et al., 2004b; Conroy et al., 2007). Although there is some debate about the role that tidal stripping plays in ICL production, it is expected that ICL substructure is correlated with the dynamical state of the cluster (Murante et al., 2007; Mihos, 2004). Since the vast majority of galaxies reside in poor groups, rather than in large clusters, it is of great interest to determine the fraction of unbound stars that reside in these environments.



Figure 2.1: Color-color diagrams of the Covey et al. (2007) sample data. Main sequence stars are marked with black dots, giants with red crosses, and supergiants with green circles. Note that the reddest giants (M7-M10), identified by the black dashed regions, are isolated from the stellar locus. This figure is a theoretical demonstration that very late M giants have the potential to be isolated in color space in the *ugriz* bands, and these dashed regions are meant only to highlight the spectral types of interest and do not represent our color selection criteria (given explicitly in equations 1 and 2). In particular, we drop the u - g and g - r cuts because objects at these distances are likely too faint in the u and g bands. The locus of sub-solar metallicity giants is generally indistinguishable from the dwarf locus, and thus not plotted here.

In the Local Group, ICL has not been observed, though deep observations and star counts have revealed a "field of streams" (e.g., Belokurov et al., 2006b). These streams have been detected out to 100 kpc and are bound to the Milky Way (e.g., Yanny et al., 2000; Ibata et al., 2001). Similarly, faint streams have been detected in the outskirts of M31 (McConnachie et al., 2009; Ibata et al., 2007). Given that the Milky Way and M31 are not yet interacting and may not even be part of the same dark matter halo, it is more likely that Local Group ICL, if it exists, would be a product of a different process altogether.

One of the more recent suggestions for ICL production is via three-body interactions (Holley-Bockelmann et al., 2005). For example, stars can be thrown out from the galaxy through tidal disruption of a binary star system by a supermassive black hole (Hills, 1988; Yu & Tremaine, 2003); this is the most common explanation for 'hypervelocity' stars such as SDSS J090745.0+0204507, with a galactic rest frame velocity of ~ 700 km/sec (Brown et al., 2005). During this process, energy and angular momentum are transferred from the black hole to one of the stars in the binary. The second star loses energy and becomes bound to the black hole while the first is ejected from the galaxy. This is expected to occur at a rate of $10^{-5}(\eta/0.1) \,\mathrm{yr}^{-1}$, where η is the stellar binary fraction (Magorrian & Tremaine, 1999).

Another three-body interaction that is likely to expel stars is a close encounter of a single star with a binary black hole (Yu & Tremaine, 2003). This is expected to occur at a rate of $10^{-4}(\eta/0.1)$ yr⁻¹ (Magorrian & Tremaine, 1999). In this case, the star gains energy from the binary black hole and is flung out of the galaxy while the black hole orbit shrinks (e.g., Quinlan, 1996; Sesana et al., 2006).

To become gravitationally unbound, stars must exceed the escape velocity of the Galaxy, now estimated to be 500-600 km/sec (e.g., Smith et al., 2007). Semi-analytic models predict that there may be approximately 100 hypervelocity stars within 8 kpc of the galactic center if the binary stars have equal masses (Yu & Tremaine, 2003). However, intragroup stars (IGS) may not be solely comprised of hypervelocity stars; they may still be bound but on very large, highly eccentric orbits– this can increase the potential number of IGS. One way to get stars on such eccentric orbits is through three-body galaxy ejections of satellites like Leo I (Sales et al., 2007; Mateo et al., 2008).

As a first attempt to probe for a population of intragroup stars, we develop a technique to search for M giant stars in between the Local Group galaxies. In this paper we present our technique for identifying candidate IGS from the Sloan Digital Sky Survey (SDSS) by applying color, distance, and proper motion cuts. Section 2.2 describes our technique. We present our results in Section 2.3, and we discuss possible sources of contamination in Section 2.4. Section 2.5 concludes and discusses methods to confirm the candidates.

2.3 Methods

During its eight years of operation, the Sloan Digital Sky Survey (SDSS; York et al., 2000a) obtained deep, multi-colored images covering more than a quarter of the sky. The SDSS uses 5 optical bandpasses (u, g, r, i, and z; Fukugita et al., 1996a; Gunn et al., 1998; Hogg et al., 2001; Gunn et al., 2006) with magnitude limits 22.0, 22.2, 22.2, 21.3, and 20.5, respectively. The DR7 data set contains 12,000 square degrees of images and a catalog of over 350 million objects with spectra of 460,000 stars.

As individual red giant stars in M31 have been observed down to the SDSS magnitude



Figure 2.2: Left: Color-color diagram similar to Figure 2.1, including L and T dwarfs from Hawley et al. (2002). The region of color space containing our IGS candidates is marked by the black rectangle. Blue triangles represent dwarfs, green squares represent supergiants, red dots represent giants, blue crosses and green stars represent L and T dwarfs. The blue and green error bars in the bottom left corner are representative of the typical 1- σ error bars for L and T dwarfs, respectively. Note that late M (M7-M9) and L dwarfs also contaminate the space. Right: Same as figure on the left, using SDSS Stripe 82 stars (filled triangles) to represent the spread in the stellar locus (the errors are contained within the size of the point) (Ivezić et al., 2007) and 677 extinction-corrected IGS candidates (filled diamonds). The cyan cloud results from the sum of gaussiandistributed errors on each candidate, i.e. the darkest cyan region represents the most probable location of the data. Similarly as in the left panel, the cyan error bar in the bottom left corner is representative of the typical 1- σ error bars for each candidate.

limits (e.g., Zucker et al., 2007), intragroup red giant stars will be observable. Indeed, at a distance of 300 kpc, all supergiants and roughly half of any giants would be detectable in the r, i, and z bands.

We developed our technique using the synthetic SDSS and 2 Micron All Sky Survey (2MASS) photometry (Covey et al., 2007) of flux calibrated spectra of *solar metallicity* main sequence, giant, and supergiant standard stars (Pickles, 1998). With simulated SDSS and 2MASS colors (Schlegel et al., 1998), we obtained a stellar locus to search for giants or supergiants that are isolated in a 'clean' region of color space. These color-color diagrams are shown in Figure 2.1. We choose solar metallicity standards to probe for IGS generated by three-body interactions within the central regions of the Galaxy, as discussed in Section 2.1.

Unfortunately, most supergiants and giants lie along the main sequence locus, and would therefore be indistinguishable by SDSS colors. However, there is a small area in each of the color-color diagrams, shown with dashed boxes in Figure 2.1, where the rarest M giants (M7-M10) are isolated from dwarfs and supergiants of the same spectral types. Considering the distances we are probing and the very red colors of these spectral types, we restrict our color selection to the r, i, and z bands as follows:

$$2.1 < r - i < 3.4,\tag{2.1}$$

$$1.3 < i - z < 2.2 \tag{2.2}$$

Again, we drop the u - g and g - r cuts because objects at these distances are likely too

faint in these bands.

Since these objects are so red, they may be confused with other nonstellar objects¹. However, we find that even quasars with z > 4.6 are too blue in i - z to fall in our color space (Richards et al., 2002; Fan et al., 2001).

A more worrisome source of contamination comes from L and T dwarfs. To investigate this, we compared our color selection region to the colors of L and T dwarfs (Hawley et al., 2002) and find that they also are contained in the color region, albeit with large uncertainties, as shown in Figure 2.2. From current observational studies, we estimate that there may be O(1000) early L dwarfs in the SDSS footprint within the magnitude limits of SDSS (Burgasser et al., 2010). Since we expect more dwarfs than giants at these faint magnitudes, it is likely that a greater number of dwarfs are scattered into the selection area through large errors than the number of giants scattered out. Objects in our selection box that are not IGS are, nevertheless, likely interesting astrophysical objects yielding more distant L dwarfs than currently known. We discuss ways to differentiate between IGS and contaminants in section 2.5.

From the DR7 Star Table, we identified all stars that satisfy our color criteria and are positioned at |b| > 20 to exclude potential disk contamination (including disk L and T dwarfs). To ensure that each candidate has a stellar point spread function, we confirmed that the object type flag in all 5 bandpasses were stellar². We then removed all objects

¹Background red galaxies have colors roughly 0.35 < r - i < 0.45 and 0.15 < i - z < 0.3 at z=0.1 (Blanton et al., 2003). These colors will become redder with increasing redshift. For example, the brightest and reddest galaxies in SDSS, LRGs at z=0.5, have r - i = 0.7 and i - z = 0.4 (Blanton et al., 2003). Since these colors do not approach our color selection region and since we have ensured stellar point spread functions for our candidates, we do not consider background galaxies as a significant source of contamination.

²Although we restrict our color criteria to the r, i, and z bands, in order to be conservative, we still require the objects to have stellar point spread functions in the u and g bands as well.

that would be nearer than 300 kpc and further than 2 Mpc in the redder bandpasses, assuming an M giant average absolute magnitude in each bandpass ($M_r = 0.8, M_i =$ $-1.5, M_z = -3.5$) yielding the following magnitude cuts: 23.2 < r, 20.9 < i, 18.9 < z— since the r band is effectively below the magnitude limits of SDSS, we searched for null detections in this band as well. We also eliminated objects that triggered any of the following flags: BADSKY, BLENDED, CHILD, COSMIC_RAY, EDGE, MAYBE_CR, MAYBE_EGHOST, MOVED, NODEBLEND, and PSF_FLUX_INTERP.

We cross-checked our candidates with the 2MASS J-H color cut of Majewski et al. (2003) since dwarfs and giants have distinctive J-H colors: $J - K_s > 0.85$, $J - H < 0.561(J - K_s) + 0.36$, $J - H > 0.561(J - K_s) + 0.22$ (top panel of Figure 2.3). Although, note that the Majewski et al. (2003) color cut selects sub-solar metallicity M giants with $[Fe/H] = -0.4 \pm 1.1$ dex (Chou et al., 2007). This will make the giants selected by the Majewski et al. (2003) cut bluer than the solar-metallicity giants selected in our sample.

The middle panel of Figure 2.3 shows a clear separation of the synthetic spectral standard giants and dwarfs from Pickles (1998) at a J-H color around 0.8 (Bessell & Brett, 1988). In general, J, H, and K_s magnitudes for our IGS candidates (16.86, 15.78, and 15.56, respectively, for an M7III according to Covey et al. (2007)) are too faint to appear in the 2MASS catalog. Since appearing in 2MASS would indicate that the source is too bright or too close to be an IGS M giant we removed any candidate that did appear in the 2MASS catalog, and to verify that they are nearby dwarfs, we plotted their J-H colors shown in the bottom panel of Figure 2.3.

The completeness limits of J, H, and K filters in the UKIDSS (Lawrence et al., 2007) survey are appropriate for our targets, however our search through the publicly-released database (DR4) did not result in any matches to our candidates since we expect our candidates lie in parts of the sky not yet covered by this release. Moreover, once the UKIDDS survey is complete it will only cover about half of the area in the SDSS footprint.

After removing known dwarfs in 2MASS, we cross-referenced the IGS candidates with the USNO-B catalog, an all-sky catalog that contains positions, proper motions, and magnitudes in various optical passbands. Here, we removed those candidates with discernible proper motions, ranging from 3 to 1412 mas/year, since the distances determined by the proper motions indicate that they are likely objects closer than 353 pc. This eliminated 782 candidates.

2.3.1 Testing the color selection region with real data

Since our current technique for finding very late M giant IGS is based on idealized solar metallicity spectra and synthetic colors, it is important to determine how robust these colors are for known M7 III-M10 III stars. Unfortunately, there are no confirmed very late M giant stars in the SDSS DR7 (or DR8) database with both spectra and colors. While it is true that these stars are rare, the real difficulty is that spectral classification of M giants is notoriously difficult, and the latest M giants can be spectral type variables, as well. Since the SDSS database did not explicitly include the latest M giants, we decided to take a three-pronged approach in checking M giant colors.

First, we searched through Simbad for spectroscopically-confirmed M7-M10 giants, finding 53, 28, 4, and 0, respectively; many of these were listed in the Catalogue of Stellar Spectral Classifications (Skiff, 2010). For each object, we cross-checked the spectral type through all other publically-available catalogs on VizieR to determine if the star was a known significant spectral variable, and if so, we discarded it. We then searched DR7 for any star within 2 arcseconds of the target and obtained the photometry. Many of these objects were nearby and saturated, and therefore appeared as several non-stellar sources in the DR7 database – these were also discarded if the composite, corrected photometry was not available. Only one of the remaining stars returned a result in the CAS and appears as the blue star in Figure 2.4, which lies within our color cut.

Second, we acquired spectra of late-type giants observed as part of a study to quantify the effects of gravity on the spectra of cool objects. These spectra were obtained with the Low Resolution Imaging Spectrometer (LRIS, Oke et al., 1995) on the 10-m W. M. Keck Observatory as part of a campaign to construct a systematic surface gravity "grid" to further constrain spectral classifications of brown dwarfs (Kirkpatrick, in prep.). We estimated the spectral type and log g parameters by eye and separated the latest M giant subset for analysis, totaling 13 spectra. Examples of these spectra are shown in Figure 2.5. Using the SDSS transmission curves, we calculated the colors in the Sloan bands, and as can be seen by the red stars in Figure 2.4, most of these stars do indeed lie in the predicted color space.

Finally, as a further check of our M giant colors we used the Bruzual-Persson-Gunn-Stryker (BPGS, hereafter; Laidler, V. et al, Synphot User's Guide, Version 5.0 (Baltimore STScI), 2005) stellar atlas of standard stars to obtain synthetic colors using the IRAF Synphot task *calcphot* ³; the green points in Figure 2.4 represent these standard starsthere were no available M9III-M10III standards in the atlas, which is expected because

 $^{^{3}}$ The i-z colors returned by *calcphot* for the M dwarfs are unreliable due to negative flux values in the BPGS stellar atlas.

these objects are all spectrum variables.

Figure 2.4 shows that the color cut we defined from synthetic SDSS colors is consistent with the colors of known M giants from all three techniques.

If the reader is interested in seeing the colors of the stars we used from the BPGS stellar atlas in the Johnson-Cousins filter set, a figure will be available online ⁴. The UBV colors of these stars are consistent with the colors reported by Worthey & Lee (2006), as well.

2.4 Results

We found a small region of color space, shown in Figure 2.2, in which the reddest solar-metallicity M giants are isolated from the rest of the stellar locus. This region hosts M7III-M10III stars, along with L dwarfs (Hawley et al., 2002).

Using our color selection criteria, we found 159,108 extinction-corrected objects in SDSS DR7. After applying the distance cut and checking the data flags, we narrowed the sample to 4181 objects. We then cross-correlated our sample with the 2MASS and USNO-B surveys, removing any stars with dwarf-like J-H colors and any stars with non-zero proper motions. Our final sample contains 677 IGS candidates. Table II.1 lists positions, asinh magnitudes, r - i, and i - z colors for all 677 candidates. The right panel of Figure 2.2 shows the location of the final set of IGS candidates with errors in r - i/i - z color space.

⁴http://astro.phy.vanderbilt.edu/~palladl2/

2.5 Discussion

As discussed in Section 2.1, IGS could be formed from several different methods. Considering the Local Group's current level of interactions, this population may likely be comprised of high metallicity hypervelocity stars (HVS) ejected through the three body mechanism. However, not enough is known about HVS and Local Group formation to say this definitively, so probing the IGS sample may help us to constrain either or both of these.

If every candidate were a solar metallicity IGS giant, they would be rare tracers of a large underlying IGS population. Assuming a single burst of star formation 10 Gyrs ago and a Salpeter initial mass function (IMF), these candidates represent $O(10^{-4})$ of the total number of IGS and $O(10^{-3})$ of the total mass in IGS, and varies only slightly for differing choices of IMF.

It is useful to compare this to theoretical predictions of stellar ejections from the Milky Way (Kollmeier et al., 2009). Stars ejected from the galaxy center through three-body interactions with a SMBH will typically have much higher metallicity than stars that were stripped from satellite galaxies originating in the outskirts of a galaxy halo (e.g., Jacoby & Ciardullo, 1999a; Kirby et al., 2008). For example, if we assume that all of the IGS are solar metallicity hypervelocity stars ejected by three-body interactions with a binary black hole consisting of a SMBH and an intermediate mass black hole (IMBH), then the total mass in stellar ejecta will be roughly equal to the mass of the IMBH (Yu & Tremaine, 2003; Quinlan, 1996). Given the Milky Way SMBH mass of $4 \times 10^6 M_{\odot}$ (Ghez et al., 2008), we would require an IMBH with mass roughly $10^5 M_{\odot}$ as the companion, independent of initial mass function. This IMBH mass is similar to the mass of the IMBH proposed to be responsible for ejecting stars in the central region of the Galaxy (Lang et al., 2011). This yields several IGS per square degree of sky and roughly tens of red giant hypervelocity IGS in the SDSS footprint⁵.

Similar back of the envelope calculations suggest that there are O(1000) L dwarfs located in the SDSS footprint, and realizing that late-type dwarfs are more common in general, we anticipate that the majority of our IGS candidates are likely L dwarfs. If these IGS candidates do turn out to be L dwarfs, then we have identified L-type dwarfs at distances of 100-200 pc, which is up to 4 times farther than currently known (Schmidt et al., 2010).

In an attempt to determine if the IGS sample has a distinct spatial distribution, we conducted a set of 2-dimensional K-S tests that compared our candidates with template samples drawn from 3 characteristic distributions: 1) an exponential disk with a scale height of 300 pc to mimic an old stellar population, and a distance cut off of 200 parsec to resemble an L dwarf distribution; 2) a random distribution; 3) and a set of observed hyper-velocity stars (Brown et al., 2009). We convolved each template data set with the SDSS footprint and employed the same galactic latitude cut as in our IGS sample. Assuming that these are very cool dwarfs, the IGS sample should exhibit the same distribution on the sky as the exponential disk, but the 2-d K-S test revealed otherwise: the probability that these two samples come from the same underlying distribution is only 10^{-4} . This is ultimately not surprising since we removed any objects with a measurable proper motion, strongly selecting against L dwarfs within the disk. The second test with a random

⁵We expand on the expected number of total IGS and red giant IGS in Appendix A.

distribution resulted in an even smaller probability of 10^{-5} , while the third test with the hypervelocity sample yielded a somewhat higher probability of 10^{-2} . Figure 2.6 shows the position of our IGS candidates on the sky compared to the hypervelocity sample used here.

2.6 Summary and Conclusions

We identified 677 intragroup stellar candidates from the SDSS DR7 using color cuts based on solar metallicity spectral standards. These are extremely red stars with 1.3 < i - z < 2.2, though the color would shift bluer with lower metallicity. As shown in Figure 2.2, the M giants in the region are not completely isolated. The latest M dwarfs and early L dwarfs are possible sources of contamination.

Followup photometric observations of our candidates in the near to mid-infrared wavelengths may differentiate between late dwarfs and M giants. Future followup photometric observations with a 4m class telescope may be promising, albeit impractical. For example, the FLAMINGOS instrument on the NOAO 4m telescope could image all of our candidates with a 113 hour total exposure time in each of the J and H bands, and over 600 hours of total exposure time in the K band for a 10-sigma detection, while this likely would not be sufficient to distinguish M giants from dwarfs. Similarly, the J, H, and K magnitude limits of NIRI on Gemini are appropriate for our targets, although would require a prohibitively long total exposure time of 2031 hours to achieve a S/N of about 12.

Also, it is possible with long-term photometric followup observations on a 1m class telescope to differentiate between dwarfs and giants based on variability, as late-type M

Table II.1. IGS candidates remaining after all criteria cuts [Complete version provided in Appendix C].

Object ID	RA	DEC	u	u err	g	g err	r	r err	i	i err	\mathbf{z}	z err	r-i	i-z
758882136836343139	60.5013	80.8760	24.4	1.3	24.7	0.6	24.3	0.7	22.0	0.2	20.3	0.2	2.3	1.7
758882137910740030	62.9212	82.5298	25.2	1.2	25.0	0.6	23.2	0.4	21.0	0.1	19.6	0.1	2.2	1.4
758882626993718584	63.3170	81.6950	26.5	0.6	24.8	0.6	25.3	0.7	22.0	0.2	20.3	0.2	3.3	1.6
758877527803168329	94.6323	63.7038	25.1	0.9	24.1	0.3	24.5	0.5	21.2	0.1	19.8	0.1	3.3	1.4
758877527266231607	94.8017	63.2633	24.6	0.7	24.4	0.5	24.5	0.6	21.8	0.1	20.4	0.1	2.7	1.4
758877527266493955	95.9369	63.5501	25.2	0.7	24.8	0.6	25.6	0.5	22.7	0.3	20.7	0.2	2.9	2.0
758878272976455144	105.3459	66.8553	25.3	0.9	24.4	0.4	24.1	0.4	21.8	0.1	20.2	0.1	2.3	1.6
758878271902778853	105.8427	66.0785	23.8	0.7	24.7	0.4	24.6	0.5	22.3	0.2	20.8	0.2	2.3	1.5
758884768580109918	107.9936	38.3022	24.8	1.1	25.1	0.7	24.4	0.6	21.9	0.1	20.4	0.2	2.6	1.4
587738067260998978	109.5721	39.4395	25.3	0.9	25.7	0.6	24.6	0.7	22.1	0.2	20.0	0.1	2.5	2.1

giants tend to be highly variable.

Naturally, low resolution spectroscopic follow-up observations of these IGS candidates would be ideal to confirm their luminosity class. The Calcium II Triplet (CaT) feature at 8498Å, 8542Å, and 8662Å is particularly useful for distinguishing late M dwarfs from giants, being much more prominent in the spectra of late-type dwarfs (Reid & Hawley, 2005). In addition, the strength of the Calcium Hydride (CaH) feature between 6800Å and 7000Å is a good indicator of luminosity class (Cohen, 1978).

Once the confirmation is complete, we can test the efficiency of our color selection technique, which will be useful for large data surveys like LSST. In fact, two surveys set to launch in the coming year will be particularly well-tuned to find IGS M giants. The DECam survey on the CTIO 4-meter telescope will observe over 1000 square degrees, with magnitude limits of r=23.4, i=24.0, and z=22.9 – over two magnitudes deeper than SDSS in z. An even deeper survey will launch on the Subaru telescope; the HyperSuprimeCam plans to observe 2000 square degrees down to z=24.9 and y=23.7. Eventually, deeper observations of IGS can reveal the metallicity of these stars – an important clue to their original birthplace within the group or in situ in the intergalactic medium.



Figure 2.3: Top: Blue dots are SDSS Stripe 82 stars (Ivezić et al., 2007) that satisfy the Majewski et al. (2003) M giant color cuts. Larger red dots are M giants from Covey et al. (2007). The point located outside the color cut corresponds to an M10 spectral type. Also note that the late M giants (M7-M10) are located in the part of this color region that is least populated, so they will be least likely to be identified by this cut. Middle: Red open circles are M giants from Covey et al. (2007). Blue dots are M dwarfs from Covey et al. (2007). Bottom: IGS candidates with 2MASS JHK colors. The solid line represents the J-H = 0.8 separation between dwarfs and giants.



Figure 2.4: Same as the left panel of Figure 2.2. The green stars overplotted here represent M-giant standards, with spectral types between M6III and M8III, from the BPGS stellar atlas. The larger green dots are dwarfs with spectral types O through M, from the same atlas. The photometry for these stars was obtained by implementing IRAF *symphot* tasks. The blue star represents the M giant identified from the Catalogue of Stellar Spectral Classifications (Skiff, 2010) with confirmed SDSS photometry. The red stars are the M giant contaminants in the Kirkpatrick (in prep.) data for which we received spectra.



Figure 2.5: Four of the spectra used to compare colors to our IGS candidates. The resulting colors are shown in Figure 2.4.



Figure 2.6: The relative positions of our IGS candidates and the hypervelocity stars of Brown et al. (2009) that were compared with the 2-d K-S test. The comparison was made between 522 IGS candidates, red dots, and 22 HVS, blue stars. Notice the higher density of IGS candidates at the edges of the footprint signifying larger numbers of these stars at lower Galactic latitudes.

Chapter III

HYPERVELOCITY STAR CANDIDATES IN THE SEGUE G & K DWARF SAMPLE

Here we reprint, in its entirety, work published in the Astrophysical Journal, 2014, Vol. 780, Article ID 7.

3.1 Abstract

We present 20 candidate hypervelocity stars from the Sloan Extension for Galactic Understanding and Exploration (SEGUE) G and K dwarf samples. Previous searches for hypervelocity stars have only focused on large radial velocities; in this study we also use proper motions to select the candidates. We determine the hypervelocity likelihood of each candidate via Monte Carlo simulations, considering the significant errors often associated with high proper motion stars. We find that nearly half of the candidates exceed their escape velocities with at least 98% probability. Every candidate also has less than a 25% chance of being a high velocity fluke within the SEGUE sample. Based on orbits calculated using the observed 6-d positions and velocities, few, if any, of these candidates originate from the Galactic center. If these candidates are truly hypervelocity stars, they were not ejected by interactions with the Milky Way's supermassive black hole. This calls for a more serious examination of alternative hypervelocity star ejection scenarios.

3.2 Introduction

Hypervelocity stars (HVSs) are believed to be ejected by three-body interactions with the supermassive black hole (SMBH) at the Galactic center (e.g., Hills, 1988; Yu & Tremaine, 2003; Brown et al., 2005). During this process, energy and angular momentum are transferred from the black hole to one of the stars in a binary system. The second star loses energy and becomes bound to the black hole while the first is ejected from the Galaxy. In this scenario HVSs can probe conditions in the Galactic center such as the binary fraction, and even place limits on the existence of a second, tightly bound SMBH. Semianalytical models predict that there may be approximately 100 HVSs within 8 kpc of the Galactic center due to the break up of equal-mass binaries (Gould & Quillen, 2003; Yu & Tremaine, 2003).

While the SMBH at the Galactic center remains the most promising culprit in generating HVSs, other hypervelocity ejection scenarios are possible, such as a close encounter of a single star with a binary black hole (Yu & Tremaine, 2003). In this case, the star gains energy from the binary black hole and is flung out of the Galaxy while the orbit of the black hole binary shrinks (e.g., Quinlan, 1996; Sesana et al., 2006). Another alternative hypervelocity ejection model involves the disruption of a stellar binary in the Galactic disk; here a supernova explosion in the more massive component can accelerate the companion to hypervelocities (e.g., Blaauw, 1961; Leonard & Dewey, 1993; Napiwotzki & Silva, 2012).

At least 18 HVSs have been discovered in the Milky Way within the past decade with velocities as high as 700 km s⁻¹ (e.g., Brown et al., 2005; Edelmann et al., 2005; Hirsch

et al., 2005; Brown et al., 2009, 2012). So far, all confirmed HVSs are massive B-type stars such as those observed around the central SMBH (e.g., Brown et al., 2009, 2012). However, since the ejection mechanisms described above apply to any stellar mass, it is important to search for HVSs among the larger set of longer-lived, lower mass stars (e.g., Quinlan, 1996). If a SMBH ejection mechanism is at play, then metal-rich stars originating from the Galactic center ought to pollute the metal-poor halo. Previous attempts to mine Sloan Extension for Galactic Understanding and Exploration (SEGUE) and SEGUE-2 stellar halo data found no metal-rich, old ejected stars (Kollmeier et al., 2009, 2010). The lack of a significant population of old, metal-rich HVSs suggests that the initial mass function at the Galactic center is mildly top-heavy. Alternatively, hypervelocity ejection mechanisms may be more complex than previously thought.

In this paper we identify the first set of G- and K-type candidate HVSs from SEGUE. We discuss candidate selection in Section 3.2, including a description of the G and K dwarf sample, and we address the significant proper-motion errors in Section 3.3. Section 3.4 contains orbital parameters for the HVS candidates, and Section 3.5 discusses possible alternative origin scenarios. Finally, we summarize and conclude in Section 3.6.

3.3 Identifying HVS Candidates

Our candidates are drawn from the G and K dwarf stars in SEGUE (Yanny et al., 2009) from the Sloan Digital Sky Survey (SDSS) Data Release 9 (DR9; SDSS-III Collaboration et al., 2012). As part of SDSS (York et al., 2000b), SEGUE provides medium-resolution $(R \approx 1800)$ spectroscopy over a broad spectral range (3800–9200 Å). Probing more than 150 lines of sight, SEGUE covers $\approx 3500 \text{ deg}^2$ of the sky, with spectroscopy of $\approx 240,000$



Figure 3.1: Transverse versus radial velocities of our HVS candidates, in kilometers per second. Red lines indicate a transverse velocity $\sqrt{2}$ times higher than the radial velocity, as expected for an isotropic stellar distribution. Blue lines represent a transverse velocity 5 times higher than the radial velocity. The majority of our candidates show large transverse-to-radial velocity ratios, characteristic of a sample strongly affected by large proper-motion errors. We caution that some of our HVS candidates may be high-velocity flukes, and we calculate the likelihood of this in Section 3.3.2.

stars over a range of spectral types. Technical information about SDSS has been published on the survey design (York et al., 2000b; Eisenstein et al., 2011), telescope and camera (Gunn et al., 1998, 2006), and spectrographs (Smee et al., 2013), as well as the photometric system (Fukugita et al., 1996b) and astrometric (Pier et al., 2003) and photometric (Ivezić et al., 2004) accuracy.

G and K dwarfs are selected from the SDSS photometric data using simple color and magnitude selection criteria. The 42,901 SEGUE G dwarfs are defined as having 14.0 $< r_0 < 20.2$ and $0.48 < (g - r)_0 < 0.55$, while the 28,332 K dwarfs have $14.5 < r_0 < 19.0$ with $0.55 < (g - r)_0 < 0.75$ (Yanny et al., 2009). The subscript zero indicates that the color and magnitude have been corrected for dust extinction, using estimates derived from Schlegel et al. (1998). Each spectrum is analyzed with the DR9 SEGUE Stellar Parameter Pipeline (SSPP), which provides estimates of effective temperature, surface gravity (log g), [Fe/H], and [α /Fe] (Lee et al., 2008a,b; Allende Prieto et al., 2008; Smolinski et al., 2011; Lee et al., 2011). We follow the quality protocol of Schlesinger et al. (2012) to remove targets with poor signal-to-noise ratio (S/N <10), incalculable atmospheric parameters, excessive reddening (greater than 0.5 mag in r), saturated photometry ($r_0 < 15$), or flags indicating temperature or noise issues. We also use the SSPP log g estimates to ensure the stars are dwarfs, using a cut on log g as a function of [Fe/H] to isolate dwarf stars (K. J. Schlesinger et al. 2014, in preparation).

For each star that satisfies these criteria, we determine its distance using the isochronematching technique described by Schlesinger et al. (2012). Briefly, each star is matched in [Fe/H] and $(g - r)_0$ to 10 Gyr isochrones from the empirically corrected Yale Rotating Stellar Evolution Code set (An et al., 2009). There are systematic distance uncertainties introduced by using 10 Gyr isochrones, as well as the possibility of undetected binarity; this leads to a systematic shift in distance of -3% for the most metal-rich stars, while metal-poor stars are largely unaffected; this is factored into our distance estimates. There are also random distance errors from uncertainties in photometry, [Fe/H], [α /Fe], and, finally, isochrone choice. The total random distance uncertainty is dominated by uncertainties in [Fe/H] and ranges from around 18% for stars with [Fe/H] > -0.5% to 8% for more metal-poor stars.

To identify HVSs, we convert the radial and tangential velocities to Galactic Cartesian coordinates, as described below in Section 3.4.1, and choose a simple but conservative initial total velocity threshold of 600 km s⁻¹ to identify stars that exceed the Galaxy's escape velocity (e.g.; Smith et al., 2007). We then verify that each candidate exceeds the escape velocity at its current location within the Galaxy. This procedure yields 42 preliminary HVS candidates of varying quality, which we further glean as described in Section 3.3.1 below.

3.4 Estimating the Fidelity of Our Candidates

3.4.1 Proper-Motion Quality Cuts

The proper-motion distribution in SDSS is skewed toward large proper-motion errors. We must ensure that the extreme velocities of our candidates are real rather than the product of large errors. We describe our technique to ensure the robustness of our candidates in this section. The first step in defining a clean hypervelocity sample is to assess the quality of the proper-motion measurement. To determine the proper motions for SDSS targets, Munn et al. (2004, 2008) matched each SDSS point source to the USNO-B catalog. The resulting SDSS+USNO-B catalog is 90% complete to g < 19.7 and has statistical errors of approximately 3–3.5 mas yr⁻¹ and systematic errors of ≈ 0.1 mas yr⁻¹ for each component.

Munn et al. (2004) defined a number of criteria to ensure that the SDSS+USNO-B proper motions are reliable; these conditions resulted in a version of the USNO-B catalog with a contamination of less than 0.5%. The criteria were later revised by Kilic et al. (2006) and are as follows¹:

- The number of objects in USNO-B within a 1" radius of the SDSS target should be 1 (match=1).
- The rms residual for the proper-motion fit in right ascension and declination must be less than 525 mas (sigRA < 525 and sigDEC < 525).
- There must be at least six detections (including the SDSS observations) used to determine the proper motion (nFit = 6).
- The distance to the nearest neighbor with g < 22 must be greater than 7" (dist22 > 7).

Only three of our 42 preliminary candidates met all of the proper-motion quality criteria, and we categorize these as "Clean." We performed an in-depth analysis for the remaining 39 stars, assigning each a likelihood of proper-motion contamination based on the same criteria adopted by Kilic et al. (2006) for their white dwarf sample. They found

¹The parameters listed in parentheses are available in the DR9 proper motions catalog in the SDSS Catalog Archive Server (CAS).
that the chance of contamination for a target with

- Six detections and a neighbor within 7" is less than 1.5%;
- Five detections and no near neighbors is 1.5%;
- Five detections and a neighbor within 7" is 35%;
- Four detections and no neighbor within 7" is 51%; and
- Three detections and no neighbor within 7" is 89%.

We further checked for any potential blending issues by visually inspecting each candidate.

We categorize 17 stars as having "Reliable" proper motions, with 1.5% or less chance of contamination and no visual blending. We categorize 10 stars as "Possible," meaning they have between 35% and 51% chance of contamination. Twelve stars were removed because of visual blending.

We choose to consider only those candidates with "Clean" and "Reliable" designations. Thus, our final sample contains 20 HVS candidates, all with greater than 98.5% probability of robust proper-motion estimates.

We do expect that this final sample contains false-positive HVS detections. One way to illustrate this is with Figure 3.1, which compares the transverse and radial velocities of our HVS sample. For a random isotropic stellar distribution, the transverse velocity should be roughly $\sqrt{2}$ times higher than the radial velocity, and the fact that this sample is predominantly composed of stars with much larger transverse-to-radial velocity ratios is a classic signature of contamination by large proper-motion errors. While this does not prove that all of the candidates are spurious, it does indicate that many of them may be. Of course, a true hypervelocity sample would not be well represented by a random, isotropic distribution, but we caution that it is premature to say that we have identified 20 HVSs. We conducted further statistical tests, described below, to evaluate the likelihood that each HVS candidate is real.

3.4.2 Monte Carlo Sampling

Although the typical error in proper motion in the SEGUE database is ~ 10 mas yr^{-1} , proper-motion errors can, for some stars, be much larger than expected for a normal distribution, especially at the high-velocity end (Gould, 2003; Gould & Salim, 2003; Gould & Kollmeier, 2004; Munn et al., 2004, 2008; Bond et al., 2010). With this in mind, we consider the possibility that these HVS candidates may have true velocities much lower than can be explained by the reported errors and that they are in fact bound to the Galaxy.

In order to determine the true range of velocities for our HVS candidates, as well as the probability that these candidates are bound given a more realistic error distribution, we built a Monte Carlo simulation to sample possible orbital parameters for each HVS candidate. Dong et al. (2011) obtained a proper-motion error distribution for the SDSS+USNO-B catalog by compiling proper motions for a sample of SDSS quasars that met the Kilic et al. (2006) criteria. We randomly resampled a million realizations of each HVS candidate's kinematics from the Dong et al. (2011) non-Gaussian proper-motion error distribution and Gaussian radial velocity errors. We also resampled each candidate's position, assuming Gaussian errors in the distance determinations as well. We find that 13 of the 20 candidates remain hypervelocity with greater than 90% probability. Figure 3.2 shows the distribution of velocities drawn randomly from the errors for the three least and most bound candidates. In most cases, the drawn velocity well exceeds the escape velocity, represented by the vertical dashed lines.

We performed a second Monte Carlo test to quantify the chance that these high velocities are simply the extreme tail end of the velocity error distribution within the entire SEGUE G and K dwarf sample. Here we construct a new mega-SEGUE sample built from 1000 realizations of each SEGUE star, in which each realization is drawn from the error distribution in proper motion, radial velocity, and distance as described above. We then calculate the "interloper likelihood" for each candidate with respect to the mega-SEGUE sample; this is the probability that a slow, noncandidate star within our sample could have had the observed velocity of a particular candidate, given the errors:

$$P(\text{interloper}, i) = 1 - \frac{n_{\text{HVS}}(v \ge v_{\text{cand},i})}{n_{\text{tot}}(v \ge v_{\text{cand},i})},$$
(3.1)

where $n_{HVS}(v \ge v_{\text{cand},i})$ is the number of stars in the mega-SEGUE sample with velocity greater than or equal to the observed velocity of candidate *i* that were originally tagged as hypervelocity in the data and $n_{tot}(v \ge v_{\text{cand},i})$ is the total number of stars in the Monte Carlo sample with velocity greater than or equal to the candidate's velocity. All candidates have less than a 25% "interloper likelihood," and more than half have less than 10%. Together, these two tests indicate that some of our candidates may in fact be the result of a statistical fluke. However, we expect the bulk of the candidates to remain hypervelocity.

Table III.1 lists the velocity of each candidate determined from the proper motions

reported in DR9, the minimum velocity calculated from a million realizations of the proper motion, radial velocity, and distance errors, the escape velocity for each candidate in a spherically symmetric Galaxy, the probability that the candidate may be bound given the escape velocity at its position, and the interloper likelihood as described above.

3.5 Orbits of HVS Candidates

3.5.1 Galaxy Model

We construct an analytical, multicomponent model of the Milky Way gravitational potential to predict the orbits of stars in the Galaxy based on the initial six-dimensional observed position and velocity. The model is easily modifiable and can be tuned to reflect the observed Galactic structural parameters.

Our model includes the following components: a central SMBH with $M_{\rm SMBH} = 4 \times 10^6 M_{\odot}$; a spherical Hernquist bulge (Hernquist, 1990) with $M_{\rm bul} = 4.5 \times 10^9 M_{\odot}$ and $r_{\rm bul} = 2.5$ kpc; Miyamoto–Nagai thin and thick disks (Miyamoto & Nagai, 1975) with $M_{\rm thin} = 6 \times 10^{10} M_{\odot}, M_{\rm thick} = 6 \times 10^9 M_{\odot}, 0.3$ kpc thin-disk scale height, 1 kpc thick-disk scale height, and 3 kpc scale lengths for both; and a Navarro–Frenk–White (NFW) dark matter halo (Navarro et al., 1997) following the formalism of Lokas & Mamon (2001); we chose $M_{\rm NFW} = 10^{12} M_{\odot}, R_{\rm vir} = 200$ kpc, and c = 10 for the Milky Way. Recent studies have argued for a shorter thick-disk scale length (e.g.; Cheng et al., 2012; Bovy et al., 2012b; Bensby et al., 2011); however, this change would have a negligible effect on our results because of the comparatively low mass of the thick disk component.

The model can also be tuned for varying degrees of axisymmetry or triaxiality in the

halo. For this study we use both spherical and triaxial models. For the triaxial model, we adopt the axis ratios b/a = 0.99 and c/a = 0.72 (Law & Majewski, 2010).

The Galactic Cartesian coordinate system used here is centered on the Galactic center; the x-axis points from the center toward the Sun (located at x = 8.2 kpc (Schönrich, 2012), the y-axis points along the direction of Galactic rotation, and the z-axis points toward the North Galactic Pole. To calculate velocity in this coordinate system, we convert radial velocity, distance, and proper motions to U, V, and W in the Galactic coordinate system. Note that issues with astrometry in DR8, as explained in Section 3.3.5 of Aihara et al. (2011b) and the associated erratum (Aihara et al., 2011a), have been resolved for the DR9 astrometry used here. We choose the velocity of the local standard of rest to be 238 km s⁻¹, and the motion of the Sun with respect to that is $U = -13.8 \text{ km s}^{-1}$, $V = 12.24 \text{ km s}^{-1}$, and $W = 7.25 \text{ km s}^{-1}$ (Schönrich, 2012; Schönrich et al., 2010). This Galactic model is consistent with the measured proper motion of $SgrA^*$, $6.379\pm0.026~{\rm mas}~{\rm yr}^{-1}$ (Reid & Brunthaler, 2004). Then, $U,\,V,$ and W are transformed into the Galactic Cartesian coordinate system, and we calculate the orbits backward in time for 1 Gyr using a fourth-order Runge-Kutta integrator. The choice of 1 Gyr is sufficient to discern the direction of origin while not being significantly influenced by a changing Galactic potential.

We examine the variation of each candidate's orbit given the errors described in Section 3.3.2. Figure 3.3 shows the 1σ and 2σ orbits for HVS 20, indicating that for some of the candidates the velocity errors are sufficiently large that the candidate itself may be bound. We find that the differences between orbits in the spherical versus the triaxial model are negligible for unbound orbits, since the stars have little time to respond to the halo

potential. Therefore, for simplicity, when discussing unbound orbits we show only those in the spherical case. However, in instances when the orbit may be bound, as for HVS 20 in Figure 3.3, the halo shape definitely influences the candidate's trajectory, suggesting that marginally bound stars may help constrain halo triaxiality.

3.5.2 Origins

As shown in Figure 3.4, the trajectories of these HVS candidates do not originate from the Galactic center, which would be expected if the stars were ejected by three-body interactions with the SMBH. Instead, they appear to be coming from all directions, which suggests that other ejection processes may be at play.

We considered the SMBH at the center of M31 as a possible source (Sherwin et al., 2008), and given the velocities of the stars, we find the required flight time to reach the solar neighborhood would be approximately 1 Gyr. Figure 3.4 shows the orbits corresponding to the seven most unbound candidates, each with a 1σ "wedge" of possible orbits. It can be seen that the candidates could not have come from M31's SMBH position 1 Gyr ago (dashed lines). The orbits of the other 13 candidates are consistent with not arriving from M31. Therefore, the SMBH at the center of M31 is not responsible for ejecting these stars. However, this does not exclude other Galactic and extragalactic sources such as globular clusters, satellite galaxies, or the centers of distant galaxies within ~ 10 Mpc.

3.6 Discussion

3.6.1 Chemical Tracing

Since it is more difficult to trace the past orbits of globular clusters and known satellite galaxies because of tidal stripping, shocks, and other mass-loss effects, we cannot say with certainty whether these stars could have originated in the Galactic disk, the bulge, or globular clusters. Another approach to determine whether these candidates belong to a particular population is to examine their chemical compositions.

We compared the metallicities of our candidates with the metallicity distribution functions (MDFs) of known globular clusters (Harris, 1997), the SEGUE G and K dwarf samples representative of the Milky Way disk population (Schlesinger et al., 2012), the Galactic bulge (Sadler et al., 1996), and by extension the bulge of M31, assuming a peak metallicity of +0.23 (Jacoby & Ciardullo, 1999b). We also compared the metallicity distribution of our candidates with the MDF of the Galactic halo, although with a peak at [Fe/H] < -1 (An et al., 2013) it is clearly inconsistent with our candidates.

The MDFs for each population are shown in Figure 3.5. As perhaps expected, the metallicity of the HVS candidates is consistent with the G and K dwarf samples in the disk. Their metallicities are also largely consistent with the high-metallicity end of the globular cluster population and the low-metallicity end of the Galactic (and M31) bulge. Similarly, the stars' $[\alpha/Fe]$ -values are broadly consistent with these stellar populations. Unfortunately, based on the information here, none of these populations can be decisively ruled out as a possible source, although it is clear that these HVSs do not originate from the metal-poor globular cluster system.



Figure 3.2: Velocity distribution for a million random samples of the velocity error distribution for the three least and most bound HVS candidates. Dashed lines show escape velocity of each candidate.



Figure 3.3: Orbit of HVS 20, a candidate with a "Reliable" proper-motion measurement but the largest probability of being bound, shown by the black lines. Also shown are the resulting orbits for the same candidate with 1σ (red lines) and 2σ (blue lines) velocity errors from the million Monte Carlo realizations. Left, two-dimensional projections in the spherical dark matter halo; right, the same for the triaxial model. The black dots and plus signs represent the locations of the Galactic center and the Sun, respectively, while the pale blue ellipses provide a rough scale for the extent of the disk. The five-pointed star in each panel marks the current position of HVS 20. The top row is a top-down view of the Galaxy while the middle and bottom rows are side views along the disk. Here we show that some candidates may in fact live on very bound orbits, and in such cases the shape of the orbit is strongly influenced by the triaxiality of the halo.



Figure 3.4: Orbits of the seven HVS candidates that are unbound with at least 98% proability, over the past 1 Gyr. As in Figure 3.3, the black dots and plus signs represent the locations of the Galactic center and the Sun, respectively, while the pale blue ellipses provide a rough scale for the extent of the disk. The shaded regions flanking the orbits of the same color represent the "wedges" of possible orbits given the 1σ velocity errors for the corresponding candidate. The like-colored stars mark the current positions of the candidates. None of the orbits plotted here intersect near the Galactic center, suggesting a different origin for these stars. In additional, if these stars had been traveling for 13 Gyr, they may have originated from as far as tens of megaparsecs away. The dashed lines, highlighted in yellow for visibility, point toward M31's location 1 Gyr ago, assuming M31 proper motion, radial velocity, and distance from Sohn et al. (2012). This interval was chosen to roughly coincide with the travel time of a hypervelocity star originating in M31. None of these HVS candidates seem to be coming from M31, which therefore is ruled out as a possible origin.

		р						%	"Interloper	
Fe/H [α/Fc	e] 9	(kpc)	$v_r^{\rm b}$	$v_t^{\ c}$	v^{q}	v_{\min}^{e}	$v_{\rm esc,Sph}^{\rm f}$	Bound	Likelihood"	Rating
-0.06 0.10	0	3.70	-76.0	56.1	802.2	92.2	533.6	6.35	0.02	Clean
-0.15 0.0	6	5.68	-25.6	15.7	628.6	290.0	517.3	7.43	0.18	Clean
-0.91 0.4	0	4.06	31.7	23.7	641.8	195.1	588.9	34.88	0.15	Clean
-0.83 0.4	6	1.04	-54.9	237.7	1296.7	587.4	552.3	0.0	0.00	Reliable
-0.34 0.19	6	3.19	58.0	61.5	1086.8	378.9	576.5	0.04	0.00	$\mathbf{Reliable}$
-0.55 0.26		3.13	80.2	46.2	916.3	286.6	534.5	0.07	0.01	Reliable
-0.55 0.35		3.06	20.4	38.1	793.9	285.0	530.2	0.30	0.02	Reliable
-0.65 0.26		2.48	6.26	51.8	769.1	376.3	570.3	1.01	0.03	Reliable
-0.67 0.22		3.31	-36.9	65.5	937.3	185.2	536.3	1.20	0.00	Reliable
-0.67 0.39		3.82	-2.2	59.7	992.9	233.5	591.0	1.34	0.00	Reliable
-0.23 0.12		3.70	78.2	28.8	712.9	285.4	527.3	2.89	0.07	Reliable
-0.70 0.16		3.14	-62.4	42.8	731.4	265.1	532.9	3.77	0.05	Reliable
-1.27 0.44		2.91	2.5	44.0	715.7	270.5	557.0	4.42	0.07	Reliable
-0.56 0.37		4.60	-97.8	44.9	840.7	121.8	550.0	5.86	0.01	Reliable
-0.80 0.28		2.22	1.6	59.2	649.8	248.7	546.5	15.98	0.14	Reliable
-0.24 0.13		3.26	43.9	58.1	672.8	55.3	530.7	19.70	0.11	Reliable
-0.33 0.21		1.78	6.2	49.1	601.4	305.4	540.9	20.01	0.24	Reliable
-1.14 0.48		4.35	-205.1	33.0	649.1	170.0	562.3	21.30	0.14	Reliable
-0.99 0.35		4.56	131.2	30.1	622.7	162.0	545.8	23.69	0.19	Reliable
-0.59 0.12		5.88	88.0	16.5	606.7	193.1	579.8	43.24	0.23	Reliable

Table III.1. Stellar and kinematic parameters for the 20 HVS candidates.

 $^{\mathrm{a}}$ Note there are large uncertainties in these measurements at the S/N of these candidates.

 $^{\rm b}{\rm Radial}$ velocity, in km s⁻¹, is taken straight from the SSPP without any corrections.

^cTotal proper motion, in mas yr⁻¹, is calculated from the μ_{RA} and μ_{Dec} values listed in the CAS without any corrections.

 d This is the total velocity of the candidate, in km s⁻¹, after conversion to the Galactic Cartesian coordinate system described in §4.1, before any error consideration.

 $^{\circ}$ The value here is the minimum total velocity, in km s⁻¹, determined from a million realizations of proper motion, radial velocity, and distance errors drawn from their distributions described in Section 3.3.

 $^{\rm f}{\rm The}$ escape velocity, in km ${\rm s}^{-1},$ at the star's position in a spherical potential.

3.6.2 Alternative Origins

As shown in Figure 3.4, none of the HVS candidates are coming from the Galactic center or from the direction of M31. The popular ejection mechanisms described inSection 3.1 involve a central SMBH and cannot explain these stars. The question where these stars originated, and how they gained such high velocities, remains.

One of the best-known hypervelocity mechanisms involves a binary system in the disk, in which a supernova explosion ejects the companion star (e.g.; Blaauw, 1961). There are many lesser known hypervelocity ejection mechanisms as well. For example, multibody ejections from the dense central regions of globular clusters (e.g.; Poveda et al., 1967) including globular clusters that may have dissipated over the lifetime of the Galaxy (e.g.; Gnedin & Ostriker, 1997; Chernoff & Weinberg, 1990; McLaughlin & Fall, 2008) may boost a star to hypervelocities. In addition, there could be a three-body interaction involving an intermediate-mass black hole or otherwise very massive star (e.g.; Gvaramadze et al., 2009). A final hypervelocity ejection mechanism involving a stellar dynamical process could be the partial tidal disruption of a single star around a SMBH (Manukian et al., 2013).

Furthermore, three-body interactions between galaxies, such as M31 (e.g.; Caldwell et al., 2010) and the Large and Small Magellanic Clouds (e.g.; Chandar et al., 2010), have been suggested as possible hypervelocity ejection mechanisms, although we have already ruled out M31 specifically. HVSs may also receive an energy boost during the tidal stripping process as long streams are stripped from an accreted satellite (e.g.; Abadi et al., 2009; Caldwell et al., 2010; Piffl et al., 2011; Fouquet et al., 2012; King et al., 2012).

3.6.3 Follow-up Analysis

A significant fraction of the candidates failed the dist22 > 7 requirement, meaning that photometric blending from a near neighbor may have affected the proper-motion determination. A larger number of the candidates suffer from too few detections in the SDSS+USNO-B catalog. Confirming these candidates as HVSs would require additional astrometric analysis in order to verify their proper motions; the *Hubble Space Telescope* Fine Guidance Sensor may be appropriate.

There is also the possibility that these candidates are unresolved spectroscopic binaries, which could imprint a large radial velocity signal. Future, higher resolution spectroscopic observations could easily decide this issue and would also allow a more detailed chemical analysis to shed light on their origins.

3.6.4 Constraints on the Initial Mass Function

The fact that we find no low-mass HVSs coming from the Galactic center continues to pose a problem for a universal initial mass function and an unbiased binary ejection mechanism. If we simply assume a Salpeter initial mass function and a mass-blind dynamical process, we would naively expect roughly 150 HVSs in the 0.6–1.2 solar mass range in our sample, compared with the 14 known 3–4 M_{\odot} HVSs (Brown et al., 2009). Either the initial mass function near the Galactic center is top-heavy or the process acting at the Galactic center ejects over 10 times more high-mass stars than low-mass ones. There is tentative evidence from the Arches and other young star clusters at the Galactic center that the initial mass function is top-heavy, with a slope of about -1.6 (Figer et al., 1999), although this is a matter of debate. If we adopt this slope for our initial mass function, then we still should have observed roughly 40 HVSs with spectral types G and K, which would require an ejection mechanism that favors massive stars by more than factor of 3^2 .

Our constraints on the initial mass function are consistent with the findings of Kollmeier et al. (2010), who searched for metal-rich F/G halo HVSs. This earlier study placed stricter limits on the ejection mechanism, however, because F/G stars would be expected to accumulate in the halo over their main-sequence lifetimes, while our sample probes only stars passing through the solar neighborhood; stars ejected from the Galactic center through stellar binary disruption, for example, would reach and pass through our sample in mere tens of millions of years. Our results are also consistent with the constraints from Zhang et al. (2013), who considered the S stars at the Galactic center to be the captured companions of binary star tidal disruption, a process that ejects the second star.

²We expand on the determinations of the expected number of HVSs in Appendix D.

3.7 Summary and Conclusions

We report a set of 20 hypervelocity candidates from the SEGUE G and K dwarf sample. These candidates have velocities greatly exceeding the escape velocity at their respective positions in the Galaxy, albeit with large proper-motion errors. Monte Carlo estimates of the position and kinematics of these stars show that seven of the 20 exceed the escape velocity at their respective locations within the Galaxy with at least 98% probability and that each candidate's interloper likelihood is less than 25%.

Surprisingly, an orbit analysis indicates that none were ejected from the Galactic center. The confirmation of these candidates as HVSs argues for a more careful exploration of alternative ejection mechanisms such as interactions within globular clusters, dwarf galaxies, or tidal tails, as well as ejections from supernovae in the Galactic disk.

If these stars are truly hypervelocity, their spectra could already contain clues to their origin. For example, abundance patterns indicative of supernova contamination would confirm or rule out a candidate's having been ejected from a high-mass binary system (Przybilla et al., 2008).

One remaining question is why these stars were not identified in previous HVS campaigns. A possibility is that prior searches focused on extreme radial velocities (e.g.; Brown et al., 2005). While the radial velocities of our candidates are relatively modest, it is the addition of proper motions that boosts these stars into hypervelocity candidacy. Naturally, our sample also explored a cooler spectral type than previous work. Future surveys may be more successful in identifying HVSs with both radial velocity and propermotion measurements. We are expanding our search for HVS candidates to the entirety of the SDSS DR9 sample in order to include all spectral types. Analysis of any additional candidates identified in this search, as well as follow-up, is deferred to a future paper.



Figure 3.5: Normalized metallicity distribution functions of our candidates (shaded), compared with G dwarfs (red), K dwarfs (green), globular clusters (blue), and the Galactic bulge (cyan).

Chapter IV

WORK IN PROGRESS

4.1 A Global Search for G/K-type Hypervelocity Stars

4.1.1 Motivation

The G- & K-dwarf sample of Schlesinger et al. (2012), used in Palladino et al. (2014), was carefully selected to be representative of the disk around the Solar neighborhood. This sample, while statistically robust, introduces selection biases in a search for hypervelocity stars. Even if a HVS is ejected through interactions within the disk, the timescales on which HVSs traverse the Galaxy are such that we should not necessarily expect to find them lingering in the Solar neighborhood.

A true HVS found within the disk may be explained by a very recent, nearby, disk ejection (e.g.; Blaauw, 1961). Alternatively, the HVS could be 'just passing through' as a hypervelocity interloper originating from another region of the Galaxy, likely the bulge (e.g.; Hills, 1988; Brown et al., 2005, 2012). The area of the Solar neighborhood targeted by the Schlesinger et al. (2012) sample is such that either scenario would not contribute significantly to the total HVS population. Conversely, assuming an un-biased ejection geometry, regardless of origination site, it is more likely that HVSs are quickly flung out into the halo. So, while there are many advantages of using the G- and K-dwarf sample of Schlesinger et al. (2012), this section is dedicated to performing a global search for G/Ktype HVSs within the SEGUE database by loosening the selection criteria to include stars outside the disk as well.

4.1.2 Candidate Selection

The global search for G/K-type hypervelocity stars was performed with a query of the twelfth and final data release of SDSS-III by selecting stars with extinction corrected $(g-r)_0$ colors and r_0 magnitudes indicative of G and K-type stars, as described in Section 3.3. The query returned 120,679 G-type stars and 173,790 K-type stars.

Subsequently, we remove any star for which effective temperature, surface gravity, spectral type, metallicity, or radial velocity were not determined (ie. values of -9999 or "unknown"). This cut results in 130,990 targets. We follow the additional quality assurances of Schlesinger et al. (2012) which also considered poor signal-to-noise, excessive reddening, saturation, and flags assigned by the SSPP. We do not, however, consider the magnitude limit and distance cuts of Schlesinger et al. (2012). Then, we apply the proper motion quality criteria of Kilic et al. (2006) and retain only those stars that satisfy the "clean" designation, described in section 3.4.1. This step results in a sample of 103,238 G/K-type stars.

An estimation of the distance to each star was determined by employing the distance modulus with average absolute magnitudes assigned by the following equation (Ivezić et al., 2008):

$$M_r = -2.85 + 6.29(g-i) - 2.30(g-i)^2 + \Delta M_r, \tag{4.1}$$

where

$$\Delta M_r = 4.50 - 1.11 [Fe/H] - 0.18 [Fe/H]^2.$$
(4.2)

Absolute magnitudes determined from this method are accurate to within 0.1 magnitude

only for stars with $0.3 < (g-i) \le 1.0$ (Ivezić et al., 2008). Finally, dwarfs were separated from giants with a simple log(g) cut at 3.5; dwarfs with log(g) > 3.5 and giants with log(g) < 3.5. There are 85,268 G- and K-dwarfs (17,970 giants).

To identify hypervelocity star candidates, we convert the radial and tangential velocities to Galactic Cartesian coordinates and choose the same initial total velocity threshold of 600 km/s as in Section 3.3, and a maximum velocity cut of 10,000 km/s. Again, as in Palladino et al. (2014), we require that each candidate exceed the escape velocity at its current location within the Galaxy. This procedure yields 28 HVS candidates.

4.1.3 Preliminary Results

As an initial comparison of the results of the global search to the sample of Schlesinger et al. (2012), we compare the positions of the samples on the sky in Figure 4.1. We note that both samples cover the SDSS footprint. The differences in the bottom panels of Figure 4.1 are due to the distance cuts placed by Schlesinger et al. (2012) to isolate a population of disk stars.

The HVS candidates presented in Palladino et al. (2014) have transverse velocities significantly greater than their radial velocities (cyan 5-pointed stars in Figure 4.2), an indication of a sample contaminated by large proper motion errors. We perform the same comparison of transverse and radial velocities for the HVS candidates identified from the global search for G/K-type stars (magenta 5-pointed stars in Figure 4.2). We find that this sample is not biased in velocity– if anything, there are more candidates with velocities dominated by the radial component, which is a reasonable expectation for a population of HVSs zipping out of the Galaxy. Finally, at this stage of the investigation we look at the distributions of total velocity, Figure 4.3, and metallicity, Figure 4.4. We find the majority of the global-search HVS candidates have comparable velocities to the HVS candidates published in Palladino et al. (2014).

We also note that the metallicity distribution is bimodal (potentially trimodal) and does not strongly resemble any of the comparison populations– published HVS candidates, disk G/K-dwarfs, globular clusters, and bulge stars. The high-metallicity peak at [Fe/H] >-0.5 is especially interesting. Perhaps we are seeing a composite population of artificially velocity-boosted halo field stars and ejected bulge stars.

4.1.4 Discussion

The results presented in the previous section are preliminary and require further consideration. One critical question that needs to be addressed before publication is: If the Schlesinger et al. (2012) sample of G/K-dwarfs is only a subset of the global search for G/K-type stars, why did we not recover the HVS candidates of Palladino et al. (2014)?

Other, marginally less perplexing, questions include: Why did we need to enforce a maximum velocity cut of 10,000 km/s? and: Are the distances determined by equation 4.1 compared to the distances determined by Schlesinger et al. (2012), especially considering that we did not restrain the g - i color of the sample to the range deemed valid for equation 4.1 (Ivezić et al., 2008)?

Once the initial candidate selection is refined, we can start to answer the more interesting science questions. To establish the probability that the candidates are HVSs requires errors on the distance determinations, which at present, we do not have (but can obtain by convolving the errors on the stellar parameters). Orbits for each candidate will need to be calculated to determine their place of origin, and whether a SMBH ejection can explain any or all of the new candidates. Furthermore, as an initial pass, we segregated the giants out for a fair dwarf-to-dwarf comparison. However, we ultimately want to consider the giants in the global sample as well and what they might contribute to the unbound stellar population in the Milky Way.

4.2 F-type Hypervelocity Star Candidates

4.2.1 Motivation

Since there are 18 confirmed O/B-type HVS, consistent with SMBH ejection (Brown et al., 2012) we know that the Hills mechanism operates at the center of the Galaxy. If we add that Palladino et al. (2014) found zero low-mass HVS candidates originating from the Galactic Center, we have compelling evidence that the ejection mechanism may preferentially boost only the most massive stars to hypervelocities and/or that the initial mass function at the Galactic center is extremely top-heavy. Now, we search for HVS candidates within a sample of SDSS F-type stars as an attempt to populate the intermediate mass range to constrain the mass-dependence of the ejection mechanism.

The sample we use for this search is adapted from the sample of Allende Prieto et al. (2014), originally selected to study abundances in the Sloan survey. I combine the distances determined by Fernandez-Alvar et al. (in prep) with proper motions, radial velocities, and stellar parameters from the SDSS Catalog Archive Server (CAS). The sample contains 12,673 F-type stars.

4.2.2 Candidate Selection

From the initial sample of 12,673 F-type stars, we perform a similar cleaning procedure as described in the sections above. First, we remove stars with undetermined effective temperature, metallicity, or surface gravity. Then, we apply the proper motion quality criteria of Kilic et al. (2006) and retain only those stars that satisfy the "clean" designation, described in section 3.4.1. This step results in a sample of 11,779 F-type stars.

Similarly, we separate dwarfs from giants with a cut on surface gravity at 3.5. This cut returns 10,590 F-dwarfs and 1,189 F-giants. For this sample, we rely on the distances determined by Fernandez-Alvar et al. (in prep).

F-dwarf hypervelocity star candidates are identified via the same procedure outlined in section 4.1.2. We find 98 preliminary HVS candidates with v > 600 km/s and 95 candidates with $v > V_{esc}$. With both distance and distance errors provided by Fernandez-Alvar et al. (in prep) we are able to initiate the statistical analysis stages of the investigation. Via Monte Carlo simulations, described in detail in section 3.4.2, we find that 34 of the preliminary F-dwarf HVS candidates are unbound at least 90 % of the time and 8 are unbound at least 95% of the time.

4.2.3 Preliminary Results

We identified 95 F-dwarf HVS candidates with total velocities exceeding the escape velocity at their respective locations. We performed the first statistical test to determine each candidate's probability of being bound to the Galaxy, described in Palladino et al. (2014) and section 3.4.2. We identified 34 F-dwarf HVS candidates that are unbound with at least 90% probability and 8 F-dwarf HVS candidates with at least 95% probability.

Figure 4.5 shows the velocity distribution of the 8 highly unbound candidates compared to the sample of 95 F-dwarf candidates. The F sample yields an overall similar distribution to that of the global G- and K-dwarf sample in section 4.1– the majority of the candidates with $v_{tot} < 2000$ km/s and a tail extending to much higher velocities at ~ 6000 km/s.

Figure 4.6 shows the metallicity distribution of the 95 F-dwarf HVS candidates and the 8 candidates that comprise the highly unbound subgroup. The F-dwarf HVS candidates are much more metal poor than even the F sample from which they were selected. However, the 8 highly unbound candidates, with a peak at $[Fe/H] \sim -1.7$, roughly align with the metal-poor peaks of both the F sample and globular cluster distributions.

4.2.4 Discussion

The results presented in the previous section are preliminary and require further consideration. We are currently awaiting updated distance determinations for the F-star sample. Once in hand, we will be able to fully characterize any HVS candidates within the sample as was done in Palladino et al. (2014), including orbits and all statistical tests. We would also like to compare the distances determined by the methods of Allende Prieto et al. (2006, 2014) and Fernandez-Alvar et al. (in prep) to the distances determined by equation 4.1.

Again, we plan to ultimately consider F-giants as well, but for reasons outlined in Section 4.1.2 we postpone their analysis to a later date.



Figure 4.1: Top: Equatorial, RA and Dec, positions of the G/K-type stars returned by the global search (black) and the G/K-dwarfs of Schlesinger et al. (2012) (red). Both samples probe all areas of the SDSS footprint. Bottom left: Cartesian, R and z, positions of the two samples. Bottom right: Same as bottom left, zoomed in to see the extent of the Schlesinger et al. (2012) sample. The differences in the bottom two panels are due to the distance cuts applied by Schlesinger et al. (2012).



Figure 4.2: The transverse velocity, v_t , compared to the radial velocity, v_r , for the HVS candidates of Palladino et al. (2014) (cyan stars) and the global-search HVS candidates (magenta stars). The red lines indicate $v_t \sqrt{2}$ times larger than v_r , expected for an isotropic stellar population. The blue lines represent v_t 5 times larger than v_r , indicative of a sample dominated by the transverse velocity. The vast majority of the cyan points fall between the blue lines, suggesting that those candidates may suffer from large proper motion errors. Conversely, the magenta points are generally distributed throughout, suggestive of an unbiased sample. Upon close inspection, roughly half of the global-search HVS candidates fall below the red lines– a strong indication that they follow radial orbits, not unexpected for HVSs.



Figure 4.3: The velocity distribution of the global-search HVS candidates. All of the Palladino et al. (2014) HVS candidates fall within the first bin, with v < 1500 km/s. The majority of the HVS candidates represented here have similar velocities to the published candidates.



Figure 4.4: The metallicity distribution of the global-search HVS candidates (magenta) and the published HVS candidates (grey) compared to disk G/K-dwarfs (top), globular clusters (middle), and bulge stars (bottom). While the published HVS largely follow the distribution of the disk sample, the globalsearch candidates do not closely resemble any of the populations shown here.



Figure 4.5: The velocity distribution of the 95 F-dwarf HVS candidates (blue) and the 8 candidates with at least 95% probability of being unbound.



Figure 4.6: The metallicity distribution of the 95 F-dwarf HVS candidates (grey) and the 8 candidates that are unbound with at least 95% probability (magenta) compared to the entire F-star sample (black, top), globular clusters (blue, middle), and bulge stars (cyan, bottom). We note that the majority of the F-dwarf HVS candidates are extremely metal poor with only a couple relatively metal rich at [Fe/H] ~ -0.5 . We caution, however, that the grid edge of the model atmospheres is at [Fe/H] = -2.5; any candidate more negative than the vertical red line is to be ignored (Allende Prieto et al., 2014, private communication).

Chapter V

CONCLUSION

In this thesis I described two complementary techniques to mine the Sloan Digital Sky Survey (SDSS) data to search for stars that are unbound to the Milky Way.

First, we developed a technique to identify distant M-giant intragroup stars (IGSs) within SDSS based solely on color. Using this technique, we identified 700 IGS candidates, between 300 kpc and 2 Mpc (Palladino et al., 2012). These IGS may constitute rare tracers of an underlying intracluster light (ICL) population surrounding the Milky Way. One possible explanation for the origin of these IGS is that they were ejected from the center of the Galaxy through interactions with our supermassive black hole as hypervelocity stars (HVSs).

Secondly, we identified candidate HVSs from a sample of SEGUE G- and K-dwarfs. We found that nearly half of the candidates exceed their escape velocities with at least 98% probability and no candidate's orbit is consistent with a Galactic Center origin (Palladino et al., 2014). The lack of HVS candidates originating from the Galactic Center is an indication that either the ejection mechanism is mass-dependent or the initial mass function at the center of the Galaxy is even more top-heavy than suggested in the literature (e.g.; Figer et al., 1999, Section 3.6.4).

We plan to continue the search for stars that comprise the unbound stellar population in the Milky Way. We are currently searching for HVS candidates in an extended G- and K-type sample that includes both dwarfs and giants and probes regions of the Galaxy beyond the disk. We are also searching for more HVS candidates within a sample of Fstars, containing dwarfs and giants, to constrain the mass-dependence of the supermassive black hole ejection mechanism operating at the center of the Milky Way.

Appendix A

ESTIMATION OF EXPECTED HVS VIA THE BINARY DISRUPTION MECHANISM.

Here we estimate on the number of G/K-type HVS we should expect to observe within SDSS resulting from the supernova binary disruption mechanism described in Section 1.4 and Tauris (2015).

In order to be consistent with the scenario described in Tauris (2015), we consider only type Ia supernovae for this estimation. We begin with the known supernova rate of 1 supernova per 100 years and the ratio of type II to type Ia supernova of ~ 3.5 (Sato et al., 2007), which gives us a rate of 1 type Ia supernovae per 350 years. Then, adopting an average velocity for our HVS candidates of 1,000 km/s, we determine it would take 5 million years for a star to travel from the disk to a distance of 5 kpc (consistent with our observations). Therefore, over a period of 5 million years, we would expect approximately 14,000 type Ia supernovae in the Milky Way.

From here, we consider the fraction of the sky observed by SDSS, and determine that there would be $\sim 3,000$ type Ia supernovae in the SDSS footprint from the past 5 million years.

Since, by definition, type Ia supernovae exist in binary systems, we must consider how many of these 3,000 binaries contain a low mass companion. Standard population synthesis models adopt a flat distribution for assigning the mass to the secondary stellar component of binaries (Abt, 1983)– if we group all spectral types (O/B, A/F, G/K, M/later), we will assume that roughly 1/4 of these binary systems contain G/K-type companions. This assumption determines that we should expect ~ 750 G/K-type companions to type Ia supernovae within the SDSS footprint from the past 5 million years.

According to Tauris (2015), only 1% of binary disruption scenarios with a low mass companion will result in an ejection speed of v > 600 km/s. This velocity requirement yields that we should expect to observe O(10) low mass HVS from a supernova binary disruption scenario within the SDSS footprint, which is consistent with our result of 20 candidates.

Appendix B

EXPANSION ON THE EXPECTED IGS CALCULATION.

Here we elaborate on the calculation outline in Section 2.5 determining the expected number of IGS in the SDSS footprint.

If we assume that the SMBH in our galaxy had a merger with an intermediate mass black hole (IMBH) with mass $3 \times 10^5 M_{\odot}$, it would 3-body scatter stars and eject them from the system isotropically. This means that the total mass of ejected stars will also be $3 \times 10^5 M_{\odot}$ (Yu & Tremaine, 2003). Assuming an average stellar mass of 1 M_{\odot} , then 3×10^5 stars will be ejected.

The total sky has 41,253 square degrees, which yields approximately 7 ejected IGS per square degree. Assuming the IGS formed 10 Gyr ago with a Salpeter IMF, then there would be 0.00124 red giant HVS per square degree. The area of sky observed by SDSS is 8000 square degrees, so this yields the expected tens of red giant IGS in the SDSS footprint resulting from SMBH ejection.

Appendix C

IGS CANDIDATES REMAINING AFTER ALL CRITERIA CUTS.

Table C.1. Complete and unabridged.

Object ID	RA	DEC	u	u err	g	g err	r	r err	i	i err	\mathbf{z}	z err	r-i	i-z
758882136836343139	60.5013	80.8760	24.4	1.3	24.7	0.6	24.3	0.7	22.0	0.2	20.3	0.2	2.3	1.7
758882137910740030	62.9212	82.5298	25.2	1.2	25.0	0.6	23.2	0.4	21.0	0.1	19.6	0.1	2.2	1.4
758882626993718584	63.3170	81.6950	26.5	0.6	24.8	0.6	25.3	0.7	22.0	0.2	20.3	0.2	3.3	1.6
758877527803168329	94.6323	63.7038	25.1	0.9	24.1	0.3	24.5	0.5	21.2	0.1	19.8	0.1	3.3	1.4
758877527266231607	94.8017	63.2633	24.6	0.7	24.4	0.5	24.5	0.6	21.8	0.1	20.4	0.1	2.7	1.4
758877527266493955	95.9369	63.5501	25.2	0.7	24.8	0.6	25.6	0.5	22.7	0.3	20.7	0.2	2.9	2.0
758878272976455144	105.3459	66.8553	25.3	0.9	24.4	0.4	24.1	0.4	21.8	0.1	20.2	0.1	2.3	1.6
758878271902778853	105.8427	66.0785	23.8	0.7	24.7	0.4	24.6	0.5	22.3	0.2	20.8	0.2	2.3	1.5
758884768580109918	107.9936	38.3022	24.8	1.1	25.1	0.7	24.4	0.6	21.9	0.1	20.4	0.2	2.6	1.4
587738067260998978	109.5721	39.4395	25.3	0.9	25.7	0.6	24.6	0.7	22.1	0.2	20.0	0.1	2.5	2.1
587738066187126136	110.3269	38.8813	23.9	1.2	24.8	0.6	24.3	0.6	22.0	0.2	20.6	0.2	2.2	1.5
758884821194311430	111.7235	31.6748	23.3	0.5	24.9	0.5	23.8	0.4	21.4	0.1	20.0	0.1	2.4	1.4
758884803478619431	112.4325	31.0977	25.2	1.1	24.5	0.5	24.1	0.6	21.9	0.1	20.5	0.1	2.2	1.4
587737808490071190	112.6070	38.1797	25.4	1.0	25.2	0.6	24.8	0.8	22.2	0.2	20.7	0.2	2.6	1.6
587728906096674179	112.7248	27.9531	24.8	0.9	25.1	0.6	25.1	0.6	21.9	0.1	20.5	0.2	3.2	1.4
587725551190869074	113.3506	36.8066	23.7	0.7	24.8	0.3	25.1	0.8	21.9	0.1	20.4	0.1	3.3	1.5
587727866178372734	113.9236	31.5228	23.8	0.7	25.3	0.6	24.2	0.5	22.0	0.2	20.4	0.1	2.2	1.5
587725552265790733	114.2540	39.4958	23.7	0.7	24.7	0.6	25.0	0.7	22.1	0.2	20.8	0.2	2.9	1.4
587725775066563723	114.3135	39.3165	25.3	0.8	25.4	0.6	24.8	0.7	21.8	0.1	20.1	0.1	3.0	1.7
588013382718391798	115.5088	23.9862	25.6	0.6	25.1	0.5	24.2	0.5	22.0	0.1	20.5	0.1	2.2	1.4
587732152555078919	115.7199	22.3048	23.7	0.9	26.6	0.4	23.8	0.5	21.2	0.1	19.3	0.1	2.7	1.9
587725775067677784	115.9952	41.5387	25.5	0.8	25.3	0.6	24.0	0.5	21.6	0.1	20.3	0.1	2.3	1.3
587732054308947237	116.6473	26.0150	23.7	0.7	24.9	0.6	24.1	0.6	21.3	0.1	19.9	0.1	2.7	1.4
587725774530872383	116.6594	41.4087	24.1	0.8	23.9	0.4	24.0	0.6	21.8	0.1	20.4	0.1	2.3	1.3
588016878822295117	116.7229	18.5509	25.8	0.6	24.8	0.6	24.4	0.6	21.6	0.1	20.0	0.1	2.9	1.5
587737826210219010	116.7903	45.1563	24.7	1.0	24.7	0.6	24.8	0.7	21.9	0.1	20.5	0.2	3.0	1.3
587728906099623039	117.2504	33.4457	24.1	0.9	25.3	0.6	23.9	0.5	21.0	0.1	19.2	0.1	2.9	1.7
Table C.1 (cont'd)

Object ID	RA	DEC	u	u err	g	g err	r	r err	i	i err	\mathbf{z}	z err	r-i	i-z
587738372742121166	117.5023	18.4613	24.8	0.8	25.5	0.5	24.9	0.5	22.7	0.3	20.8	0.2	2.2	1.9
588013383793771999	117.5227	27.3834	24.2	0.8	24.5	0.5	24.9	0.6	22.7	0.2	21.1	0.2	2.1	1.6
587728906636952618	117.5359	34.5913	24.9	0.9	24.5	0.5	24.3	0.6	21.6	0.1	20.2	0.1	2.7	1.4
587731679573509088	117.6023	27.9611	25.8	1.2	25.7	1.0	23.5	0.9	20.9	0.1	19.0	0.1	2.6	1.9
587735043069707709	117.7701	20.8646	24.7	0.9	25.2	0.6	25.6	0.5	22.4	0.2	20.9	0.2	3.3	1.4
587728931334653093	117.9390	33.6761	25.0	0.9	24.1	0.4	23.9	0.5	20.9	0.1	19.4	0.1	3.0	1.4
587735043069969790	118.3476	21.2348	25.0	0.8	25.2	0.6	24.5	0.6	22.3	0.2	20.9	0.2	2.2	1.4
587737809568531492	118.6204	48.1006	23.7	0.8	25.4	0.5	24.3	0.6	21.9	0.1	20.4	0.2	2.5	1.4
588016878286406831	118.8410	19.7030	25.0	1.4	25.0	0.8	25.1	0.7	22.0	0.2	20.5	0.2	3.2	1.4
587731873385481740	118.9846	32.3811	25.5	0.6	24.6	0.4	24.8	0.5	21.7	0.1	20.1	0.1	3.2	1.5
587735235806561437	118.9877	21.3134	23.5	0.8	24.9	0.6	24.7	0.7	21.4	0.1	20.0	0.1	3.3	1.4
587735236343759987	119.1972	22.2818	25.2	0.8	25.4	0.6	25.3	0.6	22.3	0.2	20.7	0.2	3.0	1.5
587738066730484820	119.4568	52.3041	24.7	0.9	24.7	0.5	23.9	0.4	21.6	0.1	20.2	0.1	2.3	1.4
758877528344757598	119.8005	65.7236	24.3	0.7	25.2	0.4	25.0	0.5	22.4	0.1	20.4	0.1	2.6	1.9
587739115234788850	119.8104	17.9973	24.0	0.8	24.7	0.4	24.0	0.4	21.6	0.1	20.1	0.1	2.4	1.5
587727867256243248	119.8520	39.6771	24.5	0.9	24.4	0.5	24.1	0.5	20.9	0.1	19.4	0.1	3.2	1.5
587742010042287707	120.2735	10.3882	23.7	1.0	24.7	0.6	24.8	0.7	22.1	0.2	20.4	0.2	2.7	1.6
588297863634682833	120.3658	24.4537	25.7	1.0	25.2	0.8	24.3	0.8	22.0	0.2	20.6	0.2	2.4	1.4
587725470665081929	120.4194	43.3815	23.5	0.6	25.6	0.5	24.2	0.4	21.5	0.1	19.9	0.1	2.7	1.6
587737826749645877	120.5762	50.4355	24.1	1.0	24.8	0.3	23.7	0.4	21.5	0.1	20.1	0.1	2.2	1.4
587741708859803246	120.6010	12.5824	25.5	0.6	25.1	0.5	23.5	0.3	21.3	0.1	19.9	0.1	2.2	1.4
587742010042549539	120.7911	10.6056	23.1	0.6	23.7	0.4	24.0	0.6	21.8	0.1	20.4	0.2	2.2	1.4
588013382184469831	120.8817	28.7355	23.8	0.4	25.0	0.6	25.7	0.4	22.9	0.2	21.1	0.2	2.8	1.8
587732469846377586	120.9695	26.5794	25.7	0.8	24.0	0.4	23.6	0.4	20.9	0.1	19.5	0.1	2.7	1.4
587745244689073605	121.0114	9.6150	24.5	0.8	24.5	0.4	23.7	0.3	21.1	0.1	19.8	0.1	2.5	1.3
587728669878912102	121.1313	40.8448	25.2	0.8	24.7	0.5	25.0	0.6	22.3	0.2	20.9	0.2	2.7	1.4
588023046933120424	121.1645	13.0654	24.1	0.8	24.5	0.4	24.9	0.6	22.0	0.2	20.7	0.2	2.8	1.3

Table C.1 (cont'd)

Object ID	RA	DEC	u	u err	g	g err	r	r err	i	i err	\mathbf{Z}	$z \ err$	r-i	i-z
587745243615266573	121.2201	8.7283	24.6	0.8	25.3	0.5	25.0	0.5	21.8	0.1	20.2	0.1	3.3	1.6
587739115235509353	121.2323	18.8467	25.1	0.9	24.7	0.5	25.2	0.6	21.9	0.1	20.2	0.1	3.2	1.8
587745403603453761	121.4298	8.4593	24.9	0.7	24.7	0.4	25.2	0.4	22.6	0.2	21.2	0.2	2.6	1.4
587745402529711807	121.5256	7.7839	25.0	0.7	24.8	0.4	24.2	0.4	22.0	0.1	20.4	0.1	2.2	1.6
587744637488596380	121.5423	11.1817	25.5	1.0	25.2	0.6	24.2	0.7	21.7	0.1	20.3	0.1	2.5	1.4
587731520662930498	121.5799	31.5473	25.4	1.0	25.3	0.7	24.0	0.6	21.3	0.1	19.6	0.1	2.7	1.7
587741387273275005	121.6111	16.0509	25.1	0.8	24.6	0.4	24.8	0.5	21.9	0.1	20.5	0.2	2.9	1.5
587734623238227199	121.6396	26.8599	25.5	0.7	24.7	0.5	23.9	0.5	21.3	0.1	19.8	0.1	2.6	1.4
587738066194465721	121.7083	53.7864	23.7	0.9	25.4	0.5	24.9	0.7	22.1	0.2	20.5	0.1	2.8	1.6
587738372744217890	121.7230	21.3338	24.3	0.8	26.0	0.5	24.9	0.6	22.3	0.2	20.9	0.2	2.6	1.4
587741532764898777	122.0871	14.6230	24.4	0.8	25.7	0.4	25.0	0.6	22.2	0.2	20.9	0.2	2.7	1.4
587731887346025354	122.1712	36.4783	25.4	0.7	25.1	0.6	24.5	0.5	22.2	0.2	20.8	0.1	2.4	1.4
587725550659240849	122.3177	46.7939	25.6	0.7	24.1	0.5	23.2	0.4	20.9	0.1	19.3	0.1	2.3	1.6
587728906102506800	122.3662	38.6413	24.7	0.9	25.3	0.5	24.1	0.5	21.0	0.1	19.6	0.1	3.1	1.4
587739377229628867	122.4993	18.3702	24.9	1.1	25.5	0.6	23.6	0.4	21.1	0.1	19.5	0.1	2.5	1.5
587745243615921560	122.5645	9.1305	25.5	0.6	24.6	0.4	24.8	0.5	22.7	0.2	21.0	0.2	2.1	1.7
587744874248799448	122.5791	9.5182	24.9	1.1	25.1	0.6	23.8	0.5	20.9	0.1	19.5	0.1	2.9	1.5
587734621627614164	122.5965	25.8422	24.8	1.1	24.6	0.5	24.7	0.7	22.2	0.2	20.7	0.2	2.4	1.5
587738565479368364	122.6182	7.3218	24.8	1.0	24.6	0.5	25.3	0.5	22.5	0.2	21.0	0.2	2.8	1.6
587742010043467053	122.8561	11.3085	25.9	0.6	25.8	0.5	25.2	0.6	22.3	0.2	20.9	0.3	2.9	1.4
588007005234004760	122.8710	45.4871	24.4	1.1	24.8	0.7	24.7	0.7	22.0	0.2	20.6	0.2	2.7	1.4
587739153354851262	123.2158	19.9350	25.5	0.9	25.4	0.7	24.0	0.6	21.5	0.1	19.9	0.1	2.5	1.6
587742008969790819	123.3029	10.5153	25.1	1.2	24.0	0.4	24.2	0.6	21.8	0.1	20.0	0.1	2.3	1.8
587732471458694255	123.3611	30.2143	24.3	1.0	25.6	0.6	25.1	0.7	22.6	0.3	20.9	0.2	2.5	1.7
587742009506792701	123.3705	11.1259	24.7	1.0	25.8	0.6	24.3	0.6	21.9	0.1	20.4	0.1	2.4	1.5
587739153354916942	123.4933	19.9022	23.5	0.9	25.7	0.6	24.9	0.7	22.3	0.2	20.6	0.2	2.6	1.8
587741386737255953	123.5510	16.6804	25.2	0.7	25.0	0.4	24.5	0.5	22.1	0.1	20.7	0.2	2.4	1.4

Table C.1 (cont'd)

Object ID	RA	DEC	u	u err	g	g err	r	r err	i	i err	\mathbf{Z}	z err	r-i	i-z
587745403067565543	123.6343	8.6650	24.6	0.9	24.2	0.3	23.9	0.4	21.2	0.1	19.9	0.1	2.7	1.3
587734949653120190	123.7966	6.6361	24.9	1.2	24.1	0.5	24.5	0.9	22.3	0.3	20.4	0.2	2.2	1.8
587731522275312616	123.8046	35.2063	24.6	1.1	25.7	0.5	25.4	0.7	22.2	0.2	20.6	0.2	3.2	1.6
587745244690384353	123.8545	10.4778	25.0	0.7	25.1	0.5	23.9	0.4	21.7	0.1	20.3	0.1	2.1	1.5
588016878825964493	123.8681	23.6659	25.1	1.2	23.9	0.4	23.9	0.6	21.2	0.1	19.6	0.1	2.7	1.7
587735236346315834	124.2462	25.9692	25.5	0.8	25.9	0.6	25.7	0.7	22.7	1.0	20.8	0.2	2.9	1.9
587734949653382475	124.4021	6.8419	23.4	0.7	24.4	0.6	24.0	0.7	21.8	0.2	20.2	0.2	2.2	1.6
587742008970380597	124.5343	11.0003	25.2	1.2	25.1	0.6	24.0	0.5	21.9	0.1	20.5	0.1	2.2	1.3
588016879363359634	124.5851	24.7195	23.2	0.6	25.5	0.6	23.8	0.5	21.5	0.1	20.1	0.1	2.3	1.4
588010358525134744	124.6703	1.9613	24.7	1.2	25.5	0.6	25.2	0.5	22.7	0.2	20.7	0.2	2.5	2.0
587735236883645397	124.8579	26.9178	25.4	0.9	25.3	0.6	25.3	0.7	22.0	0.2	20.4	0.2	3.2	1.7
587738948271473620	125.1085	23.5913	25.0	1.1	25.5	0.6	24.5	0.7	22.0	0.2	20.6	0.2	2.5	1.3
588016839634584912	125.2505	23.8631	25.8	0.7	25.3	0.6	23.5	0.3	21.3	0.1	19.7	0.1	2.2	1.6
587737808497673062	125.4675	53.0010	25.4	1.1	24.6	0.6	24.6	0.9	21.3	0.1	19.7	0.1	3.3	1.6
587735235809969032	125.5140	26.3690	23.9	1.1	25.9	0.5	23.9	0.6	21.7	0.1	19.9	0.1	2.3	1.7
587741816771773524	125.5536	13.5298	26.0	0.7	24.1	0.4	24.4	0.6	22.2	0.2	20.8	0.2	2.2	1.5
587739152819029103	125.7847	20.7065	25.9	1.0	26.5	0.9	24.8	0.7	22.3	0.2	20.9	0.2	2.5	1.4
587741532229731745	125.8905	16.0249	24.1	1.0	25.1	0.6	24.3	0.6	21.7	0.1	19.8	0.1	2.6	1.9
587731885736723588	125.9387	37.9468	24.1	1.0	25.3	0.6	25.0	0.6	22.1	0.1	20.8	0.1	2.9	1.4
588016879901017043	125.9927	26.2273	23.8	0.8	25.0	0.7	24.3	0.6	21.8	0.1	20.1	0.1	2.6	1.7
587741387812373747	126.0358	18.8876	25.2	0.6	25.3	0.4	24.4	0.5	21.9	0.1	20.0	0.1	2.5	1.9
587732701777364516	126.4451	3.4755	23.6	0.9	25.2	0.7	24.1	0.6	22.0	0.1	20.6	0.2	2.1	1.4
587739376157590638	126.4929	19.4975	23.4	0.7	24.7	0.6	24.1	0.5	21.7	0.1	20.3	0.1	2.4	1.4
587741817309234100	126.8135	14.3474	23.8	0.8	24.7	0.6	23.5	0.4	21.2	0.1	19.8	0.1	2.3	1.4
587734621629776894	126.9443	29.0573	25.3	1.0	24.1	0.4	23.5	0.4	21.0	0.1	19.7	0.1	2.5	1.4
588848899352561319	126.9608	-0.5883	24.8	1.0	24.8	0.6	23.7	0.4	21.5	0.1	20.1	0.1	2.2	1.5
758877529419810240	127.3261	66.1607	24.4	0.7	25.3	0.4	23.5	0.3	20.9	0.1	19.1	0.0	2.5	1.8

Table C.1 (cont'd)

Object ID	RA	DEC	u	u err	g	g err	r	r err	i	i err	\mathbf{Z}	z err	r-i	i-z
588010358526444861	127.6105	2.2769	24.4	1.4	25.1	0.7	24.2	0.6	21.6	0.1	20.0	0.1	2.6	1.6
587742008971887729	127.7826	12.1561	23.8	1.0	25.5	0.6	24.7	0.6	22.0	0.1	20.5	0.1	2.7	1.6
588848900426761686	127.9468	0.2102	23.6	0.6	24.4	0.5	24.8	0.6	22.2	0.2	20.8	0.2	2.6	1.4
758885526635152958	128.0022	-3.5678	24.2	1.0	24.4	0.4	24.3	0.6	21.1	0.1	19.5	0.1	3.2	1.6
587741709937149050	128.0561	16.6285	23.3	0.4	25.3	0.5	24.3	0.6	21.3	0.1	19.7	0.1	3.0	1.5
587739406759560244	128.1348	20.5569	24.8	1.0	24.2	0.4	24.8	0.8	21.5	0.1	20.1	0.1	3.3	1.5
588013382188270667	128.2241	34.6591	25.2	0.9	24.7	0.5	24.1	0.4	21.7	0.1	20.1	0.1	2.4	1.6
587735343181989199	128.2882	6.4828	24.2	1.0	25.7	0.4	24.9	0.5	22.3	0.2	20.8	0.1	2.7	1.5
587739115775985007	128.5542	23.6851	23.8	0.7	24.8	0.5	24.1	0.4	21.6	0.1	20.1	0.1	2.5	1.5
587739406222820115	128.5628	20.3347	24.2	1.2	25.9	0.5	23.6	0.5	21.0	0.1	19.5	0.1	2.5	1.6
587725552273458067	128.6043	54.1334	24.9	1.1	24.5	0.6	24.7	0.8	21.8	0.2	20.4	0.2	3.0	1.4
587727900001437041	128.7573	0.7449	24.4	1.1	24.6	0.6	24.3	0.5	22.0	0.1	20.1	0.1	2.3	1.9
588010136263263191	129.1525	47.2654	25.0	0.9	24.1	0.5	23.9	0.5	21.7	0.1	20.3	0.2	2.2	1.4
587735240099890134	129.1914	28.7956	24.4	1.3	24.1	0.6	24.6	0.7	22.1	0.2	20.8	0.2	2.5	1.3
587745243618936116	129.2938	11.3592	25.0	0.9	24.7	0.5	25.2	0.5	22.7	0.2	20.7	0.2	2.5	2.0
587732702315611255	129.5683	4.1795	24.8	1.1	25.7	0.8	23.8	0.5	21.0	0.1	19.2	0.1	2.7	1.8
588017979409171295	129.7783	26.2118	25.6	0.6	24.5	0.5	25.1	0.6	22.3	0.2	20.8	0.2	2.9	1.5
587744637492462684	130.0276	14.4030	23.2	0.7	25.7	0.5	24.1	0.7	22.0	0.1	20.6	0.2	2.1	1.5
587741532768634037	130.1910	18.3188	24.6	0.9	25.2	0.5	24.1	0.5	21.4	0.1	19.7	0.1	2.7	1.7
588023239936181347	130.8433	15.1661	24.7	0.9	24.9	0.6	23.5	0.3	21.2	0.1	19.9	0.1	2.2	1.4
588010360138499214	130.8541	3.6405	25.7	0.8	25.4	0.7	23.7	0.5	21.2	0.1	19.9	0.1	2.4	1.3
587735343183168623	130.9081	6.8343	25.5	0.8	24.2	0.4	24.6	0.6	22.2	0.2	20.2	0.1	2.5	2.0
587738947200287788	130.9399	26.1878	24.5	1.0	25.5	0.6	24.5	0.6	22.3	0.2	20.9	0.2	2.2	1.3
588023239130940550	131.1204	14.8844	24.3	1.1	25.2	0.6	23.6	0.9	21.4	0.1	19.5	0.1	2.2	1.9
587732048940958642	131.1430	41.9367	24.5	1.4	25.1	0.5	24.1	0.6	21.9	0.1	20.3	0.1	2.2	1.7
587732578297316598	131.4154	4.5721	23.5	0.6	25.5	0.4	24.9	0.5	22.6	0.2	20.5	0.1	2.3	2.1
587731885739410417	131.7574	42.1904	24.9	1.3	25.2	0.6	24.5	0.7	22.3	0.2	20.8	0.2	2.1	1.5

Table C.1 (cont'd)

Object ID	RA	DEC	u	u err	g	g err	r	r err	i	i err	\mathbf{Z}	z err	r-i	i-z
588016879367095353	132.6092	29.4826	23.6	0.8	25.2	0.6	24.6	0.6	21.9	0.1	20.6	0.2	2.7	1.3
587745243083638011	132.8636	11.8510	23.8	0.9	26.4	0.4	23.6	0.4	21.4	0.1	19.9	0.1	2.2	1.5
587745244157445136	132.8730	12.7470	24.2	1.0	25.3	0.5	24.3	0.6	22.2	0.2	20.8	0.3	2.2	1.4
587734623243994049	133.7811	35.0230	25.2	0.8	24.6	0.6	24.3	0.7	22.0	0.2	20.0	0.1	2.3	2.0
587741390494172383	133.7948	20.9534	24.7	1.0	25.3	0.5	25.2	0.5	22.2	0.1	20.7	0.2	3.0	1.5
587739115241407869	133.8194	25.8542	24.9	0.9	24.7	0.5	24.6	0.5	22.4	0.1	20.9	0.2	2.2	1.5
587738067810517573	134.0615	64.5391	25.2	0.9	24.8	0.5	23.5	0.3	21.4	0.1	20.0	0.1	2.2	1.4
587728880868459565	134.1800	3.7857	25.6	0.8	25.3	0.7	24.8	0.7	22.2	0.2	20.7	0.2	2.6	1.5
587728931879781318	134.3119	48.0796	24.0	0.9	25.9	0.5	24.2	0.6	22.0	0.1	20.5	0.2	2.3	1.5
587738065662444681	134.9270	62.3799	24.5	1.0	26.1	0.5	25.4	0.6	22.4	0.2	20.8	0.2	2.9	1.6
587725551201682191	135.0932	56.9195	23.9	1.0	25.2	0.7	24.3	0.7	21.3	0.1	20.0	0.1	3.0	1.4
588848900966974589	135.5404	0.7874	25.3	0.7	24.5	0.6	24.8	0.6	22.6	0.3	20.8	0.2	2.2	1.8
588010931370787667	136.2548	-1.1901	25.9	0.9	25.9	0.8	23.9	0.8	21.1	0.1	19.6	0.1	2.8	1.5
587725074990892118	136.4803	0.0084	24.8	0.8	24.1	0.4	23.5	0.4	21.1	0.1	19.7	0.1	2.5	1.3
587739157111767895	136.4962	26.2652	23.1	0.7	24.9	0.6	23.4	0.4	21.1	0.1	19.5	0.1	2.3	1.6
587731887352513242	136.4992	46.7866	23.3	0.7	25.5	0.7	24.3	0.8	21.6	0.1	20.1	0.1	2.6	1.6
587735661006685535	136.6245	30.3807	25.2	0.8	24.9	0.5	24.6	0.6	22.1	0.1	20.7	0.1	2.5	1.5
587731521207075766	136.6842	43.0042	25.0	1.1	25.1	0.6	24.7	0.8	21.8	0.2	20.5	0.2	2.8	1.4
587737810113070028	137.0808	62.1737	25.4	0.7	24.5	0.4	25.0	0.7	21.9	0.1	20.4	0.1	3.1	1.5
588848899893953565	137.2123	-0.1872	25.2	0.7	24.7	0.6	24.7	0.6	21.7	0.1	20.3	0.1	3.0	1.5
587726031692956510	137.2299	0.9463	24.0	1.4	25.7	0.7	24.2	0.6	22.0	0.2	20.5	0.2	2.2	1.5
587732771572941745	137.6327	6.8926	23.8	0.9	24.5	0.5	24.2	0.6	21.7	0.1	20.2	0.2	2.5	1.5
588010360141644747	137.9911	4.1146	25.6	0.7	25.9	0.5	23.2	0.3	20.9	0.1	19.4	0.1	2.3	1.5
587732578837267406	138.4298	5.7164	25.2	0.9	25.8	0.6	24.8	0.8	21.9	0.2	20.5	0.1	2.9	1.4
587745402000376763	138.6059	11.5764	25.0	1.2	25.4	0.6	24.9	0.6	22.3	0.2	20.5	0.1	2.6	1.8
587745403611120842	138.7042	12.7301	24.4	0.8	24.6	0.5	23.4	0.3	21.0	0.1	19.6	0.1	2.4	1.5
588010931372229514	139.5752	-1.3058	24.8	1.4	24.7	0.8	25.0	0.9	21.9	0.2	20.5	0.2	3.2	1.3

Table C.1 (cont'd)

Object ID	RA	DEC	u	u err	g	g err	r	r err	i	i err	\mathbf{z}	$z \ err$	r-i	i-z
587739116317901722	139.6774	29.3276	24.7	1.0	24.8	0.5	23.7	0.4	21.6	0.1	20.2	0.1	2.1	1.4
587739377237623666	139.9974	26.7931	25.1	1.0	24.0	0.4	25.0	0.6	21.8	0.1	20.4	0.2	3.2	1.4
588009366936683268	140.1298	53.4593	23.9	1.0	25.3	0.6	24.6	0.7	22.1	0.2	20.2	0.2	2.5	1.8
587741532236350353	140.7038	21.8742	25.0	1.0	25.3	0.5	24.8	0.6	22.1	0.1	20.3	0.1	2.7	1.7
588023046942032976	140.8789	20.6767	23.8	0.9	24.9	0.6	24.3	0.5	22.0	0.1	20.6	0.2	2.3	1.4
587742062128530595	141.0222	17.1453	25.0	0.8	26.0	0.4	24.8	0.6	22.5	0.2	21.0	0.2	2.3	1.5
587735661008782311	141.4131	32.7343	24.5	0.9	25.3	0.5	25.1	0.5	22.5	0.2	21.2	0.2	2.5	1.4
587725084112126989	141.8161	-1.2311	24.9	0.9	23.9	0.3	24.7	0.7	21.8	0.1	20.2	0.1	2.9	1.6
587732048408216470	141.9283	47.6861	25.3	1.0	26.1	0.5	23.4	0.4	20.9	0.1	19.6	0.1	2.5	1.3
587741816779048079	142.0450	18.8744	24.1	0.9	25.2	0.4	23.6	0.3	20.8	0.0	19.3	0.1	2.8	1.6
587725082501710856	142.2355	-2.5049	23.1	0.5	25.2	0.7	24.9	0.6	22.1	0.1	20.5	0.1	2.8	1.6
587742061055509536	142.9936	16.6777	23.6	0.6	25.3	0.5	23.4	0.3	21.3	0.1	19.6	0.1	2.1	1.7
587741490361992228	143.6895	23.6026	24.5	0.9	24.9	0.5	24.2	0.4	21.7	0.1	20.2	0.2	2.5	1.6
588009368548868956	143.8196	56.6921	25.0	1.2	24.5	0.5	25.1	0.8	22.0	0.2	20.3	0.2	3.1	1.8
587741821600007462	143.8346	20.0941	25.0	0.6	24.9	0.4	25.5	0.4	23.2	0.3	21.2	1.0	2.3	2.1
587745541052171128	143.8352	14.0716	24.0	0.7	25.2	0.5	25.1	0.7	22.3	0.2	21.0	0.2	2.7	1.3
587735239569244837	143.8536	35.9067	24.6	0.9	25.0	0.8	24.2	0.6	21.2	0.1	19.9	0.1	3.0	1.3
587735241717253077	144.2055	37.9672	24.1	0.7	25.3	0.5	24.9	0.6	22.7	0.3	21.1	0.2	2.2	1.5
587741532774990849	144.5733	23.6227	25.2	0.8	24.7	0.5	25.3	0.6	22.3	0.2	20.2	0.1	3.0	2.1
587725469600514982	144.7247	58.0650	23.3	0.5	24.1	0.4	24.7	0.7	22.1	0.2	20.6	0.2	2.6	1.5
588010135194960753	144.8458	54.8227	25.2	1.0	24.8	0.6	23.5	0.4	21.2	0.1	19.8	0.1	2.4	1.3
587746028521653096	145.0208	80.8068	24.9	1.7	25.6	0.9	25.1	1.0	22.0	0.2	20.3	0.2	3.1	1.7
587735347484492866	145.4473	11.0259	23.5	0.6	24.9	0.5	24.7	0.5	21.8	0.1	20.1	0.1	2.9	1.7
587745245236888254	145.6147	16.5869	23.9	0.9	24.3	0.4	24.3	0.6	20.9	0.1	19.6	0.1	3.3	1.4
587729149307519797	145.7727	-2.8949	24.1	1.1	24.4	0.8	24.3	0.9	21.7	0.2	20.0	0.1	2.6	1.7
588010135732225093	145.8303	55.6668	23.0	0.4	24.6	0.6	24.2	0.5	21.7	0.1	20.2	0.1	2.4	1.6
587745964634277002	146.0568	80.7485	25.6	1.1	24.3	0.5	25.3	0.9	22.2	0.2	20.1	0.1	3.1	2.1

Table C.1 (cont'd)

Object ID	RA	DEC	u	u err	g	g err	r	r err	i	i err	\mathbf{z}	z err	r-i	i-z
587742062130823973	146.4056	18.4595	25.1	0.8	25.1	0.7	24.0	0.4	21.6	0.1	20.2	0.1	2.4	1.4
587739114172974115	146.6519	30.3160	25.6	0.8	25.1	0.6	23.3	0.3	21.2	0.1	19.3	0.0	2.1	1.9
587745404151661507	147.1223	14.8640	25.0	0.9	25.4	0.5	24.0	0.5	21.4	0.1	19.9	0.1	2.6	1.5
587729387680760562	147.1262	54.4118	25.5	0.7	25.3	0.7	23.9	0.5	21.8	0.1	20.0	0.1	2.2	1.7
587725082503873175	147.1817	-2.6797	24.0	1.3	25.6	0.9	23.6	0.5	21.2	0.1	19.5	0.1	2.4	1.7
587739407304885266	147.3420	29.0520	25.6	0.7	24.4	0.5	24.9	0.7	21.9	0.1	20.4	0.2	3.0	1.5
588016892783755943	148.5550	37.3050	25.6	0.7	25.1	0.6	23.9	0.6	21.7	0.1	20.2	0.1	2.2	1.5
587739377778164705	148.6406	30.2270	25.8	0.6	24.3	0.5	23.5	0.4	21.4	0.1	20.0	0.1	2.1	1.4
587731500257903502	149.5388	52.1156	25.5	0.8	24.4	0.4	23.5	0.3	21.0	0.1	19.4	0.1	2.6	1.6
587732771041313734	149.6479	7.8360	25.0	1.0	24.3	0.5	23.4	0.3	21.0	0.1	19.7	0.1	2.4	1.3
587734622176346991	149.6934	41.2482	24.1	1.0	24.3	0.4	24.0	0.7	21.8	0.2	20.5	0.2	2.2	1.4
587738410862183251	150.0329	12.6704	23.7	0.9	23.8	0.3	23.9	0.5	21.3	0.1	20.0	0.1	2.5	1.4
587732049484907326	150.0537	52.0423	24.2	0.9	24.8	0.6	24.2	0.6	21.6	0.1	20.0	0.1	2.6	1.6
588848900436788207	150.8812	0.2572	25.5	0.6	24.2	0.4	23.9	0.4	21.5	0.1	20.2	0.1	2.4	1.3
587741491438879653	151.0424	26.5829	25.1	0.9	25.0	0.5	23.8	0.4	21.6	0.1	20.1	0.1	2.2	1.5
587731500795429936	151.2333	53.2276	25.6	0.7	25.4	0.5	23.8	0.4	21.7	0.1	19.9	0.1	2.2	1.7
587735661012780068	151.5337	36.5869	24.5	1.0	25.1	0.5	23.4	0.3	21.1	0.0	19.3	0.0	2.3	1.8
587746028522308335	151.8511	79.8267	25.6	1.2	26.5	0.5	24.1	0.8	21.0	0.1	19.5	0.1	3.0	1.6
587735661012911027	152.0080	36.6516	25.0	0.9	24.7	0.5	24.7	0.5	22.0	0.1	20.6	0.1	2.7	1.4
587741815709762781	152.7671	20.7417	24.9	0.9	24.3	0.3	23.2	0.2	20.8	0.0	19.5	0.1	2.4	1.4
587745401469863009	153.3206	13.9345	22.9	0.3	25.3	0.7	23.7	0.4	21.2	0.1	19.7	0.1	2.5	1.5
587733079193813860	153.6774	49.4685	22.4	0.3	25.2	0.9	24.0	0.7	21.7	0.1	20.3	0.1	2.4	1.4
587739376706323289	153.8388	30.7842	25.3	0.9	24.2	0.4	24.4	0.7	22.0	0.1	20.6	0.2	2.4	1.4
587738618093962166	154.5962	34.5953	24.9	1.2	24.1	0.4	24.8	0.7	22.1	0.2	20.3	0.1	2.7	1.7
587734623788991419	154.6451	44.3487	25.3	0.9	24.9	0.5	23.5	0.3	21.3	0.1	20.0	0.1	2.2	1.3
587728879803827157	155.0198	3.9675	24.5	0.9	24.5	0.6	24.6	0.6	22.3	0.2	20.8	0.2	2.3	1.5
588023046411125759	155.0974	23.9575	25.5	0.6	24.4	0.4	23.9	0.4	21.6	0.1	20.1	0.1	2.3	1.5

Table C.1 (cont'd)

Object ID	RA	DEC	u	u err	g	g err	r	r err	i	i err	\mathbf{Z}	$z \ err$	r-i	i-z
587735661014287541	155.7114	37.7758	23.9	0.8	25.0	0.5	24.3	0.6	21.5	0.1	19.7	0.1	2.8	1.8
587731500260066168	157.0133	54.2306	23.8	0.8	25.6	0.5	23.5	0.4	20.9	0.1	19.3	0.1	2.7	1.6
587732482741568799	157.5530	46.1316	24.7	0.9	24.8	0.4	23.5	0.3	21.1	0.1	19.5	0.1	2.4	1.6
587734863221883710	157.9937	9.2045	25.9	0.6	25.2	0.6	24.3	0.6	21.6	0.1	19.9	0.1	2.7	1.7
588017605750096661	159.5751	42.8911	24.8	0.8	25.1	0.5	23.9	0.4	21.4	0.1	19.9	0.1	2.4	1.5
587745539985179956	159.6697	15.7868	25.4	0.5	24.9	0.6	23.8	0.3	21.5	0.1	20.1	0.1	2.3	1.4
587732578309702669	159.8156	7.0035	24.8	1.0	24.9	0.5	24.2	0.5	21.6	0.1	20.3	0.1	2.6	1.4
587739407846802324	160.0626	32.8519	25.1	0.8	24.6	0.5	23.8	0.4	21.0	0.1	19.6	0.1	2.8	1.3
587732578309833768	160.1011	6.8967	24.3	0.9	24.5	0.5	23.8	0.4	21.4	0.1	19.7	0.1	2.3	1.7
587735348564788049	160.5537	13.7753	24.0	0.8	25.3	0.5	24.4	0.4	22.1	0.1	20.2	0.1	2.3	1.9
587742863935997065	160.6133	15.4069	26.4	0.5	24.0	0.4	23.6	0.4	21.0	0.1	19.6	0.1	2.6	1.4
588009370149061217	160.6377	58.9619	25.1	0.9	25.6	0.7	23.8	0.5	21.1	0.1	19.8	0.1	2.7	1.3
588017978885014291	161.5637	37.6308	24.9	1.1	24.7	0.6	24.7	0.7	21.4	0.1	20.0	0.1	3.3	1.4
587731498650567700	161.8517	54.0201	24.4	1.0	24.3	0.5	25.0	0.7	22.4	0.2	20.9	0.2	2.6	1.5
588848900978770884	162.5416	0.7967	23.9	0.6	25.3	0.5	24.1	0.6	21.7	0.1	20.1	0.1	2.4	1.6
588017111826367362	162.7343	45.7594	25.1	0.9	25.4	0.6	23.5	0.3	21.2	0.1	19.6	0.1	2.3	1.6
587738615412360104	162.7542	34.2673	24.3	0.8	24.3	0.5	24.1	0.6	21.8	0.1	20.2	0.1	2.4	1.6
587745517955974092	163.6358	-23.0362	23.9	1.6	25.6	0.7	24.5	0.8	22.1	0.2	20.2	0.1	2.5	1.8
587741531709309855	164.6244	27.4117	25.7	0.8	25.1	0.7	23.4	0.4	21.0	0.1	19.3	0.1	2.4	1.7
587726032241820432	164.6555	2.0157	23.1	0.6	24.3	0.5	24.7	0.7	22.5	0.3	20.4	0.2	2.2	2.1
588011098872087170	164.9986	61.8172	24.9	1.2	25.3	0.7	23.8	0.6	21.2	0.1	19.8	0.1	2.6	1.4
587739294550655835	165.0833	32.7535	24.4	1.2	25.5	0.5	24.8	0.6	22.1	0.2	20.7	0.2	2.7	1.4
588017110216475538	165.4542	45.0014	24.2	1.1	25.1	0.6	24.8	0.7	22.2	0.2	20.6	0.2	2.6	1.6
587745419172512941	165.5052	-20.0161	24.6	1.3	24.7	0.5	24.9	0.6	22.2	0.2	20.9	0.2	2.7	1.3
587742949567759253	165.6178	62.9000	24.9	1.3	25.5	0.6	24.1	0.7	21.8	0.1	20.1	0.1	2.4	1.7
588017720101962625	166.1909	40.1517	25.5	0.9	24.8	0.6	25.2	0.6	22.3	0.2	20.5	0.1	2.9	1.8
587742061065405423	166.4210	21.0518	25.3	0.7	25.1	0.6	23.9	0.4	21.7	0.1	20.3	0.1	2.2	1.5

Table C.1 (cont'd)

Object ID	RA	DEC	u	u err	g	g err	r	r err	i	i err	\mathbf{z}	z err	r-i	i-z
587745419173299268	166.4666	-18.4174	24.4	1.6	24.9	0.6	23.9	0.4	21.5	0.1	20.0	0.1	2.4	1.5
587739097519883343	166.5179	36.0247	25.0	1.0	25.1	0.6	23.5	0.9	20.8	0.0	19.4	0.1	2.7	1.4
587745517421659160	167.3495	-18.3025	25.4	1.1	25.2	0.9	23.9	0.5	21.7	0.1	20.4	0.2	2.2	1.4
588017702388368164	167.7542	10.9681	25.0	1.2	25.7	0.8	24.7	1.0	22.2	0.3	20.4	0.2	2.5	1.8
588017627224015847	168.5089	44.2847	25.1	0.7	24.5	0.5	24.0	0.5	21.8	0.1	20.1	0.1	2.2	1.7
587745421325763640	170.8485	-7.7415	24.5	0.9	24.5	0.4	23.9	0.4	21.7	0.1	20.0	0.1	2.2	1.8
588010358008710228	171.3568	3.8177	23.3	0.5	24.7	0.6	23.2	0.2	21.0	0.1	19.3	0.0	2.2	1.7
587726032781706246	171.5645	2.4977	24.1	1.2	23.5	0.3	23.9	0.5	21.7	0.1	20.2	0.2	2.2	1.5
588017110755509046	172.2477	46.4072	23.9	0.8	24.6	0.5	23.7	0.4	21.2	0.1	19.6	0.1	2.5	1.6
587748878767293348	172.4064	-4.6113	25.1	1.0	24.4	0.4	24.0	0.6	21.7	0.2	20.3	0.1	2.3	1.5
588010878755930812	172.6794	4.6231	25.1	1.0	23.6	0.4	24.4	0.8	21.0	0.1	19.6	0.1	3.4	1.5
587732772662215483	173.3376	10.7932	25.1	1.1	24.6	0.5	25.0	0.7	22.6	0.3	20.8	0.2	2.4	1.7
587732483283551295	174.1708	49.3591	24.4	0.7	24.4	0.3	24.4	0.4	22.0	0.1	20.3	0.1	2.4	1.7
587741709419676538	174.5612	27.4629	23.9	0.8	24.7	0.4	23.5	0.3	21.3	0.1	19.8	0.1	2.2	1.5
587735349376451646	175.4564	15.3323	24.5	0.8	24.1	0.3	24.3	0.5	21.8	0.1	19.9	0.1	2.6	1.9
587734894361904010	175.8093	10.9364	23.1	0.6	24.0	0.5	23.8	0.7	20.9	0.1	19.6	0.1	2.9	1.3
588023668631536741	175.8511	19.2520	25.4	1.0	25.0	0.6	24.3	0.4	22.0	0.1	20.7	0.1	2.3	1.3
587732771589850177	176.5640	9.9435	25.2	1.0	23.9	0.4	23.5	0.4	20.9	0.1	19.6	0.1	2.6	1.3
587731869096477517	177.0183	54.6310	25.6	0.8	25.0	0.6	24.6	0.7	21.7	0.1	20.3	0.1	2.9	1.4
587742572151440523	177.3201	19.6050	24.5	0.9	24.7	0.5	23.5	0.3	20.9	0.0	19.3	0.1	2.6	1.5
588013381668373442	177.6184	51.0420	25.4	0.8	25.0	0.7	24.2	0.5	21.1	0.1	19.4	0.1	3.1	1.7
587729386077750152	178.0388	59.0900	24.8	1.1	24.1	0.5	24.0	0.7	21.8	0.1	20.3	0.1	2.3	1.4
587741602029831217	179.2004	27.5228	25.5	0.8	26.2	0.3	24.4	0.6	22.3	0.2	20.7	0.2	2.2	1.5
588848900449436610	179.6998	0.2298	25.3	0.7	23.7	0.3	23.6	0.3	21.2	0.1	19.9	0.1	2.4	1.4
588017979428438734	180.6327	40.3433	24.9	1.0	24.3	0.6	24.7	0.7	22.1	0.2	20.8	0.2	2.6	1.3
587742189366609038	180.7011	24.2325	24.5	1.1	24.8	0.5	23.4	0.3	21.2	0.1	19.6	0.1	2.2	1.6
587728676856660277	180.7489	63.7893	22.3	0.2	24.9	0.6	25.5	0.5	23.2	0.9	21.7	0.4	2.2	1.5

Table C.1 (cont'd)

Object ID	RA	DEC	u	u err	g	g err	r	r err	i	i err	\mathbf{z}	z err	r-i	i-z
588017949350822800	181.0926	43.1136	26.1	0.9	24.3	0.3	23.7	0.4	21.5	0.1	19.7	0.1	2.1	1.8
587742863676605501	182.0668	16.8712	25.2	1.1	25.0	0.5	23.6	0.4	21.0	0.1	19.5	0.1	2.6	1.5
587732702875419575	182.1231	7.8887	25.6	1.0	24.5	0.5	24.0	0.5	21.3	0.1	19.9	0.1	2.8	1.4
587741602031862746	184.3958	27.6086	24.3	1.0	24.9	0.5	23.4	0.3	21.3	0.1	19.9	0.1	2.1	1.4
588013381670405094	185.0387	51.4017	24.9	0.9	25.2	0.6	23.8	0.4	20.8	0.1	19.0	0.0	3.0	1.9
587741601495319326	185.2366	27.2104	25.5	0.6	24.8	0.5	23.3	0.3	20.9	0.1	19.6	0.1	2.4	1.3
588017109685896049	185.9750	46.4380	25.3	0.9	25.7	0.6	25.2	0.6	21.8	0.1	19.9	0.1	3.3	1.9
587725039557936262	186.5731	-3.2575	25.7	0.8	25.6	0.5	24.0	0.5	21.8	0.1	20.4	0.2	2.1	1.4
587741602032976873	187.2962	27.6403	24.1	1.1	25.7	0.4	23.7	0.4	21.4	0.1	19.9	0.1	2.2	1.5
587725040632071125	187.5061	-2.3313	25.1	0.8	24.1	0.4	23.4	0.3	21.0	0.1	19.5	0.1	2.3	1.5
588017720645911348	187.5605	42.1361	24.8	1.1	25.2	0.5	24.3	0.6	21.8	0.1	20.3	0.1	2.5	1.5
587731869099492352	189.0408	54.8202	23.6	0.9	24.5	0.7	23.3	0.3	21.1	0.1	19.7	0.1	2.3	1.4
587732483287942155	189.8179	49.8006	23.7	0.6	25.4	0.5	24.4	0.6	21.2	0.1	19.7	0.1	3.2	1.6
587742904474141884	190.4263	18.6385	25.8	0.5	25.0	0.5	24.6	0.6	21.9	0.2	20.2	0.1	2.6	1.8
587732772132750451	190.6511	10.5235	26.0	0.9	25.0	0.6	25.3	0.6	22.3	0.2	21.0	0.2	3.0	1.4
588017991220789974	190.9165	9.9436	25.5	0.9	24.1	0.4	24.1	0.6	21.7	0.1	20.3	0.1	2.4	1.4
587738575141209378	190.9666	40.9319	23.6	0.5	25.1	0.4	25.0	0.4	22.4	0.2	20.9	0.2	2.6	1.4
587745543729120325	191.7133	-13.2505	23.8	1.3	25.7	0.6	24.5	0.7	22.1	0.2	20.6	0.2	2.4	1.5
587746041410552832	192.0253	-8.3185	24.0	1.0	25.4	0.7	24.9	0.7	22.3	0.2	20.9	0.2	2.6	1.4
587742902327510145	192.3695	16.7756	25.2	0.8	24.5	0.4	24.9	0.5	22.4	0.2	20.9	0.2	2.6	1.5
587742575908619510	192.6829	19.1709	24.9	1.1	23.7	0.3	23.5	0.3	20.9	0.0	19.5	0.1	2.6	1.4
587732771596993679	193.1921	9.9646	25.9	0.6	24.9	0.6	23.3	0.3	21.1	0.1	19.7	0.1	2.2	1.4
588023668639007562	193.9799	19.3382	24.7	1.3	24.3	0.4	24.4	0.5	21.3	0.1	19.9	0.1	3.0	1.5
587741724971762568	194.2034	23.7828	24.5	0.8	25.1	0.6	24.9	0.6	21.7	0.1	20.2	0.1	3.3	1.5
587739720291189639	195.6271	29.1829	24.9	1.1	25.2	0.6	24.4	0.6	21.4	0.1	19.4	0.1	3.0	2.0
587733195697226654	195.8201	54.0262	25.5	0.9	25.1	0.6	24.8	0.7	22.1	0.2	20.3	0.1	2.7	1.8
587728678470157331	196.2246	64.6864	23.7	0.8	24.3	0.5	24.3	0.6	21.6	0.1	20.1	0.1	2.8	1.5

Table C.1 (cont'd)

Object ID	RA	DEC	u	u err	g	g err	r	r err	i	i err	\mathbf{z}	z err	r-i	i-z
587742062151271483	196.3262	22.5705	24.6	1.1	25.4	0.6	25.2	0.5	21.8	0.1	20.5	0.1	3.3	1.4
588017722260062969	198.7887	42.6726	25.4	1.2	24.8	0.6	24.1	0.6	21.7	0.1	20.3	0.1	2.4	1.4
588018055645299778	199.3085	51.5645	24.6	0.8	24.3	0.3	24.8	0.5	21.8	0.1	20.4	0.1	3.0	1.4
587736807225361346	201.7009	13.2873	24.3	1.0	25.0	0.7	24.2	0.5	21.9	0.1	20.6	0.2	2.2	1.3
587725816953439053	201.7913	66.0372	23.6	0.8	24.1	0.3	24.2	0.7	21.1	0.1	19.3	0.1	3.2	1.7
587729773680264258	202.0138	-2.3744	24.3	1.2	25.3	0.7	23.9	0.6	21.7	0.1	20.1	0.2	2.3	1.5
588017991225574380	202.0260	9.4413	23.7	0.9	24.5	0.5	24.9	0.7	22.5	0.2	20.9	0.2	2.4	1.6
587739504476357411	202.0431	30.4375	25.3	0.9	24.9	0.8	24.8	0.7	21.9	0.1	20.2	0.1	2.9	1.7
587729776902079384	202.2706	-2.9269	25.7	1.0	23.6	0.4	24.2	0.7	21.9	0.1	20.5	0.2	2.3	1.4
587741721216156530	202.4198	25.8089	25.3	0.9	25.2	0.6	23.6	0.4	21.3	0.1	19.4	0.1	2.3	1.9
587726014005445549	203.3734	1.3577	25.2	1.0	24.5	0.7	23.4	0.3	21.1	0.1	19.6	0.1	2.3	1.5
588017726019208408	203.7106	7.6235	25.4	0.7	25.3	0.5	25.0	0.6	22.5	0.2	20.7	0.1	2.4	1.8
588017570317075482	204.1004	11.4837	25.1	0.8	25.1	0.6	24.6	0.6	22.4	0.2	21.0	0.2	2.2	1.4
587739132951987468	204.4014	36.6419	25.2	0.6	25.0	0.4	24.1	0.4	21.9	0.1	20.3	0.1	2.2	1.6
587729773681574984	204.9153	-2.1625	23.8	0.9	25.5	0.6	24.2	0.5	22.0	0.2	20.6	0.2	2.2	1.4
587738574072317507	205.2519	38.5611	24.9	0.7	24.9	0.4	24.4	0.4	21.9	0.1	20.6	0.1	2.5	1.3
588011219135431601	205.9093	61.4749	25.2	0.7	24.0	0.3	23.6	0.3	21.3	0.1	19.8	0.1	2.4	1.5
587739707943027687	206.4155	28.6515	25.5	0.9	25.1	0.5	24.0	0.5	21.5	0.1	19.9	0.1	2.5	1.6
587742060544918633	206.5629	20.3468	25.9	0.9	24.7	0.6	24.8	0.6	22.5	0.2	21.1	0.2	2.3	1.4
587738574609647082	206.5928	38.7391	24.7	0.6	25.2	0.3	25.3	0.4	23.0	0.3	21.2	0.2	2.4	1.7
587722981751915681	206.9856	-1.0598	24.3	0.8	25.1	0.6	24.0	0.4	21.8	0.1	20.0	0.1	2.2	1.8
587746269582001411	207.0844	-3.5539	23.6	1.0	24.4	0.5	24.5	0.6	21.4	0.1	20.0	0.1	3.1	1.4
587735696449209285	207.2544	55.4942	24.2	1.0	25.4	0.5	24.6	0.6	22.2	0.2	20.5	0.2	2.4	1.7
587729774756561829	207.8376	-1.3271	25.6	0.8	24.4	0.6	23.9	0.5	21.4	0.1	20.0	0.1	2.5	1.4
588017625089115043	207.9801	41.4142	25.0	0.9	24.9	0.7	24.3	0.5	21.6	0.1	20.2	0.1	2.7	1.4
587746236297905182	209.0299	-6.4122	25.7	0.9	25.6	0.5	24.7	0.7	21.8	0.1	19.9	0.1	2.9	1.9
587739406254277723	209.2927	31.8219	23.7	0.8	25.2	0.4	24.3	0.5	21.9	0.1	20.4	0.1	2.4	1.5

Table C.1 (cont'd)

Object ID	RA	DEC	u	u err	g	$g \ err$	r	r err	i	i err	\mathbf{z}	$z \ err$	r-i	i-z
587729159509705539	209.4048	4.8979	25.7	0.9	24.5	0.6	23.6	0.4	21.5	0.1	20.2	0.2	2.1	1.3
587736586573120665	209.4356	40.8209	24.0	0.7	24.7	0.4	24.0	0.4	21.6	0.1	19.7	0.1	2.4	1.9
588298664117994712	209.4619	45.5971	24.4	0.6	25.1	0.4	23.8	0.3	21.0	0.1	19.4	0.0	2.8	1.6
587742594695693110	209.6562	16.8099	23.3	0.5	24.4	0.5	23.6	0.4	21.1	0.1	19.7	0.1	2.5	1.4
587742592548537313	210.3118	15.0032	25.4	0.9	25.1	0.8	24.3	0.6	21.1	0.1	19.6	0.1	3.3	1.4
588017991229244610	210.5254	8.9233	23.7	0.8	24.9	0.6	25.4	0.6	22.3	0.2	20.9	0.2	3.1	1.4
588017991229310016	210.6635	8.9254	25.5	0.9	25.2	0.6	23.4	0.3	21.0	0.1	19.6	0.1	2.4	1.3
587735695913321470	210.8686	54.3749	26.0	0.6	27.4	0.2	24.0	0.7	21.8	0.1	20.5	0.2	2.1	1.4
588017992303838071	212.4217	9.5701	25.1	0.8	25.1	0.6	24.2	0.6	21.4	0.1	19.9	0.1	2.8	1.5
587736478662656905	212.5009	11.6189	25.0	0.8	24.3	0.5	24.5	0.6	22.3	0.2	20.7	0.2	2.2	1.6
588011219137004633	212.6382	60.1345	24.3	0.8	24.9	0.5	24.6	0.5	22.2	0.2	20.8	0.2	2.4	1.4
587739380985562078	213.0207	31.6272	26.1	0.6	25.3	0.5	25.0	0.6	22.0	0.2	20.3	0.1	3.0	1.7
587746278167282738	213.5916	-18.0786	23.1	0.9	25.6	0.7	23.9	0.7	21.7	0.1	20.1	0.1	2.2	1.7
587739609175163770	213.6255	29.3398	24.7	1.2	25.2	0.5	24.3	0.6	21.9	0.1	20.1	0.1	2.4	1.9
588848899927442575	213.6762	-0.1211	25.5	0.6	25.8	0.5	23.8	0.5	21.5	0.1	20.1	0.1	2.3	1.4
587736542017750200	213.9824	8.3750	24.6	0.9	26.3	0.3	23.8	0.4	21.4	0.1	19.8	0.1	2.4	1.7
588011123583091390	214.1712	59.0976	25.4	0.8	24.4	0.6	24.4	0.8	21.9	0.2	20.6	0.2	2.5	1.3
587746278167938208	214.2849	-19.4582	23.9	1.8	26.3	0.5	25.4	0.8	22.0	0.2	20.7	0.2	3.4	1.4
587736585501410737	214.8308	38.5935	23.8	0.6	24.5	0.3	24.1	0.3	21.8	0.1	20.0	0.1	2.3	1.8
587742190454244421	215.2578	22.1226	25.3	1.0	25.2	0.5	25.0	0.5	22.3	0.2	20.6	0.2	2.6	1.7
587736914602951670	215.2992	12.1341	23.2	0.5	25.5	0.6	23.3	0.3	21.0	0.1	19.7	0.1	2.3	1.3
587736584428062095	215.5006	37.3992	24.1	0.6	24.1	0.3	23.6	0.3	20.9	0.0	19.5	0.1	2.7	1.4
587730022789809162	215.5008	6.5153	24.9	1.3	25.5	0.6	23.9	0.5	21.7	0.1	20.2	0.2	2.2	1.5
588017606304989959	216.2809	41.6769	25.8	0.8	24.5	0.5	23.5	0.3	21.0	0.1	19.6	0.1	2.5	1.4
587736541481927841	216.3328	7.8922	23.4	0.7	24.5	0.4	23.4	0.3	21.0	0.1	19.7	0.1	2.4	1.3
588017713661412643	216.6573	49.0901	25.6	0.7	25.1	0.4	23.7	0.3	20.8	0.1	19.1	0.0	2.8	1.7
587733410984231821	218.0394	49.9561	23.6	0.7	25.2	0.6	23.8	0.5	20.8	0.1	19.1	0.0	2.9	1.7

Table C.1 (cont'd)

Object ID	RA	DEC	u	u err	g	g err	r	r err	i	i err	\mathbf{z}	z err	r-i	i-z
587733397562262461	218.1495	49.2996	24.0	1.2	24.1	0.4	24.2	0.7	21.5	0.1	20.0	0.1	2.8	1.5
587739827129484142	218.6243	22.7365	25.4	1.2	25.6	0.5	24.1	0.5	21.3	0.1	20.0	0.1	2.7	1.3
588298662510789880	218.8043	41.5426	23.8	0.6	25.0	0.3	24.6	0.4	22.5	0.2	20.4	0.1	2.1	2.1
587736962922316671	219.5107	31.8622	24.7	0.9	24.9	0.5	24.5	0.6	22.3	0.2	21.0	0.2	2.2	1.3
588017711515370532	220.0139	46.1570	24.8	1.0	25.3	0.5	24.2	0.4	22.0	0.1	20.3	0.1	2.2	1.6
587739630630863931	220.8132	27.5135	25.3	0.9	24.7	0.5	23.7	0.4	21.5	0.1	20.1	0.1	2.3	1.3
587742629061329898	221.6043	14.8235	25.1	0.9	24.2	0.4	24.3	0.6	22.1	0.2	20.6	0.1	2.1	1.5
587722982295340447	221.9858	-0.7504	24.9	0.9	25.2	0.5	24.5	0.5	21.4	0.1	19.7	0.1	3.1	1.7
587726032267052136	222.3068	1.6773	24.4	1.2	24.4	0.6	23.4	0.4	21.2	0.1	19.9	0.1	2.2	1.4
588018055115703409	222.6902	45.3658	25.0	1.0	24.2	0.4	24.6	0.5	22.3	0.2	20.9	0.2	2.3	1.4
587729776911123462	223.0024	-2.4566	24.6	1.5	25.7	0.6	24.0	0.5	21.9	0.1	20.4	0.2	2.2	1.5
587735490821751827	223.3052	43.5829	23.6	0.5	24.4	0.4	23.3	0.2	20.9	0.1	19.5	0.1	2.3	1.4
587736585504622197	223.3183	35.7665	23.7	0.5	25.2	0.3	23.8	0.3	21.1	0.1	19.2	0.0	2.6	2.0
587729776911385690	223.6075	-2.4566	26.0	0.8	25.2	0.7	24.0	0.5	21.8	0.1	19.9	0.1	2.2	1.8
587739827131974679	224.4462	21.0559	25.6	1.0	24.7	0.5	24.3	0.5	21.9	0.1	20.4	0.1	2.4	1.5
588017979980907639	224.8475	33.7420	24.5	0.9	25.0	0.4	24.3	0.5	22.0	0.1	20.4	0.1	2.3	1.6
588017702949880900	224.9550	9.1385	24.8	1.1	24.9	0.5	24.4	0.5	21.8	0.1	20.4	0.1	2.6	1.3
588848899932554225	225.4619	-0.0569	25.3	0.7	24.9	0.4	24.9	0.6	22.3	0.2	20.8	0.2	2.6	1.5
587739131349632328	225.6101	29.7374	25.9	0.5	24.8	0.4	23.9	0.3	20.9	0.0	19.3	0.0	3.0	1.6
587739382064350410	225.8459	28.7320	23.8	0.7	25.0	0.4	24.9	0.5	21.8	0.1	20.1	0.1	3.1	1.7
587739707950957500	225.9034	23.8605	24.9	1.1	25.3	0.5	24.6	0.6	22.2	0.2	20.5	0.2	2.4	1.7
587742061090112393	226.3155	16.9165	24.6	1.0	25.0	0.5	23.9	0.4	21.8	0.1	20.5	0.1	2.1	1.3
587742782061151186	226.3166	66.5927	24.7	1.5	25.3	0.6	25.0	0.7	22.1	0.2	20.7	0.2	2.8	1.4
587739810492187739	226.4236	21.1685	25.3	0.7	23.9	0.3	24.2	0.4	21.5	0.1	20.1	0.1	2.6	1.4
587729228758057768	226.4657	61.8045	25.3	0.9	24.8	0.6	24.7	0.7	22.1	0.2	20.5	0.1	2.5	1.6
587736975270610131	226.9655	28.8498	25.3	0.9	24.8	0.5	23.3	0.2	21.1	0.1	19.8	0.1	2.2	1.3
588017605772248021	227.2607	37.3744	24.1	0.9	25.9	0.4	24.6	0.6	22.2	0.2	20.9	0.2	2.4	1.3

Table C.1 (cont'd)

Object ID	RA	DEC	u	u err	g	g err	r	r err	i	i err	\mathbf{Z}	z err	r-i	i-z
587726100415775772	227.5137	2.9317	24.5	1.7	24.9	0.8	23.7	0.5	21.2	0.1	19.8	0.1	2.5	1.4
587742610269472024	227.7051	12.5743	25.2	0.9	24.3	0.4	23.5	0.3	21.1	0.1	19.4	0.1	2.4	1.8
587733412061119193	228.1641	47.1411	24.4	1.1	25.5	0.6	24.7	0.7	21.8	0.2	20.4	0.1	2.9	1.4
588017990700303528	228.5661	6.7330	23.7	0.9	24.6	0.6	24.6	0.6	21.9	0.1	20.4	0.1	2.8	1.4
587739720304493641	228.7774	22.4037	25.5	0.8	24.6	0.5	23.5	0.4	21.3	0.1	19.9	0.1	2.3	1.4
588017626707199080	228.8637	35.9400	25.0	0.9	25.1	0.4	23.8	0.3	21.6	0.1	19.9	0.1	2.2	1.7
587729227148624895	229.2771	59.1033	25.3	1.1	24.3	0.5	24.4	0.7	21.4	0.1	20.0	0.1	2.9	1.4
587735667458310997	229.2816	49.6756	25.8	0.8	25.2	0.6	23.5	0.3	21.2	0.1	19.8	0.1	2.3	1.4
587739720304886838	229.6108	21.9413	23.9	0.9	25.1	0.5	24.1	0.4	21.3	0.1	19.7	0.1	2.9	1.5
587729746838029514	230.4316	1.0337	24.6	1.4	24.6	0.6	23.3	0.3	21.0	0.1	19.6	0.1	2.3	1.3
587736478670521399	230.5734	9.3747	24.6	0.7	24.9	0.5	25.0	0.6	22.1	0.2	20.6	0.1	2.9	1.5
587733411525559181	231.3260	45.0889	26.0	0.7	24.6	0.5	24.9	0.6	22.1	0.1	20.7	0.2	2.8	1.4
587733427626967884	231.5140	49.0118	23.8	1.0	24.8	0.6	23.6	0.3	21.5	0.1	20.0	0.1	2.1	1.5
587739811568026947	231.6018	20.3935	25.2	0.6	25.2	0.4	24.5	0.4	22.1	0.1	20.8	0.2	2.4	1.3
587739845389386773	231.6261	18.0976	25.1	0.9	24.1	0.3	24.3	0.4	21.5	0.1	19.9	0.1	2.7	1.6
587742629065851966	231.9400	12.6672	24.9	1.0	25.2	0.5	25.3	0.6	22.3	0.2	20.8	0.2	3.0	1.5
587729159519601807	232.0062	3.6837	24.1	1.0	25.1	0.7	24.4	0.6	22.0	0.2	20.5	0.2	2.4	1.5
587742628529374488	232.8739	12.0835	22.7	0.4	25.2	0.5	24.9	0.6	22.0	0.1	20.5	0.2	2.9	1.5
588017703490356395	233.2395	8.2607	23.8	0.9	25.3	0.6	24.2	0.5	21.9	0.1	20.5	0.1	2.3	1.4
587735489215398938	233.4104	37.4698	24.6	1.2	24.8	0.5	25.1	0.7	22.0	0.2	20.7	0.2	3.1	1.3
587736584972141961	233.8093	30.7206	24.8	0.7	24.9	0.4	25.0	0.4	22.4	0.1	21.0	0.2	2.6	1.4
588011217533076244	233.9039	50.5680	26.2	0.5	25.3	0.6	24.4	0.6	21.8	0.1	20.4	0.1	2.6	1.3
587739131353236873	233.9560	26.2114	25.1	0.9	24.9	0.5	23.3	0.2	21.0	0.0	19.7	0.1	2.3	1.3
587729229296894714	234.0202	59.1885	23.7	1.0	24.9	0.7	24.4	0.7	21.2	0.1	19.9	0.1	3.2	1.3
587736542026532135	234.0643	6.1930	24.9	0.9	25.3	0.5	23.3	0.3	21.1	0.1	19.6	0.1	2.2	1.5
588017991239533745	234.0934	6.3800	24.9	1.4	24.9	0.7	24.2	0.7	21.9	0.1	20.5	0.2	2.3	1.4
588011219679904764	234.0954	52.7817	22.6	0.3	24.4	0.4	24.3	0.5	21.9	0.1	20.5	0.2	2.4	1.4

Table C.1 (cont'd)

Object ID	RA	DEC	u	u err	g	g err	r	r err	i	i err	\mathbf{z}	$z \ err$	r-i	i-z
587739849672885314	234.2626	61.2775	25.3	1.0	24.7	0.5	25.0	0.6	22.3	0.2	21.0	0.3	2.7	1.3
587729407535023060	234.7772	53.2424	24.0	1.3	25.2	0.6	24.1	0.6	21.4	0.1	19.8	0.1	2.7	1.6
587733604262217093	234.9250	48.0162	25.0	0.8	25.4	0.4	25.3	0.4	23.0	0.2	20.9	0.2	2.4	2.0
587729748450608329	234.9880	2.0907	24.5	1.0	24.6	0.5	23.5	0.4	20.9	0.1	19.4	0.1	2.6	1.5
588018090541581454	235.0644	35.0280	24.2	1.0	25.1	0.5	23.7	0.4	21.0	0.1	19.6	0.1	2.7	1.4
587736543100798537	235.4796	6.8844	23.6	0.5	25.0	0.5	23.7	0.3	21.4	0.1	20.0	0.1	2.2	1.4
587736752465708166	235.5902	36.6984	25.1	1.0	24.9	0.6	24.6	0.6	22.2	0.2	20.9	0.2	2.4	1.3
588011217533731867	235.6902	49.4629	24.6	1.3	24.6	0.5	24.7	0.7	22.3	0.2	20.5	0.2	2.3	1.8
587729158984303831	235.6940	3.1472	23.9	0.9	24.5	0.7	23.6	0.4	21.2	0.1	19.8	0.1	2.4	1.4
587739815314785389	235.9708	22.2438	26.0	0.6	25.1	0.5	24.0	0.4	21.3	0.1	19.9	0.1	2.8	1.4
588848900474078591	236.0793	0.2772	23.9	0.7	25.3	0.5	24.0	0.5	21.2	0.1	19.6	0.1	2.7	1.6
588017991777322253	236.1821	6.6733	24.7	1.4	24.2	0.5	24.1	0.6	21.9	0.1	20.5	0.2	2.2	1.4
587742611346949395	236.4521	11.5199	25.2	1.0	25.4	0.5	23.4	0.3	21.1	0.1	19.4	0.1	2.3	1.6
587736478673405157	237.1816	8.2832	25.9	0.4	24.7	0.5	23.9	0.4	21.6	0.1	20.3	0.1	2.3	1.4
588017992314717506	237.4281	6.8918	23.6	0.8	25.6	0.6	24.4	0.6	21.9	0.1	20.3	0.1	2.4	1.7
587742782064231497	237.8862	61.7237	25.8	0.9	24.4	0.5	24.0	0.5	21.8	0.1	20.5	0.2	2.2	1.4
587736584437237083	237.8977	27.9870	24.7	1.0	25.3	0.4	23.7	0.3	21.3	0.1	19.4	0.0	2.4	1.9
587739845392270485	238.1536	15.9378	24.3	1.0	25.8	0.4	23.2	0.2	21.1	0.1	19.6	0.1	2.2	1.4
587736751930147927	238.1548	34.5144	25.7	0.8	25.2	0.6	24.8	0.6	22.1	0.1	20.6	0.1	2.7	1.5
587736478137124342	238.4324	7.6701	24.8	0.9	25.2	0.5	24.7	0.6	21.9	0.1	20.3	0.1	2.8	1.6
587733609086452775	238.5687	46.0634	26.0	0.8	25.0	0.7	25.1	0.7	22.3	0.2	21.0	0.3	2.7	1.4
587739652644734147	238.8518	21.2850	24.3	0.9	25.2	0.5	24.2	0.5	21.9	0.2	20.6	0.2	2.2	1.4
587736478137386319	238.9851	7.4959	24.3	0.9	25.0	0.5	23.5	0.3	21.0	0.1	19.2	0.1	2.5	1.8
588018253757678463	239.0767	37.5630	25.3	1.1	25.2	0.7	23.8	0.5	21.3	0.1	19.8	0.1	2.5	1.4
587739651034514532	239.1066	19.7811	23.7	1.0	25.2	0.6	23.3	0.3	21.0	0.1	19.3	0.1	2.3	1.7
587735665851958247	239.1080	42.5534	23.6	0.7	25.2	0.7	24.5	0.7	22.2	0.2	20.7	0.2	2.3	1.5
587736915687245537	239.6342	8.7628	25.2	0.6	24.7	0.4	23.3	0.2	21.1	0.1	19.7	0.1	2.2	1.4

Table C.1 (cont'd)

Object ID	RA	DEC	u	u err	g	g err	r	r err	i	i err	\mathbf{Z}	z err	r-i	i-z
587739720309343534	239.6798	18.1405	25.0	1.1	25.6	0.5	24.1	0.5	21.8	0.1	20.5	0.2	2.3	1.4
587736921042388393	240.0317	30.7710	25.6	0.5	24.7	0.4	25.1	0.5	22.5	0.2	20.8	0.1	2.6	1.7
587736545780958737	240.4654	3.7250	24.6	1.0	25.0	0.5	23.6	0.3	21.0	0.1	19.7	0.1	2.5	1.4
588011218609374239	241.1320	47.0361	23.6	0.6	23.9	0.3	25.1	0.6	21.9	0.1	20.5	0.2	3.2	1.4
587736619325261027	241.3784	26.7038	25.0	0.7	25.3	0.5	23.7	0.4	21.4	0.1	19.8	0.1	2.3	1.6
588018056733459669	241.4014	36.8108	24.2	0.8	24.3	0.4	24.1	0.4	21.9	0.1	20.4	0.1	2.2	1.5
587739382071231978	241.6125	21.9298	25.2	0.8	24.3	0.4	25.1	0.5	22.9	0.3	21.1	0.2	2.3	1.8
587736543103616588	241.8496	6.0972	24.7	0.8	24.6	0.5	24.3	0.4	22.0	0.1	20.6	0.2	2.2	1.4
587742610275763654	241.9947	9.2120	25.9	0.7	25.6	0.5	25.0	0.7	22.6	0.2	20.6	0.2	2.4	2.0
587742615099475377	242.0435	13.8276	25.8	0.6	25.3	0.6	23.7	0.4	21.3	0.1	19.8	0.1	2.4	1.5
587736478675633759	242.1584	7.3265	24.6	0.8	25.4	0.5	24.6	0.6	22.4	0.2	21.1	0.2	2.2	1.3
588018055660438849	242.1775	35.0422	25.1	0.9	25.4	0.5	24.3	0.5	22.1	0.1	20.5	0.1	2.3	1.6
587745969465197834	242.2062	18.6575	25.7	0.6	25.4	0.4	24.3	0.6	21.6	0.1	20.0	0.1	2.7	1.6
587735743153964425	242.3917	30.4111	24.6	0.8	25.1	0.4	24.9	0.5	22.2	0.1	20.8	0.2	2.7	1.4
587742062170867072	242.4613	13.0901	24.2	0.9	24.7	0.5	24.7	0.6	21.8	0.1	20.3	0.1	2.9	1.5
588017991780140686	242.5729	5.7330	25.3	1.1	25.5	0.6	25.1	0.6	22.7	0.3	20.7	0.2	2.4	2.0
587733398108570644	242.7150	36.9276	23.4	0.7	25.6	0.6	24.7	0.8	21.8	0.1	20.3	0.2	2.9	1.5
587733411530540062	242.8266	37.5447	25.2	1.1	25.7	0.6	24.7	0.7	21.6	0.1	20.3	0.2	3.1	1.3
587736975277753681	243.0262	21.3412	25.0	0.9	25.3	0.5	25.0	0.4	22.8	0.2	20.9	0.2	2.3	1.9
587735665853990030	243.3169	39.3194	24.5	0.9	24.4	0.4	23.4	0.3	21.2	0.1	19.9	0.1	2.1	1.4
588017704031618682	243.3674	6.9732	23.6	0.9	24.4	0.4	24.9	0.7	21.8	0.1	19.9	0.1	3.1	1.9
588017990706922819	243.6486	4.6922	25.3	1.2	25.3	0.6	24.3	0.7	22.0	0.2	20.7	0.2	2.3	1.4
587733410457519145	243.7203	35.9499	25.3	1.2	24.7	0.6	25.1	0.6	21.9	0.1	20.5	0.1	3.2	1.4
587742610276550453	243.7307	8.7167	24.5	1.2	25.5	0.5	23.9	0.5	21.3	0.1	19.7	0.1	2.6	1.6
587729227154129885	243.9273	49.8892	26.2	1.0	26.1	0.5	24.0	0.7	21.8	0.1	19.8	0.1	2.2	2.0
588017990707054071	244.0028	4.6171	24.9	1.5	23.8	0.3	24.7	0.7	22.1	0.1	20.5	0.1	2.6	1.6
587733603729540496	244.0163	40.8852	24.5	1.1	24.4	0.4	24.2	0.3	21.8	0.1	20.2	0.1	2.4	1.5

Table C.1 (cont'd)

Object ID	RA	DEC	u	u err	g	g err	r	r err	i	i err	\mathbf{z}	z err	r-i	i-z
588017978915816840	244.0491	23.3466	24.1	0.7	25.2	0.4	23.7	0.8	21.1	0.1	19.6	0.1	2.6	1.5
587742645701707491	244.1075	12.2334	25.1	0.9	25.1	0.5	24.0	0.4	21.6	0.1	20.2	0.1	2.3	1.4
587739384742544637	244.2294	19.1452	23.6	0.9	25.0	0.6	23.6	0.4	21.5	0.1	19.5	0.1	2.1	2.0
587742628534551977	244.5379	9.2653	24.6	1.6	24.9	0.7	24.1	0.6	21.2	0.1	19.8	0.1	2.9	1.5
587733604266476986	244.6289	41.1774	24.8	0.8	24.8	0.4	23.3	0.2	21.1	0.1	19.7	0.1	2.2	1.4
587736542031251048	244.7886	4.7700	26.1	0.4	24.5	0.4	24.3	0.5	22.1	0.1	20.6	0.2	2.2	1.5
587739815318848801	244.9110	18.0632	26.0	0.6	24.9	0.5	24.0	0.4	21.7	0.1	19.8	0.1	2.3	1.9
587736477603203007	245.1129	6.0467	25.1	0.8	24.8	0.5	25.1	0.5	22.0	0.1	20.2	0.1	3.1	1.8
587742550692071086	245.3298	10.1221	24.0	0.8	25.2	0.4	24.2	0.4	22.0	0.1	20.4	0.1	2.2	1.6
587739720848770485	245.3571	15.9546	24.6	1.3	25.0	0.6	23.6	0.4	21.5	0.1	20.0	0.1	2.2	1.4
588017978916472499	245.5021	22.5628	23.6	0.5	24.7	0.3	24.0	0.3	21.8	0.1	20.4	0.1	2.2	1.3
587733411531981933	245.7252	35.3316	23.1	0.6	24.8	0.6	24.8	0.6	21.9	0.1	20.5	0.2	2.8	1.4
587733441588954429	246.0699	37.4823	25.0	0.8	25.1	0.4	23.3	0.2	20.8	0.0	19.4	0.1	2.5	1.5
587739827678545312	246.3828	13.6807	23.1	0.6	24.9	0.6	24.5	0.7	21.5	0.1	20.1	0.1	3.0	1.4
587736981712339898	246.5739	54.2199	22.8	0.6	25.1	0.7	24.7	0.7	22.4	0.3	20.5	0.2	2.3	1.8
587729652884440235	246.7736	42.0935	25.3	0.9	24.4	0.5	24.4	0.6	21.4	0.1	20.0	0.1	3.0	1.4
587729653957723430	246.8709	43.4958	24.6	1.0	24.8	0.6	23.5	0.3	21.3	0.1	20.0	0.1	2.1	1.3
587729751668950320	247.0348	41.4821	23.3	0.5	25.1	0.5	24.0	0.4	21.3	0.1	19.9	0.1	2.7	1.4
587739166239753205	247.2699	19.6234	24.7	1.5	25.1	0.8	24.4	0.9	21.8	0.2	20.4	0.2	2.6	1.4
587736813139854969	247.5860	7.0929	24.3	1.3	25.4	0.6	24.8	0.6	22.1	0.1	20.6	0.2	2.8	1.5
587729751669278024	247.7063	40.8164	24.6	0.9	24.8	0.5	23.4	0.3	21.0	0.1	19.7	0.1	2.3	1.3
587735743156717021	247.9481	26.3187	23.8	0.6	25.2	0.5	24.8	0.5	22.4	0.2	21.0	0.2	2.4	1.4
587736919436166489	248.1066	23.6474	25.3	1.1	24.7	0.5	24.5	0.5	22.3	0.2	20.6	0.1	2.2	1.7
587733440516916637	248.3485	33.9695	24.6	0.9	25.3	0.4	24.8	0.5	22.5	0.2	20.9	0.2	2.3	1.7
587736751935325329	248.6440	26.9571	24.6	1.3	25.5	0.6	24.4	0.6	22.1	0.2	20.6	0.2	2.3	1.5
587736782000162004	248.9315	26.9147	23.7	1.0	25.5	0.7	23.8	0.5	21.0	0.1	19.1	0.1	2.8	1.9
758879715545319438	248.9442	-5.1978	23.5	0.8	23.8	0.4	23.3	0.3	20.9	0.1	19.3	0.1	2.4	1.6

Table C.1 (cont'd)

Object ID	RA	DEC	u	u err	g	g err	r	r err	i	i err	\mathbf{Z}	z err	r-i	i-z
587733441590724022	249.3206	34.4546	24.8	0.8	24.9	0.4	23.5	0.3	21.0	0.0	19.1	0.0	2.6	1.9
587729652885816245	249.4168	39.6499	24.7	1.0	25.1	0.5	25.2	0.6	22.1	0.2	20.8	0.2	3.2	1.3
587742782068950343	249.4335	52.7792	25.3	1.0	24.4	0.4	24.3	0.5	21.3	0.1	19.9	0.1	3.0	1.4
587736751935718633	249.4523	26.2625	24.3	1.2	24.5	0.5	25.0	0.6	22.6	0.2	20.8	0.2	2.4	1.8
758879663465040338	249.4910	-6.4353	24.5	1.3	24.5	0.5	24.7	0.6	21.7	0.1	20.1	0.1	3.0	1.5
587736946812651159	249.5874	23.8315	25.3	0.6	25.7	0.4	25.1	0.5	22.9	0.3	21.0	0.2	2.2	1.9
587733431921345875	249.6904	30.5633	25.4	0.9	24.5	0.4	24.0	0.5	21.7	0.1	20.3	0.1	2.2	1.4
587739845397514242	249.7283	11.5006	25.3	0.8	24.6	0.5	25.1	0.5	22.0	0.1	20.4	0.2	3.1	1.6
587739707961902675	250.2695	13.8930	23.9	1.2	25.0	0.7	24.0	0.5	21.7	0.1	20.2	0.2	2.2	1.5
758879662928890205	250.6202	-7.6837	24.6	1.3	24.5	0.5	24.9	0.7	22.4	0.2	20.8	0.2	2.5	1.6
587736976891774814	250.6440	18.2724	25.0	0.7	24.5	0.4	23.5	0.2	21.1	0.1	19.4	0.0	2.4	1.7
587729231979742631	250.7657	44.0368	25.2	0.7	25.2	0.5	24.8	0.6	22.1	0.1	20.3	0.1	2.7	1.8
588017978382681775	251.3875	18.0645	24.6	0.7	25.0	0.4	24.7	0.5	22.2	0.2	20.7	0.2	2.5	1.5
758879663466088597	251.4277	-7.8792	24.8	1.3	25.1	0.6	25.6	0.5	22.3	0.2	20.4	0.2	3.2	1.9
588018056202028389	251.5395	27.8247	25.5	0.8	25.5	0.4	24.9	0.5	22.2	0.1	20.5	0.2	2.7	1.7
587736753546986772	251.6243	26.1073	24.1	0.9	24.7	0.5	24.2	0.5	21.9	0.1	20.5	0.1	2.3	1.4
587739706889012739	251.6538	12.3188	23.1	0.8	24.1	0.5	24.1	0.7	21.7	0.1	20.3	0.2	2.3	1.4
587736619330438661	251.7312	19.8500	24.9	0.8	25.1	0.5	25.0	0.5	22.3	0.2	20.6	0.1	2.7	1.7
587739720851982082	252.0974	12.7373	25.4	1.0	24.5	0.5	24.3	0.5	22.0	0.2	20.6	0.2	2.3	1.5
587736586054468875	252.2020	20.3096	25.3	0.7	24.9	0.6	23.6	0.4	21.0	0.1	19.7	0.1	2.5	1.4
587736980105659250	252.4118	45.8986	24.4	1.0	24.5	0.5	24.2	0.4	22.1	0.1	20.3	0.1	2.1	1.8
587725993039037495	252.4292	41.6562	25.7	0.8	24.6	0.5	24.2	0.6	21.1	0.1	19.7	0.1	3.1	1.4
587733603734193524	252.4306	32.8219	23.8	0.9	25.7	0.5	23.7	0.3	21.5	0.1	20.1	0.1	2.1	1.4
587735743695947666	252.6627	23.0738	25.2	0.8	25.0	0.4	25.1	0.5	22.1	0.1	20.7	0.2	2.9	1.4
587735665322427744	252.6977	29.8417	23.9	0.7	24.9	0.5	24.4	0.6	21.2	0.1	19.8	0.1	3.2	1.4
588007004197422275	252.7092	38.3522	25.1	0.8	25.5	0.5	23.7	0.4	21.3	0.1	20.0	0.1	2.4	1.3
587733603734521243	252.9105	32.0392	25.5	0.9	24.1	0.4	23.4	0.3	20.8	0.1	19.2	0.1	2.6	1.6

Table C.1 (cont'd)

Object ID	RA	DEC	u	u err	g	g err	r	r err	i	i err	\mathbf{Z}	z err	r-i	i-z
587739863639262265	252.9721	52.0506	25.5	0.8	24.8	0.5	23.9	0.4	21.4	0.1	20.0	0.1	2.5	1.4
587736976356214727	253.0261	16.0940	24.9	0.9	26.1	0.3	24.4	0.4	22.0	0.1	20.4	0.1	2.3	1.6
587725994112648394	253.2654	42.3323	24.9	1.1	25.6	0.5	24.6	0.7	21.8	0.2	20.4	0.2	2.8	1.4
587742783680611670	253.2765	51.3818	25.1	0.7	25.1	0.4	24.7	0.6	22.3	0.2	20.9	0.2	2.3	1.5
587725489986733093	253.2858	62.1073	25.4	0.8	24.8	0.7	23.5	0.4	20.9	0.1	19.5	0.1	2.6	1.4
587736919439181090	253.8167	19.2425	25.5	1.2	25.7	0.7	25.0	0.9	22.8	0.3	20.6	0.2	2.2	2.2
587736919439312192	253.9703	19.0003	24.1	1.7	23.8	0.3	24.9	0.9	21.9	0.1	20.4	0.1	3.1	1.5
588018090551084753	254.0996	21.3105	23.2	0.5	25.0	0.5	25.1	0.5	22.0	0.1	20.5	0.2	3.1	1.5
587729752746427740	254.2227	35.0509	25.4	0.6	24.8	0.5	24.8	0.5	22.5	0.2	20.9	0.2	2.4	1.6
588007005271754255	254.4336	37.5274	25.2	0.7	24.7	0.4	24.9	0.5	22.6	0.2	21.3	0.3	2.2	1.4
588018055130384012	254.4947	23.8064	25.4	0.8	24.6	0.4	25.0	0.5	22.4	0.1	20.8	0.2	2.7	1.5
587736586055714227	254.5579	18.4856	25.0	0.8	25.3	0.5	24.0	0.5	21.0	0.1	19.5	0.1	3.0	1.5
588007005271819804	254.5623	37.4297	24.0	0.7	24.9	0.5	24.9	0.5	22.0	0.1	20.5	0.2	3.0	1.4
587739848072365482	254.7056	41.8400	25.1	1.1	24.5	0.6	24.5	0.5	22.2	0.2	20.3	0.1	2.3	1.9
587742781535684101	254.9394	45.4067	25.3	0.8	24.9	0.6	24.2	0.5	21.9	0.1	20.2	0.1	2.3	1.7
587736618795468652	255.0328	16.8106	23.9	0.9	24.8	0.6	24.2	0.4	22.0	0.1	20.5	0.1	2.2	1.5
587736945742382695	255.3519	18.0737	24.5	1.0	24.3	0.4	24.1	0.5	21.1	0.1	19.7	0.1	3.0	1.4
587730842593919931	255.4256	75.2287	25.6	0.7	25.3	0.6	24.8	0.6	22.1	0.1	20.5	0.1	2.7	1.6
758882760135739040	255.5078	11.5037	23.1	0.8	25.5	0.6	25.3	0.7	21.9	0.2	20.6	0.3	3.4	1.4
587742837363049620	255.5588	44.5236	26.1	0.5	24.3	0.6	23.5	0.4	20.9	0.1	19.1	0.1	2.6	1.8
588018090015262523	255.7956	19.2954	24.5	0.9	25.4	0.5	25.0	0.6	22.5	0.2	21.1	0.3	2.4	1.4
587733604809508269	255.8353	30.3520	26.4	0.4	25.7	0.4	24.1	0.5	21.9	0.1	20.4	0.2	2.2	1.5
587736980644496959	255.8479	42.1073	24.6	1.0	24.9	0.5	24.1	0.4	21.9	0.1	20.5	0.2	2.2	1.3
588018253229852213	255.8540	22.7398	24.8	1.2	25.2	0.6	25.0	0.7	22.0	0.2	20.5	0.2	2.9	1.5
587736945742710257	255.9018	17.5926	24.9	1.0	24.4	0.4	24.3	0.5	22.0	0.1	20.4	0.1	2.3	1.6
587729751137716155	255.9747	30.7506	24.4	0.9	25.2	0.5	24.0	0.4	21.9	0.1	20.5	0.1	2.1	1.4
587739862031861165	256.1489	44.8678	25.3	0.9	25.3	0.4	24.5	0.5	21.9	0.1	20.6	0.2	2.6	1.4

Table C.1 (cont'd)

Object ID	RA	DEC	u	u err	g	g err	r	r err	i	i err	\mathbf{z}	z err	r-i	i-z
587729781199799884	256.1800	33.6150	23.7	0.8	25.6	0.4	24.7	0.5	21.9	0.1	20.5	0.1	2.8	1.4
587729409155925388	256.5650	35.3264	23.8	0.7	24.6	0.5	24.4	0.6	22.3	0.2	20.9	0.2	2.2	1.4
758879771378845842	256.6357	11.7766	26.0	0.5	25.6	0.5	24.8	0.6	21.9	0.1	20.5	0.1	2.9	1.5
588018253766985235	256.6869	22.5763	24.9	1.3	24.0	0.4	24.4	0.7	21.4	0.1	19.6	0.1	3.0	1.8
587739850219849059	256.6912	42.5799	26.2	0.4	24.8	0.5	24.2	0.5	21.8	0.1	20.2	0.1	2.5	1.5
587736752476587872	256.7066	20.0545	25.2	0.9	25.6	0.5	24.6	0.6	22.3	0.2	21.0	0.2	2.2	1.4
758879771915716733	256.7997	12.0263	25.9	0.8	25.0	0.5	24.8	0.6	22.3	0.2	20.9	0.2	2.5	1.4
587725491597870175	256.8440	61.5855	25.4	0.9	24.6	0.5	24.0	0.5	21.8	0.2	20.5	0.2	2.2	1.4
587729409156253046	256.9946	34.7249	24.3	0.8	24.5	0.5	23.6	0.4	20.9	0.1	19.5	0.1	2.7	1.4
587725489989026939	257.1205	57.2346	25.9	0.5	24.3	0.6	25.1	0.7	22.1	0.2	20.5	0.2	3.0	1.6
587730842058425492	257.2105	72.1206	24.6	1.2	24.7	0.5	24.6	0.6	22.3	0.2	21.0	0.2	2.3	1.4
587733398653502709	257.3226	23.5716	25.9	1.1	25.6	0.7	24.3	0.9	22.0	0.2	20.2	0.2	2.3	1.8
587729781200717393	257.4297	31.7636	25.0	1.0	24.4	0.4	24.9	0.5	22.0	0.1	20.6	0.1	2.9	1.4
758882760673462010	257.4823	11.0434	25.2	1.2	25.0	0.8	24.3	0.7	22.0	0.2	20.7	0.2	2.3	1.4
758882759599916688	257.5127	10.2060	25.9	0.8	24.7	0.7	24.2	0.7	21.7	0.2	20.2	0.2	2.5	1.5
587739850220438969	257.5187	41.2803	25.9	0.5	24.8	0.5	24.4	0.5	22.3	0.2	20.6	0.1	2.1	1.7
758882758526371526	257.5470	9.1500	25.1	1.4	25.5	0.9	24.6	0.9	21.6	0.2	20.1	0.1	3.0	1.5
587733431926130522	257.6034	22.1850	23.9	0.5	25.2	0.5	23.9	0.4	21.1	0.1	19.6	0.1	2.8	1.5
587746214818612526	257.7310	36.2683	25.8	1.0	24.3	0.5	23.6	0.4	21.0	0.1	19.7	0.1	2.6	1.3
587736980109329796	257.7816	38.6284	24.6	0.9	25.2	0.5	23.7	0.4	21.5	0.1	20.2	0.1	2.2	1.4
587725491598591049	257.9375	59.9488	24.9	1.1	25.1	0.6	24.5	0.6	21.7	0.1	20.3	0.2	2.8	1.4
587742632274429117	258.1832	42.0965	26.1	0.6	24.5	0.5	24.9	0.7	22.5	0.2	21.0	0.2	2.4	1.5
587746214282396851	258.1880	34.6982	25.1	1.3	24.6	0.6	23.9	0.6	21.1	0.1	19.8	0.1	2.8	1.4
587729409157236396	258.2269	32.7377	25.8	0.5	25.2	0.5	24.7	0.6	22.5	0.2	21.0	0.2	2.3	1.4
758882759063504941	258.3772	9.3276	25.6	1.2	25.5	0.7	24.5	0.9	22.1	0.2	20.6	0.2	2.5	1.4
587733604811409067	258.5882	26.7106	25.3	1.1	23.6	0.3	23.8	0.4	21.6	0.1	20.2	0.1	2.2	1.4
587729652355368489	258.6758	27.2485	26.0	0.6	24.7	0.6	23.6	0.4	21.4	0.1	19.8	0.1	2.3	1.5

Table C.1 (cont'd)

Object ID	RA	DEC	u	u err	g	g err	r	r err	i	i err	\mathbf{z}	z err	r-i	i-z
758882760674183086	258.9452	10.2856	25.1	1.2	24.6	0.7	24.5	0.8	21.9	0.2	20.3	0.1	2.5	1.6
587729782275049006	259.1334	31.0371	25.4	0.8	24.3	0.4	24.6	0.5	22.5	0.2	21.0	0.3	2.1	1.4
758882758527092812	259.1883	8.4838	26.8	0.3	25.3	0.9	24.9	0.8	22.1	0.2	20.6	0.2	2.8	1.5
587739849149187654	259.5068	35.7880	23.0	0.5	25.6	0.5	24.4	0.5	22.1	0.1	20.8	0.2	2.3	1.3
587742632275412364	259.5289	40.0295	24.9	1.1	24.4	0.5	25.6	0.5	22.5	0.2	21.1	0.2	3.1	1.4
587742633348367682	259.6630	42.0777	25.4	0.8	25.6	0.5	24.8	0.7	22.4	0.2	21.0	0.2	2.4	1.4
587742632275608757	259.7201	39.6730	23.3	0.6	24.7	0.5	23.9	0.5	21.0	0.1	19.4	0.1	2.9	1.5
587729751677143147	259.7740	26.0275	24.9	0.8	24.8	0.4	24.3	0.4	21.9	0.1	20.6	0.1	2.4	1.3
587739850222536105	259.9737	36.9954	25.4	0.7	25.8	0.4	24.4	0.5	22.2	0.2	20.7	0.2	2.2	1.5
587729652892829434	259.9849	26.2638	23.3	0.6	25.3	0.6	23.8	0.5	21.5	0.1	20.1	0.1	2.4	1.4
758882761211381701	259.9916	10.4703	22.9	0.7	25.6	0.7	24.3	0.7	21.4	0.1	20.0	0.1	2.9	1.4
758879770843744759	260.1522	9.6868	22.4	0.3	25.4	0.6	24.9	0.6	21.9	0.1	20.5	0.2	3.0	1.4
758879770307005127	260.2121	9.0591	26.0	0.7	24.5	0.5	25.0	0.6	22.2	0.2	20.2	0.1	2.9	2.0
587729408622397293	260.3897	28.4186	24.8	0.9	25.3	0.5	23.6	0.4	21.0	0.1	19.4	0.1	2.6	1.6
587739863108355554	260.6001	39.5375	25.3	0.8	25.0	0.5	23.9	0.4	21.5	0.1	20.0	0.1	2.4	1.5
588011502062404545	260.7412	70.5535	24.1	1.4	23.6	0.4	24.1	0.6	21.5	0.1	20.2	0.2	2.6	1.3
587729781203208138	260.7426	26.8782	23.4	0.6	24.9	0.6	23.8	0.4	21.3	0.1	19.8	0.1	2.5	1.5
758882758527879636	260.7825	7.6438	23.6	0.9	25.4	0.8	25.2	0.7	22.3	0.2	20.5	0.2	2.9	1.8
758882760675035440	260.7964	9.5815	25.1	1.1	25.4	0.6	25.3	0.7	22.1	0.2	20.7	0.2	3.3	1.4
587729408086182141	260.8324	26.8894	24.4	0.7	24.6	0.5	24.7	0.6	22.1	0.1	20.6	0.1	2.6	1.5
758879769233853381	261.1184	7.7240	26.1	0.6	25.5	0.8	24.1	0.6	21.9	0.1	20.4	0.1	2.2	1.5
587729408086378539	261.2006	26.5870	24.2	0.7	24.7	0.5	24.5	0.6	22.2	0.2	20.5	0.1	2.3	1.7
587729781740144448	261.2399	26.9407	23.9	0.8	24.0	0.4	24.6	0.6	22.0	0.1	20.6	0.2	2.6	1.3
587725491064931490	261.2604	52.8318	24.5	0.9	26.0	0.4	24.1	0.6	21.9	0.2	20.4	0.2	2.2	1.5
587729408086837186	261.7103	25.7472	23.7	0.6	25.1	0.5	24.5	0.6	21.7	0.1	20.3	0.1	2.8	1.4
587746215358432840	261.7363	30.4520	24.1	1.1	25.3	0.6	25.8	0.8	22.5	0.3	20.8	0.2	3.3	1.7
587742632277640802	261.9570	35.2503	25.5	0.8	24.0	0.3	25.0	0.6	22.3	0.2	20.8	0.2	2.8	1.5

Table C.1 (cont'd)

Object ID	RA	DEC	u	u err	g	g err	r	r err	i	i err	\mathbf{Z}	$z \ err$	r-i	i-z
587746214285739793	262.2928	27.8783	25.9	0.6	25.5	0.5	25.1	0.6	22.4	0.2	20.9	0.2	2.7	1.5
587739850224633642	262.3707	32.6224	22.8	0.4	24.5	0.4	24.2	0.4	21.3	0.1	19.8	0.1	2.9	1.5
587725577499116558	262.3744	59.8877	24.4	1.4	24.5	0.6	24.2	0.4	22.0	0.2	20.5	0.2	2.3	1.5
758877291036084359	262.9636	25.3941	25.3	0.7	25.4	0.7	23.8	0.4	21.4	0.1	19.9	0.1	2.4	1.5
758882760139344771	263.1063	7.9430	24.2	1.4	25.0	0.7	24.6	0.7	21.6	0.1	20.1	0.2	3.1	1.5
758882758529125574	263.3234	6.3313	25.4	1.1	25.6	0.9	23.6	0.4	21.2	0.1	19.7	0.1	2.4	1.5
587739862038087456	263.3453	31.8399	24.7	1.0	24.7	0.5	24.5	0.5	22.1	0.1	20.7	0.2	2.4	1.4
758882761213086475	263.5151	8.7344	25.1	1.4	25.9	0.6	24.8	0.7	22.5	0.3	20.5	0.2	2.3	2.1
587742633352103738	263.5610	34.1023	23.6	0.6	24.6	0.5	25.1	0.5	22.6	0.2	21.1	0.2	2.6	1.4
758877273857657029	263.6971	25.2289	25.3	0.8	25.1	0.7	25.5	0.5	22.5	0.2	20.9	0.2	3.0	1.7
587742633352300273	263.6975	33.6899	23.7	0.7	24.5	0.5	25.0	0.6	22.4	0.2	21.0	0.2	2.6	1.4
758877292110088500	263.8506	25.8400	24.6	0.9	25.2	0.5	23.4	0.3	21.0	0.1	19.5	0.1	2.4	1.5
587742632279541546	263.9289	31.3187	24.2	0.8	25.5	0.5	24.8	0.5	21.8	0.1	20.3	0.1	2.9	1.6
758877275468532229	264.8312	26.2702	25.5	0.7	25.1	0.5	24.5	0.6	22.3	0.2	20.9	0.3	2.2	1.4
758877276542208313	264.9979	27.2111	25.3	0.8	25.1	0.5	23.7	0.4	21.2	0.1	19.8	0.1	2.5	1.4
587734174952326699	265.1735	49.7183	24.9	1.2	25.1	0.6	24.6	0.3	22.1	0.1	20.4	0.1	2.5	1.7
758877291037002069	265.2069	24.7046	25.2	0.8	25.9	0.6	23.6	0.3	21.1	0.1	19.6	0.1	2.6	1.5
587734175489459584	266.0173	49.1882	24.5	0.9	25.0	0.5	23.8	0.4	20.8	0.0	19.2	0.0	3.0	1.7
587730842601129417	266.3485	59.3150	24.0	0.8	24.5	0.5	24.1	0.4	21.4	0.1	20.1	0.1	2.7	1.3
587734176026527402	266.9295	48.8554	25.1	0.9	25.4	0.5	25.1	0.6	22.5	0.2	21.1	0.3	2.7	1.4
587734174955341054	267.1264	42.9341	25.5	1.1	25.5	0.9	24.0	0.6	21.3	0.1	19.8	0.1	2.7	1.5
587730807702357170	267.3163	41.1764	24.4	1.7	26.0	0.6	23.8	0.7	21.6	0.1	20.1	0.1	2.2	1.6
758877293185140800	267.3254	25.7639	24.9	0.9	25.0	0.5	24.6	0.6	21.9	0.1	20.5	0.1	2.6	1.5
588011502068368684	267.4633	57.2745	23.3	0.6	25.3	0.6	24.3	0.6	22.2	0.2	20.6	0.2	2.2	1.5
758883006032184816	267.8302	65.6736	23.4	0.8	25.3	0.5	24.7	0.6	21.6	0.1	20.2	0.1	3.1	1.3
587734176564249664	267.9171	46.9716	25.2	0.9	24.1	0.5	23.6	0.4	21.1	0.1	19.7	0.1	2.6	1.4
587730809846891561	267.9193	48.0153	25.6	1.1	24.9	0.8	23.7	0.6	21.1	0.1	19.7	0.1	2.6	1.4

Table C.1 (cont'd)

Object ID	RA	DEC	u	u err	g	g err	r	r err	i	i err	\mathbf{Z}	z err	r-i	i-z
758877291575052929	267.9698	24.1137	25.1	1.1	25.8	0.5	25.0	0.6	22.7	0.2	20.7	0.2	2.3	2.0
588011502069810430	268.1147	54.0275	25.1	1.2	25.6	0.5	24.1	0.5	21.6	0.1	19.9	0.1	2.5	1.7
758874294221145127	268.2962	78.2125	23.4	0.9	24.7	0.6	24.1	0.6	21.8	0.1	20.3	0.1	2.3	1.4
588011503140930769	268.3149	59.9652	25.1	1.3	24.9	0.7	24.3	0.7	21.8	0.2	20.4	0.2	2.6	1.4
758883003347961142	268.4961	63.6166	24.8	1.0	24.8	0.6	24.5	0.5	22.2	0.2	20.7	0.2	2.3	1.5
587734176029869411	268.6625	41.3388	25.5	1.2	25.4	0.7	24.1	0.8	21.9	0.2	20.5	0.2	2.2	1.4
758877291038575154	268.7399	23.4486	24.4	1.0	25.5	0.7	25.5	0.5	22.1	0.2	20.5	0.1	3.4	1.6
587730808777082334	268.8786	39.0188	24.3	1.5	25.6	0.7	24.6	0.8	21.9	0.2	20.5	0.3	2.7	1.5
588011503142372743	268.9401	56.7506	25.8	0.8	24.2	0.4	23.3	0.3	21.0	0.1	19.6	0.1	2.2	1.5
587734176030459448	269.0289	39.8995	25.4	1.1	24.5	0.6	24.3	0.7	22.2	0.2	20.5	0.3	2.1	1.6
758879801444206165	269.0647	44.1647	24.6	1.1	25.0	0.6	25.2	0.5	22.5	0.2	21.0	0.2	2.7	1.5
758877274933562048	269.1480	24.3495	25.2	0.8	24.8	0.5	23.7	0.4	21.3	0.1	19.9	0.1	2.4	1.4
588011502071776366	269.3381	49.6072	23.1	0.7	24.7	0.7	24.0	0.6	21.2	0.1	19.8	0.1	2.8	1.4
758879800370595475	269.3537	43.1697	23.8	0.9	24.5	0.5	25.2	0.6	22.2	0.2	20.2	0.1	2.9	2.1
587734176566937094	269.4759	40.9627	24.8	1.2	25.2	0.7	25.5	0.8	22.7	0.4	20.6	0.2	2.8	2.0
588011503143290028	269.5050	54.7412	24.6	1.2	25.3	0.7	23.3	0.3	21.1	0.1	19.7	0.1	2.2	1.4
758877274933758866	269.6641	24.2277	24.8	0.9	25.6	0.5	24.4	0.4	21.8	0.1	20.5	0.1	2.5	1.4
758877292649450148	269.9273	24.4050	24.5	1.0	24.2	0.4	24.2	0.5	21.8	0.1	20.5	0.2	2.4	1.4
758877274934086491	270.3753	24.0512	23.9	0.8	25.0	0.6	24.7	0.6	21.5	0.1	19.7	0.1	3.3	1.8
588011503145452721	270.5323	49.7575	25.2	1.3	25.9	0.6	25.1	0.8	21.9	0.2	20.2	0.2	3.2	1.7
758879743999608647	270.5526	42.5107	24.8	1.1	25.1	0.5	24.1	0.4	21.7	0.1	20.3	0.1	2.4	1.4
758877274934217530	270.7269	23.9671	25.4	0.8	25.2	0.6	24.7	0.6	21.6	0.1	20.2	0.1	3.0	1.4
588011503146042364	270.7437	48.5025	23.8	1.3	26.2	0.6	24.1	0.7	22.0	0.2	20.3	0.2	2.2	1.6
758877276007959231	270.9368	24.6711	25.7	0.6	24.5	0.5	24.1	0.5	21.5	0.1	19.9	0.1	2.5	1.6
758882837984248917	271.6909	64.9125	25.4	1.0	25.2	0.7	23.8	0.5	21.6	0.1	20.3	0.1	2.2	1.3
758879799834708180	272.1941	42.3734	24.3	1.1	24.6	0.4	24.5	0.5	21.8	0.1	20.4	0.1	2.7	1.4
758882836910703914	272.5915	64.0585	25.6	0.9	25.5	0.5	24.2	0.6	22.0	0.2	20.4	0.2	2.3	1.5

Table C.1 (cont'd)

Object ID	RA	DEC	u	u err	g	g err	r	r err	i	i err	\mathbf{z}	z err	r-i	i-z
758879800371906760	273.3558	42.5621	25.2	0.8	25.1	0.5	24.1	0.5	21.9	0.1	20.5	0.1	2.2	1.4
758879745611073714	273.6150	43.1534	23.9	0.7	25.2	0.4	23.8	0.3	21.4	0.1	19.5	0.1	2.3	1.9
758879800372299970	274.5077	42.2531	25.4	0.8	24.5	0.5	24.4	0.5	21.6	0.1	19.6	0.1	2.8	2.0
758879800372365338	274.6957	42.3814	25.6	0.7	25.1	0.6	23.9	0.4	21.3	0.1	19.8	0.1	2.6	1.5
758882835837617362	275.8669	63.1277	26.0	0.7	25.2	0.8	25.0	0.8	22.2	0.2	20.7	0.2	2.8	1.5
758879800372824046	276.0684	41.8951	23.3	0.6	24.8	0.5	24.8	0.6	22.2	0.2	20.7	0.2	2.6	1.5
758879801983436639	276.4392	43.1702	25.8	0.9	25.3	0.5	24.5	0.6	22.0	0.2	20.6	0.2	2.5	1.3
758883003886470770	276.9034	63.6557	24.5	1.2	24.5	0.5	24.2	0.5	21.7	0.1	20.3	0.1	2.5	1.4
758879745075448276	276.9994	41.9097	25.2	0.8	25.2	0.4	24.1	0.4	21.4	0.1	19.8	0.1	2.8	1.5
758879801983633403	277.1004	43.0152	23.6	0.9	25.3	0.5	23.9	0.5	20.9	0.1	19.2	0.0	3.0	1.7
758879745075776134	278.0546	41.7461	25.4	0.7	25.3	0.4	24.1	0.4	21.5	0.1	19.7	0.1	2.6	1.8
758879799299934602	278.2223	40.5655	23.8	0.9	24.4	0.6	24.4	0.6	21.7	0.1	20.3	0.1	2.7	1.4
75887979930000074	278.3825	40.4587	24.6	1.2	26.0	0.6	25.6	0.5	22.6	0.3	20.5	0.1	3.1	2.1
758879801447483544	278.9912	42.1108	24.9	1.1	25.1	0.6	24.9	0.6	22.5	0.2	20.9	0.2	2.5	1.5
758879801984551053	279.8073	42.3506	24.5	1.4	24.9	0.6	25.1	0.6	22.1	0.2	20.7	0.2	3.0	1.4
758882836376455142	285.5722	62.6958	25.0	1.5	24.8	0.5	24.3	0.6	22.1	0.2	20.5	0.1	2.2	1.6
758883004961916696	285.9520	63.8609	26.2	0.6	25.1	0.5	24.8	0.5	21.5	0.1	20.1	0.1	3.3	1.5
758883006035724115	286.5770	64.5230	25.7	0.9	25.7	0.4	25.2	0.6	22.8	0.3	21.0	0.2	2.3	1.9
758882837987198851	287.0166	63.8325	25.3	0.8	24.2	0.4	23.6	0.3	21.2	0.1	19.6	0.1	2.4	1.6
758874337169704032	288.4135	77.8720	25.7	1.0	24.7	0.7	24.0	0.7	21.8	0.2	20.4	0.2	2.1	1.5
758874337706574831	288.6738	78.3016	25.6	1.1	24.5	0.7	23.8	0.6	21.6	0.1	20.2	0.2	2.2	1.4
758883034478347963	289.1602	62.7378	24.3	1.3	25.3	0.6	24.0	0.5	20.9	0.1	19.5	0.1	3.1	1.3
758874294759851155	289.4078	78.5075	25.7	1.2	25.2	0.7	24.0	0.7	21.2	0.1	19.8	0.1	2.8	1.4
758883004425832388	289.4148	62.8697	25.4	0.9	24.5	0.5	24.5	0.5	21.4	0.1	19.9	0.1	3.1	1.5
758874293686305978	289.9318	77.5889	24.1	1.4	24.5	0.6	25.2	0.9	22.2	0.2	20.5	0.2	3.0	1.6
758883033404999281	290.2589	61.6931	23.9	1.0	24.5	0.7	25.1	0.6	22.1	0.2	20.8	0.2	3.0	1.3
758874293149500612	290.7478	77.2522	26.5	1.0	25.1	0.8	25.1	0.9	21.9	0.2	20.3	0.2	3.2	1.6

Table C.1 (cont'd)

Object ID	RA	DEC	u	u err	g	g err	r	r err	i	i err	\mathbf{z}	z err	r-i	i-z
758874293149828414 758874294760899955	$293.8706 \\ 300.5010$	$76.9509 \\ 77.8988$	$24.8 \\ 25.7$	$\begin{array}{c} 1.2 \\ 0.8 \end{array}$	$24.9 \\ 25.4$	$\begin{array}{c} 0.7 \\ 0.6 \end{array}$	$25.4 \\ 25.1$	$\begin{array}{c} 0.7 \\ 0.7 \end{array}$	$22.5 \\ 22.0$	$\begin{array}{c} 0.3 \\ 0.2 \end{array}$	$20.6 \\ 20.6$	$0.2 \\ 0.2$	$2.9 \\ 3.0$	$1.9 \\ 1.5$

Appendix D

EXPANSION ON THE EXPECTED HVS CALCULATIONS.

Here we expand on the calculations outlined in section 3.6.4, determining the expected number of HVSs originating from an interaction with a central SMBH assuming varying initial mass functions.

We start with the Salpeter IMF, given by the equation:

$$\frac{dN}{dM} = \int_{M_1}^{M_2} M^{-2.35} \tag{4.1}$$

where $\frac{dN}{dM}$ is the number of stars in a small mass range and M_1 is the lower mass limit and M_2 is the upper mass limit. If we follow exactly the prediction from Section 3.6.4, considering the 0.6 – 1.2 solar mass range for G/K-type stars and the 3 – 4 solar mass range for the known B-type HVS, we find:

$$\frac{N_B}{N_{G/K}} = \frac{\int_3^4 M^{-2.35}}{\int_{0.6}^{1.2} M^{-2.35}} = \frac{3^{-1.35} - 4^{-1.35}}{0.6^{-1.35} - 1.2^{-1.35}} = \frac{0.7}{1.2}$$
(4.2)

Once scaled for the 14 known B-type HVS that were also detected by SDSS, we get that the expected number of G/K-type HVS is roughly 240.

However, since we were uncertain of the exact masses for the confirmed high mass HVS, we instead assumed an average mass of 3.5 M_{\odot} to determine $\frac{dN}{dM}$, the number of stars in a 1 dex mass bin around this average, and to be more in line with the approach of Kollmeier et al. (2010). Similarly, at the low mass end of the IMF low mass stars dominate

significantly, and we assume that the contribution from 1.2 M_{\odot} stars is negligible. In this case, we find:

$$\frac{N_B}{N_{G/K}} = \frac{3.5^{-1.35}}{0.6^{-1.35}} \tag{4.3}$$

Scaling for the 14 known B-type HVSs, yields the predicted ${\sim}150$ G/K-type HVSs stated in Section 3.6.4.

We make similar assumptions for the top-heavy IMF (Figer et al., 1999):

$$\frac{dN}{dM} = \int_{M_1}^{M_2} M^{-1.6} = \frac{3.5^{-0.6}}{0.6^{-0.6}} \tag{4.4}$$

Assuming an average B-type mass of 3.5 M_{\odot} and a dominating contribution from 0.6 M_{\odot} stars, we get the expected 40 G/K-type HVSs.

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